

# Les fluctuations de courant : au delà du bruit

B. Reulet (1,2), D. Prober (1), J. Gabelli (2)

(1) Yale University – Department of Applied Physics (USA)

(2) Laboratoire de Physique des Solides – Orsay (France)

Introduction:

what to measure ?

# Averaged quantities

1) STATIC (thermodynamic or steady state) properties:

$\vec{M}(\vec{H})$       Magnetization vs. Magnetic field

$I(V)$       Current vs. Voltage

$\vec{n}(\vec{E})$       Molecular orientation vs. Electric field

$x(F)$       Displacement vs. Force

X ray diffraction vs. V,P, etc...

# Susceptibilities

2) DYNAMICAL properties: response to an oscillating field

$$H = H_{dc} + \delta H \cos \omega_0 t \quad \chi(H_{dc}, \omega_0) = \frac{\partial M_{\omega_0}}{\partial H_{\omega_0}} \quad \text{Magnetic susceptibility}$$
$$V = V_{dc} + \delta V \cos \omega_0 t \quad G(V_{dc}, \omega_0) = \frac{\partial I_{\omega_0}}{\partial V_{\omega_0}} \quad \text{ac Conductance}$$

Both are COMPLEX quantities: in-phase and out-of-phase response.

Interesting when:

$$\omega_0 \tau > 1$$

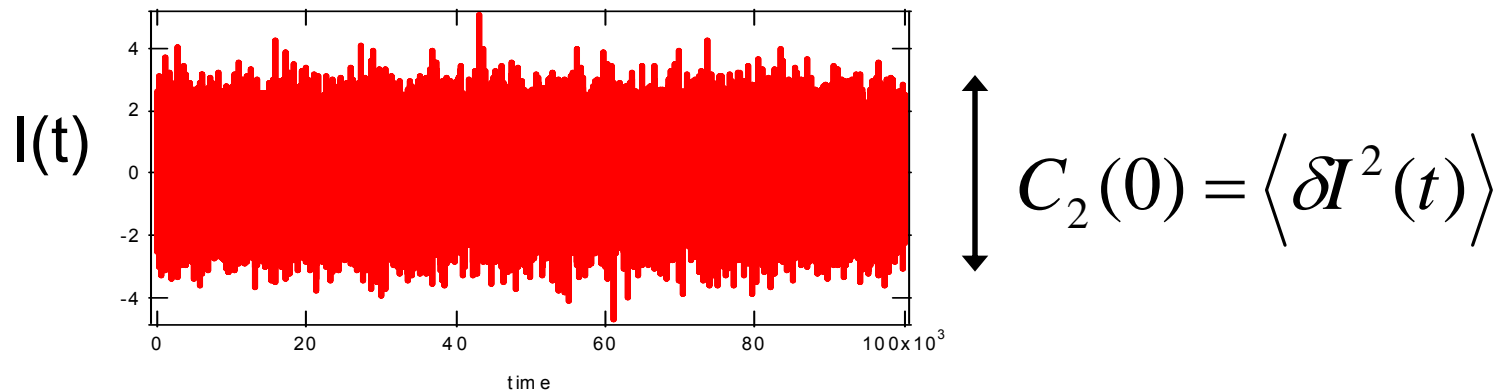
Aging of polymers or spin glasses, dielectric constant or magnetic permittivity, etc.

# Fluctuations (noise)

## 3) Intrinsic FLUCTUATIONS (or noise):

Correlation function:  $C_2(V, \tau) = \langle \delta I(t + \tau) \delta I(t) \rangle$  Time average ↙

Noise spectral density:  $S_2(V, \omega) = \langle \delta I(\omega) \delta I(-\omega) \rangle$



Magnetization noise (motion of domain walls),  $1/f$  noise of electric dipoles in amorphous systems glasses, of voltage in transistors, electric noise associated with depinning of charge density waves, etc.

# Summary

Average

$$\langle \bullet \rangle$$

Fluctuations

$$\langle \bullet \bullet \rangle$$

$$I(V)$$

$$S_2(V, \omega) = \langle \delta I(\omega) \delta I(-\omega) \rangle$$

Dynamical  
response

$$\frac{\partial \bullet}{\partial V_{\omega_0}}$$

$$G(V, \omega_0) = \frac{\partial I_{\omega_0}}{\partial V_{\omega_0}}$$

# Beyond

	$\langle \bullet \rangle$	$\langle \bullet \bullet \rangle$	$\langle \bullet \bullet \bullet \rangle$
	$I(V)$	$\langle \delta I(\omega) \delta I(-\omega) \rangle$	$\langle \delta I(\omega) \delta I(\omega') \delta I(-\omega - \omega') \rangle$
$\frac{\partial \bullet}{\partial V_{\omega_0}}$	$\frac{\partial I_{\omega_0}}{\partial V_{\omega_0}}$	$\frac{\partial S_2(\omega)}{\partial V_{\omega_0}}$	<p>THIRD CUMULANT</p> <p>↑</p>
$\frac{\partial^2 \bullet}{\partial V_{\omega_0} \partial V_{\omega_1}}$	$\frac{\partial^2 I_{\omega_0 \pm \omega_1}}{\partial V_{\omega_0} \partial V_{\omega_1}}$		

NOISE DYNAMICS

Frequency mixing (non linear transport)

The Dynamics of noise:

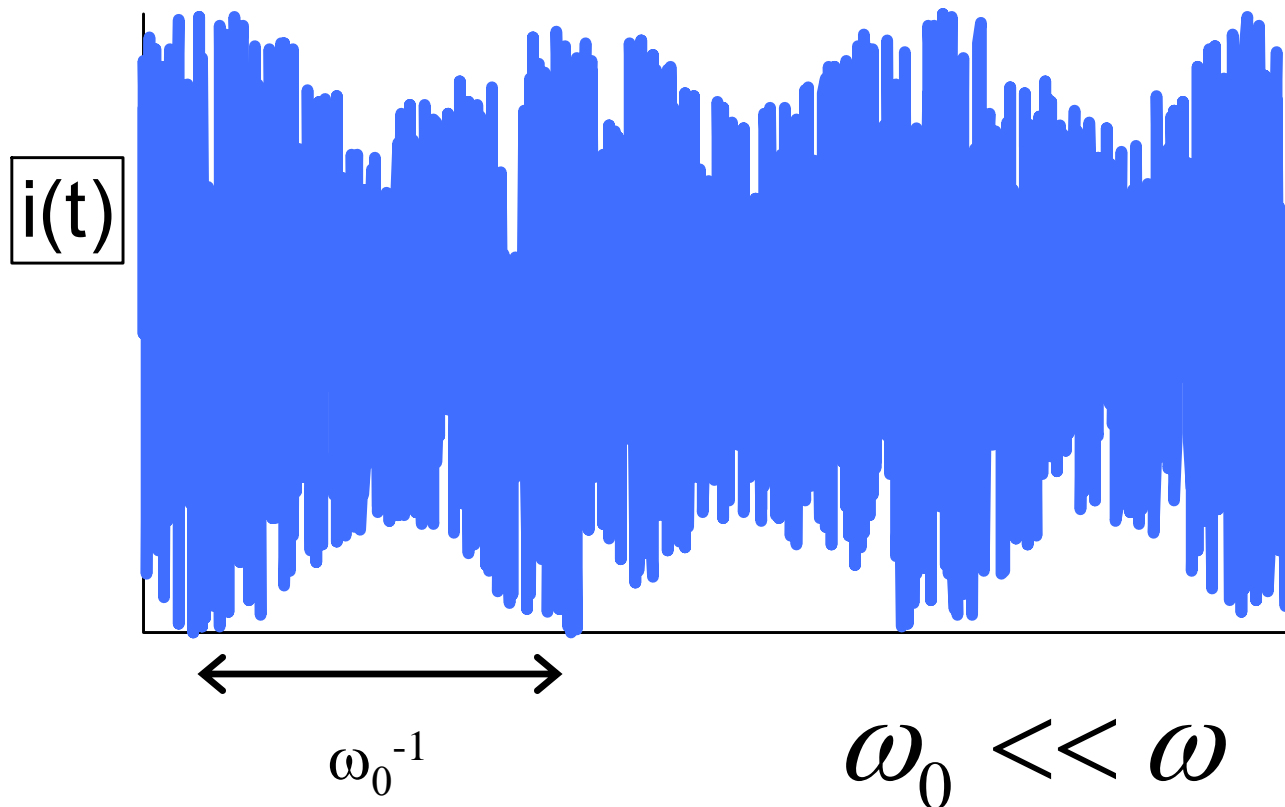
The noise susceptibility  
in the  
classical regime



# Noise susceptibility – How fast can one modulate noise ?

$$V(t) = V_{dc} + \delta V \cos \omega_0 t$$

$$\frac{\partial S_2(\omega)}{\partial V_{\omega_0}}$$



# Noise susceptibility – the case of a macroscopic conductor

Fluctuation-dissipation theorem, at equilibrium and low frequency  $\hbar\omega \ll k_B T$  :

$$S_2 = 4k_B T G$$

NOISE = electron THERMOMETER

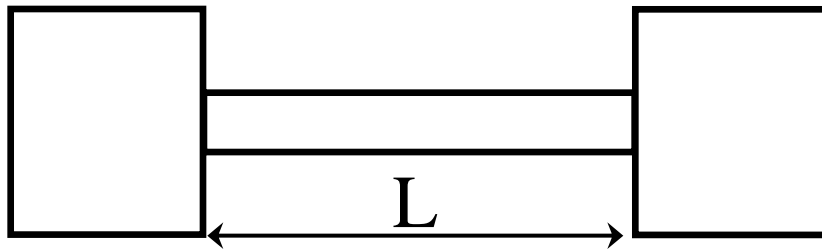
$$\delta V(t) \Rightarrow \delta P_{Joule}(t) = I \delta V(t) \Rightarrow \delta T(t) \Rightarrow \delta S_2(t)$$

$$\chi_{\omega_0}(\omega) = \frac{\partial S_2(\omega)}{\partial V_{\omega_0}} \propto Z_{thermal}^{-1}(\omega_0) \quad \omega \gg \omega_0$$

Its frequency dependence gives ENERGY RELAXATION (i.e. INELASTIC) time

# Noise susceptibility – from macro- to mesoscopic conductor

diffusive metallic wire: length  $L \gg$  mean free path



$$L^2 = D \tau_D$$

\* long wire or SNS: phonon cooling

The noise susceptibility gives the ELECTRON-PHONON time

\* intermediate wire: diffusion cooling

The noise susceptibility gives the DIFFUSION time

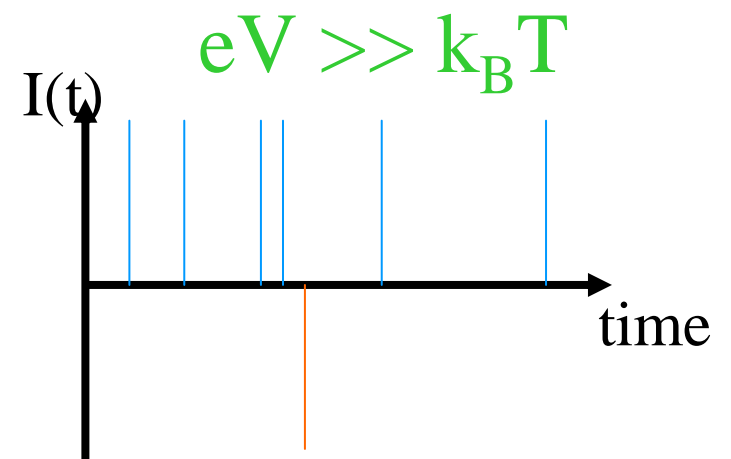
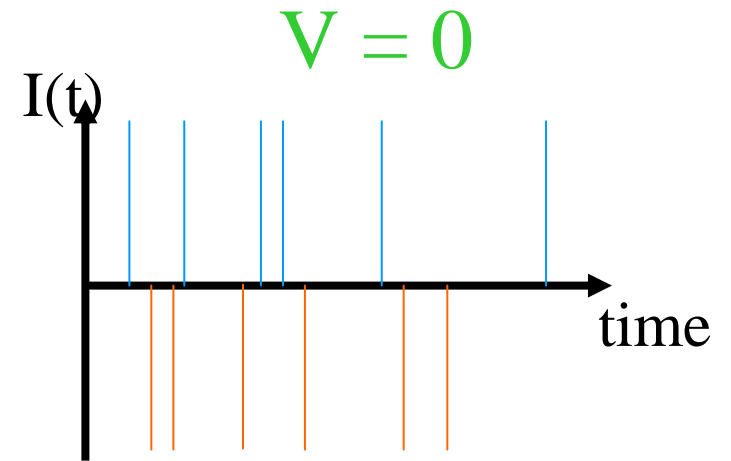
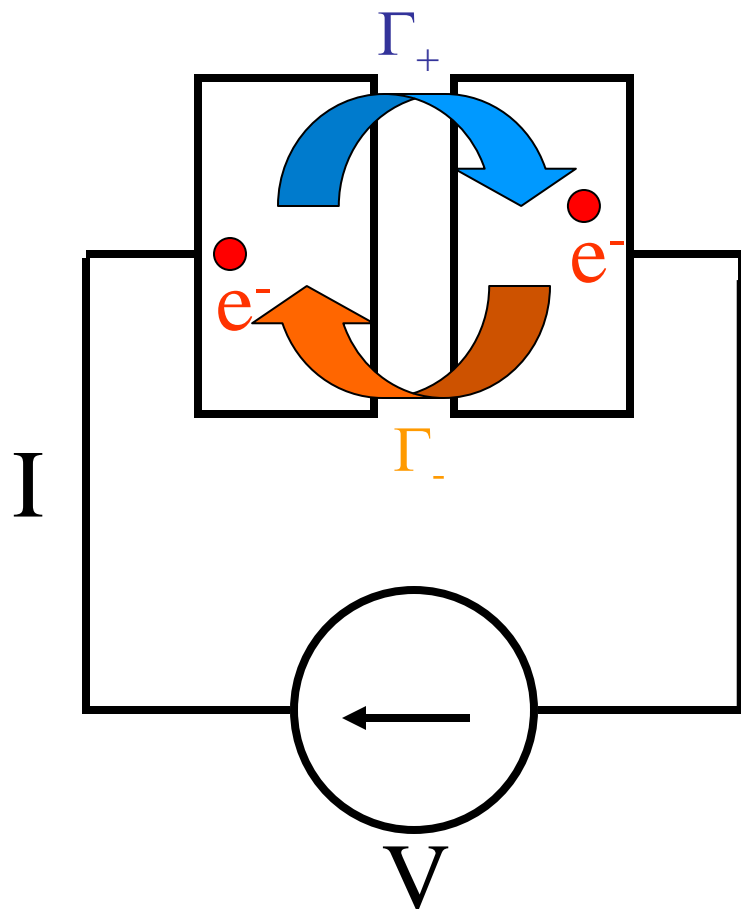
\* short wire: elastic transport (independent electrons)

The noise susceptibility gives the ELECTRON-ELECTRON time

\* ballistic wire (nanotube), quasi-crystals (sub-diffusive), ... ??

The noise susceptibility  
in the  
Quantum regime

# Current in a tunnel junction



# Current fluctuations in a tunnel junction at low frequency

$$\langle \delta I^2 \rangle = 2eIB \coth\left(\frac{eV}{2k_B T}\right) \quad \text{B=bandwidth}$$

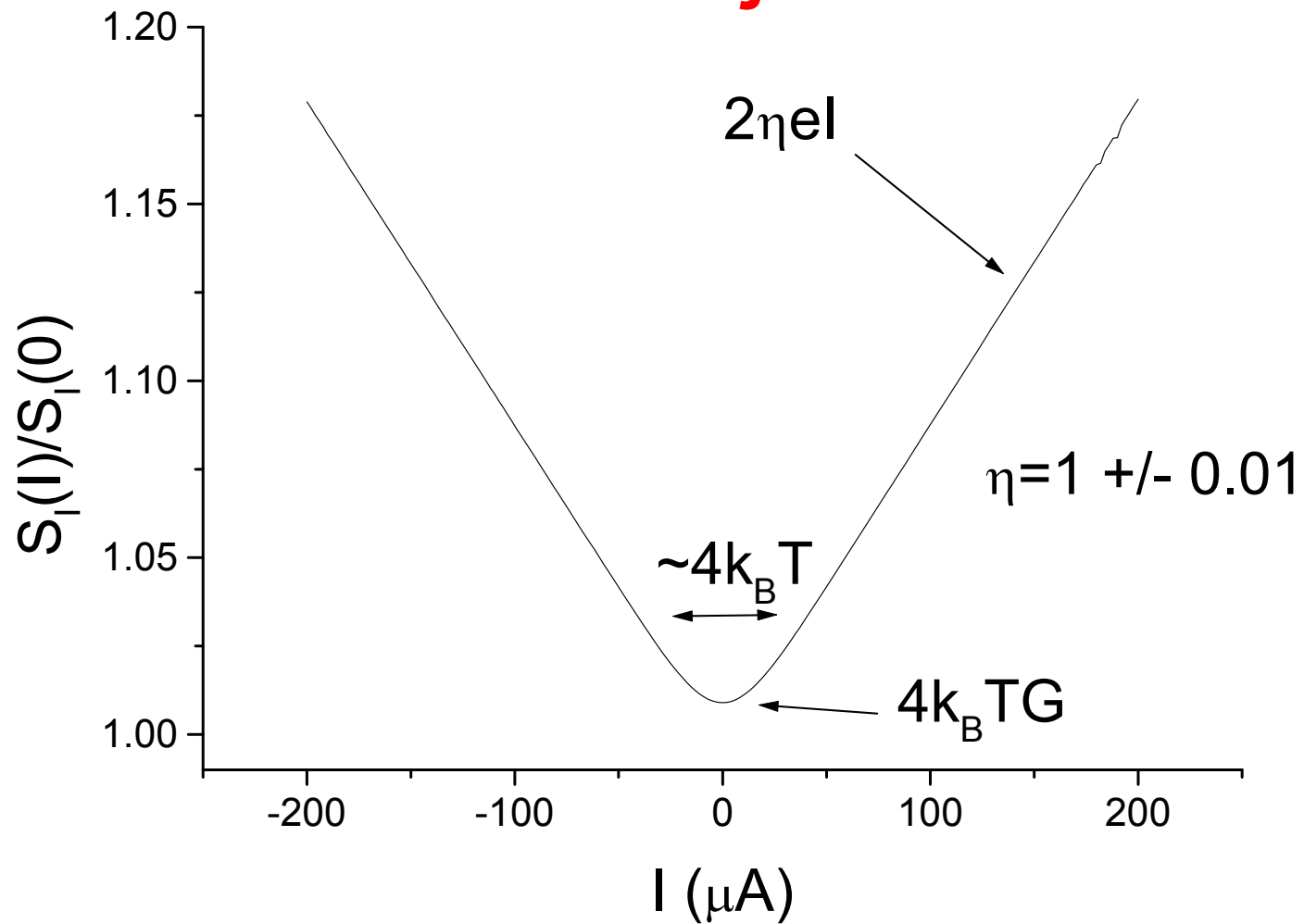
$$S_2 = \begin{cases} 4k_B T G & \text{if } eV \ll k_B T \\ 2eI & \text{if } eV \gg k_B T \end{cases}$$

Noise spectral density in A<sup>2</sup>/Hz

Equilibrium (Johnson) noise: macroscopic, fluctuation-dissipation theorem

Shot noise: discreteness of charge

# Measurement of $S_2$ on a tunnel junction



# Quantum mechanics: ordering of operators?

Average current:

$$I_{DC} = \langle \hat{I} \rangle$$

Noise  $S_2$ :

$$S_2(\omega) = \int dt e^{i\omega t} \left\{ \begin{array}{l} \langle \hat{I}(0) \hat{I}(t) \rangle \\ \langle \hat{I}(t) \hat{I}(0) \rangle \\ \frac{1}{2} \left( \langle \hat{I}(0) \hat{I}(t) \rangle + \langle \hat{I}(t) \hat{I}(0) \rangle \right) \end{array} \right.$$

Absorption

Emission

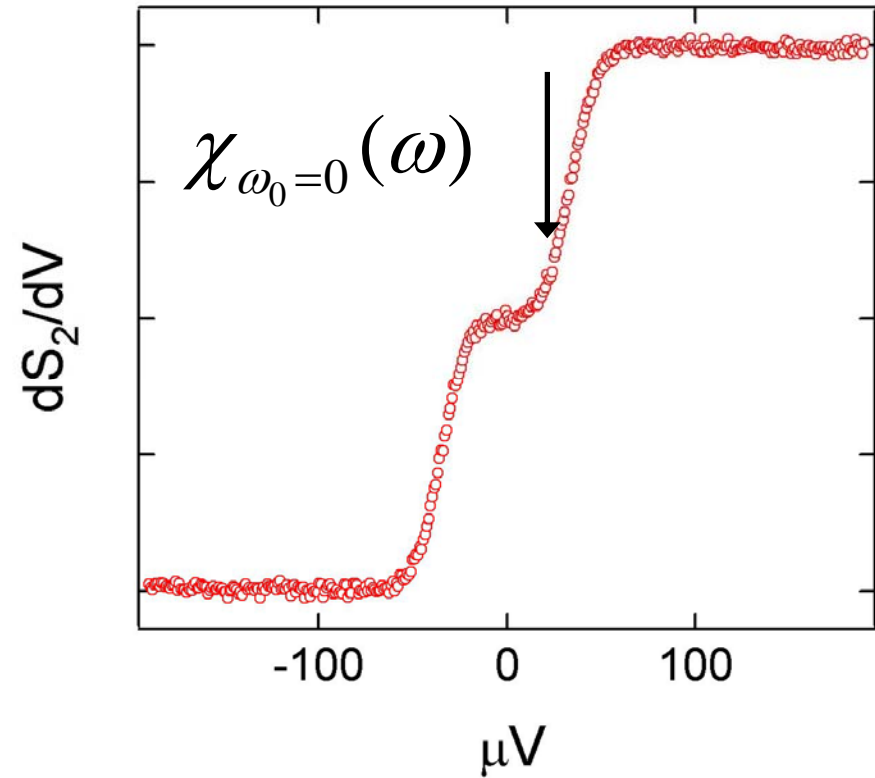
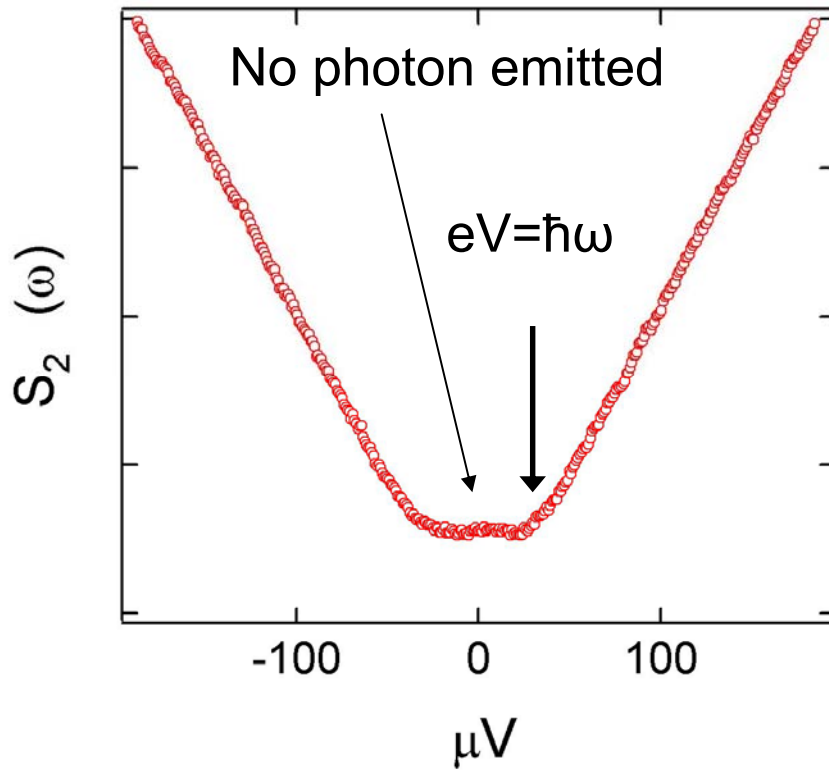
Classical

$$S_2^{abs}(\omega) = S_2^{em}(-\omega)$$

$$S_2^{sym}(\omega) = S_2^{em}(\omega) + \frac{1}{2} G \hbar \omega$$



# Reminder: noise in the quantum regime $\hbar\omega > k_B T, eV$ for a tunnel junction



Measured on a tunnel junction at  $T=35\text{mK}$ ,  $f=6\text{ GHz}$ ,  $\hbar\omega/k_B T \sim 8.5$

# Noise susceptibility – beyond the classical regime: theory

What if  $\omega_0 > \omega$  ?  
What if  $\hbar\omega > k_B T$  ?

$$\chi_{\omega_0}(\omega) \propto \langle \delta I(\omega) \delta I(\omega_0 - \omega) \rangle$$

Calculation:

- \* Landauer-Büttiker formalism
- \* SYMMETRIZATION of the operators and  $\omega \rightarrow -\omega$

The symmetrization rule depends on the experimental setup !

In particular:  $\omega \sim \omega_0$

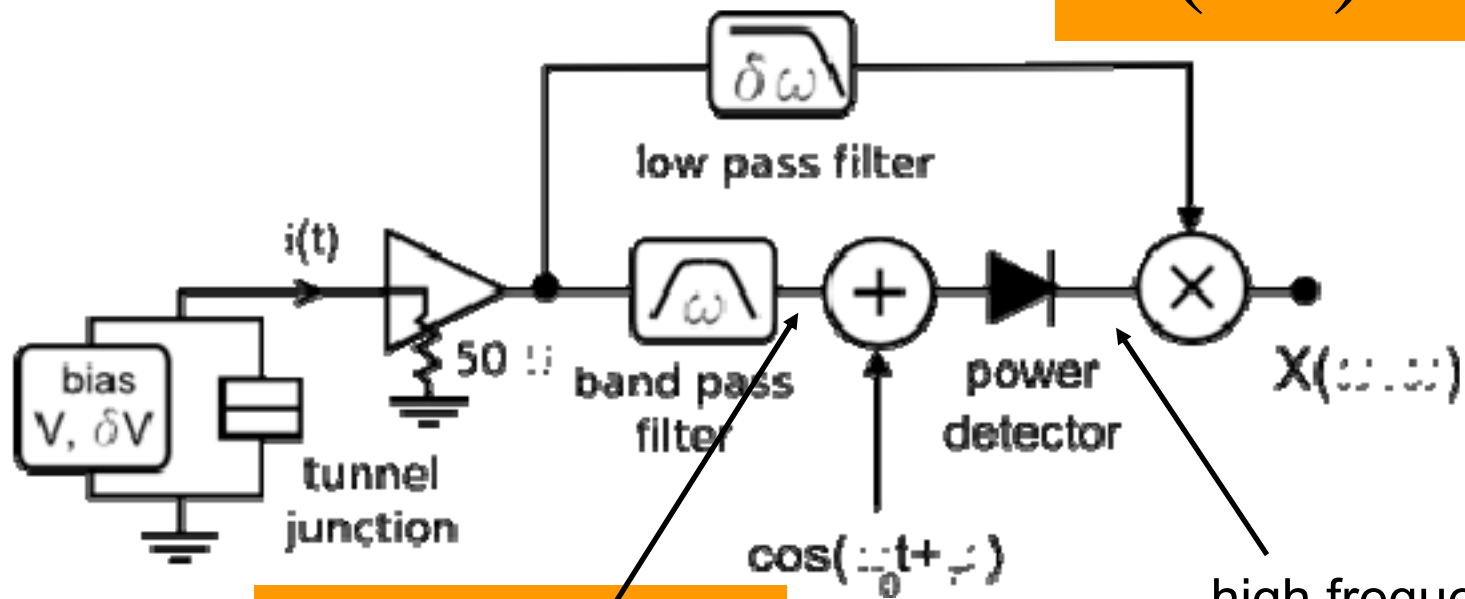
$$\chi_{\omega}(\omega) \propto \langle \delta I(\omega) \delta I(0) \rangle$$

# Noise susceptibility – the quantum regime: experiment

$\omega_0 \sim \omega \sim 6 \text{ GHz}$   
 $\hbar\omega/k_B T \sim 8.5$   
 $\delta\omega \sim 100 \text{ MHz}$

low frequency current

$$\delta I(\pm \varepsilon) e^{\pm i \varepsilon t}$$

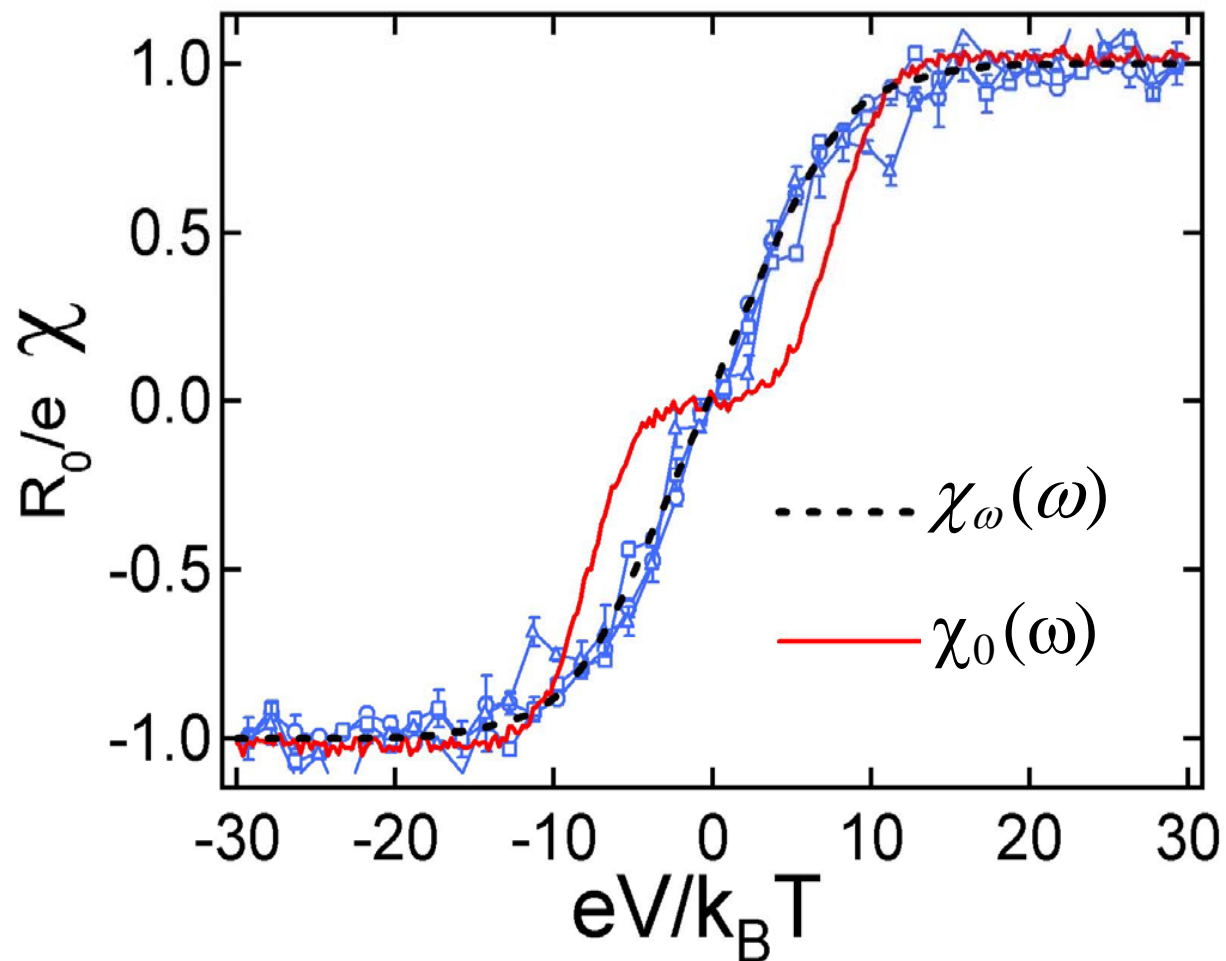


$$\delta I(\pm \omega) e^{\pm i \omega t}$$

high frequency current shifted to low freq.

# Noise susceptibility – the quantum regime: experiment

$\omega_0 \sim \omega \sim 6$  GHz  
 $T = 35$  mK  
 $\hbar\omega/k_B T \sim 8.5$

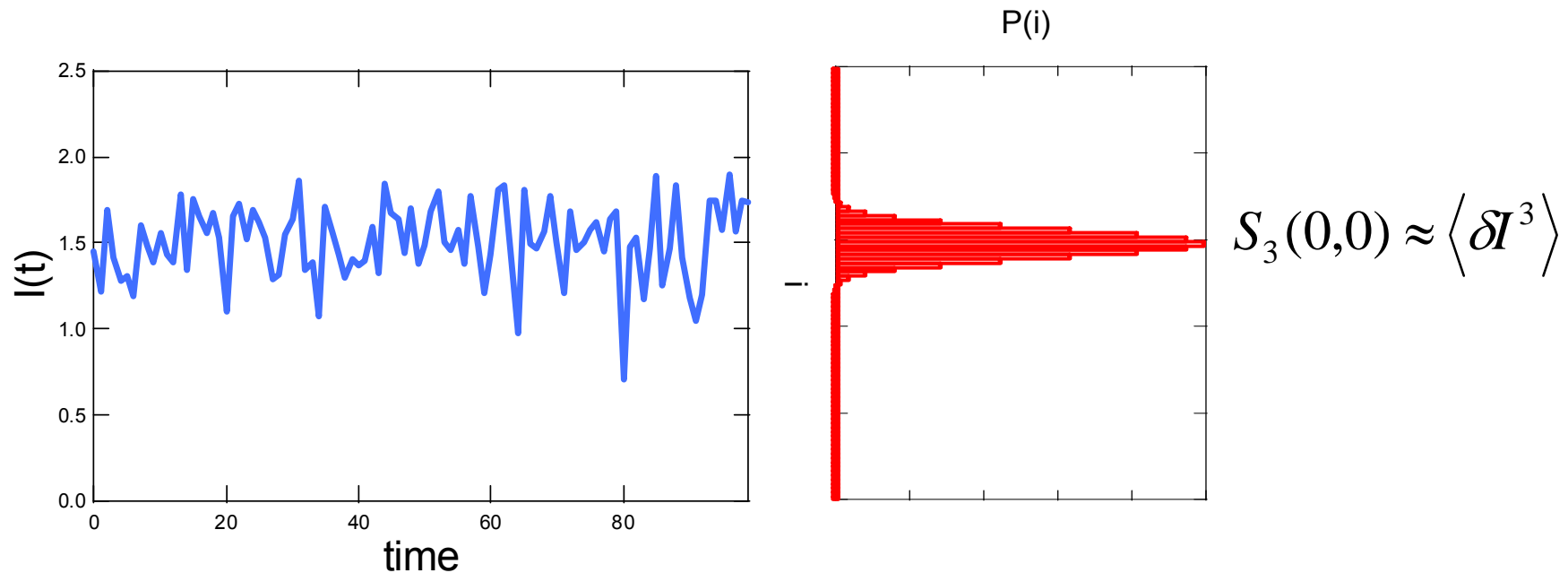


The third cumulant  
of current fluctuations  
in the classical regime

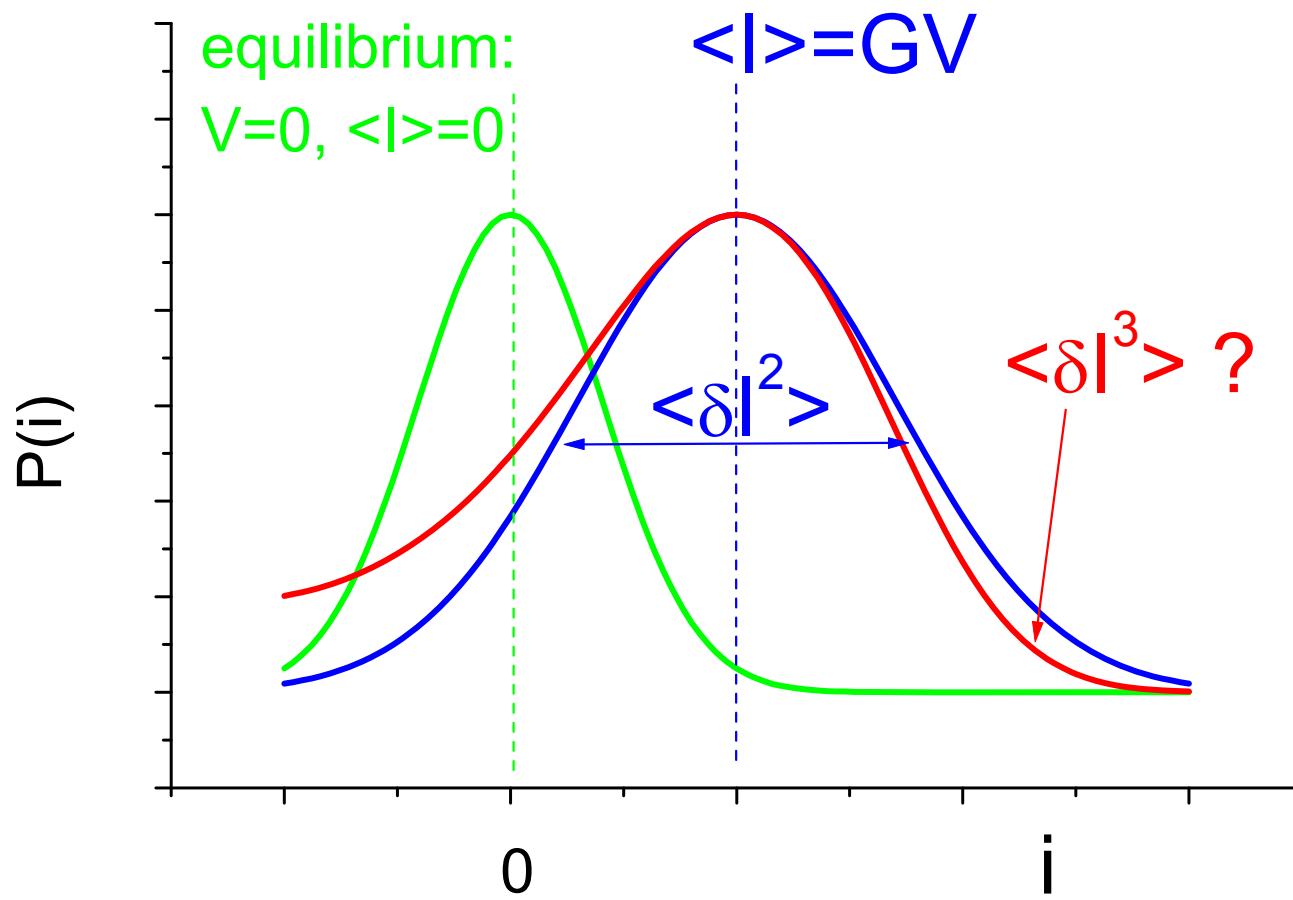
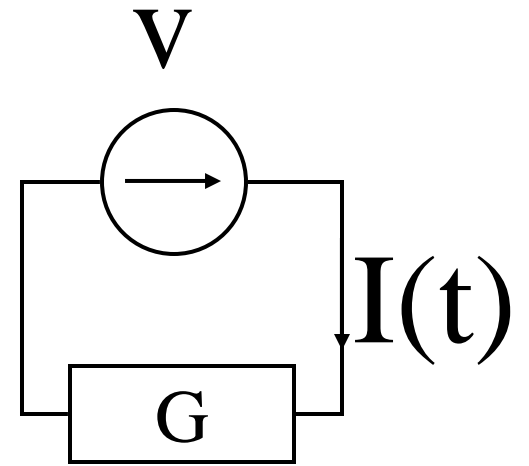
# The third cumulant of noise

$$S_3(\omega, \omega') = \langle \delta I(\omega) \delta I(\omega') \delta I(-\omega - \omega') \rangle$$

At low frequency:  $S_3 =$  **SKEWNESS** of the probability distribution of current fluctuations  $P(I)$ : **zero for gaussian noise**



# Statistics of the current

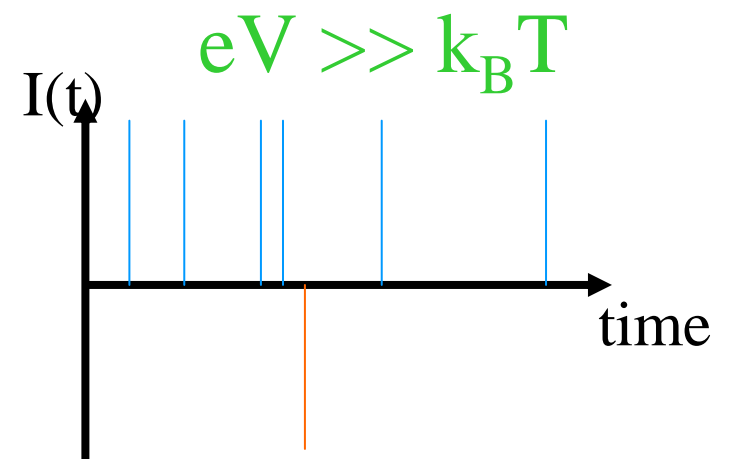
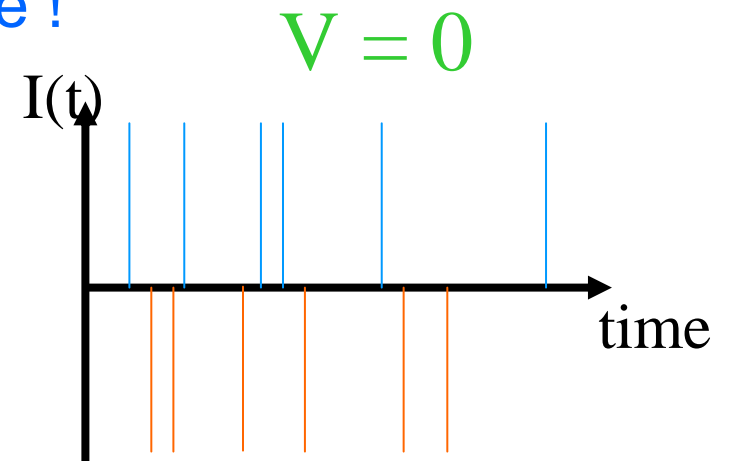
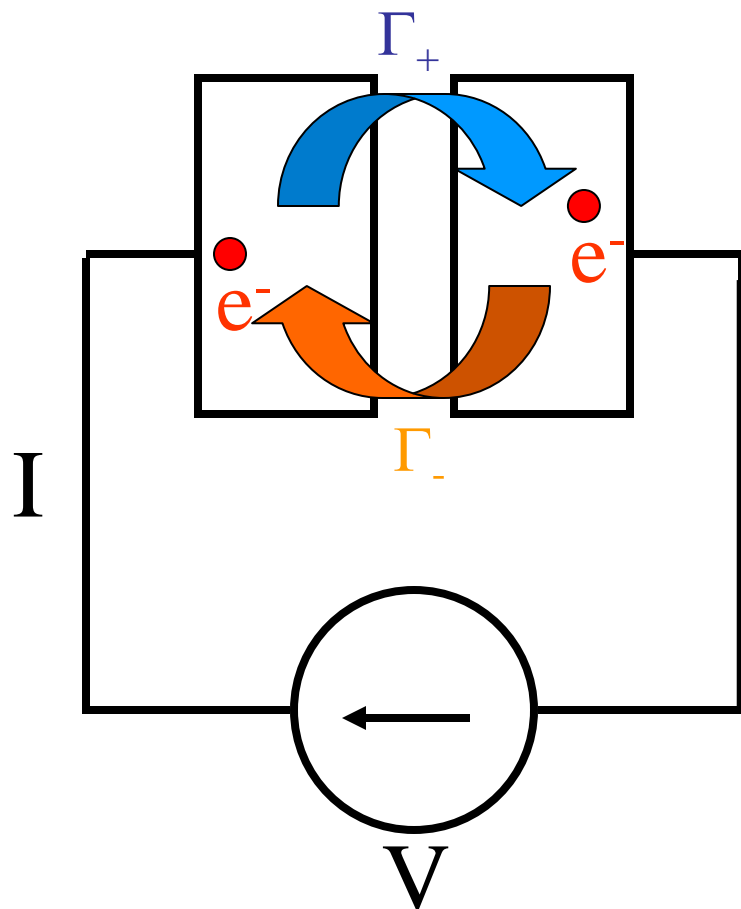


$$I(t) = \langle I \rangle + \delta I(t)$$

At equilibrium: Johnson noise:  $\langle \delta I^2 \rangle = 4k_B TGB$ , (B=bandwidth)

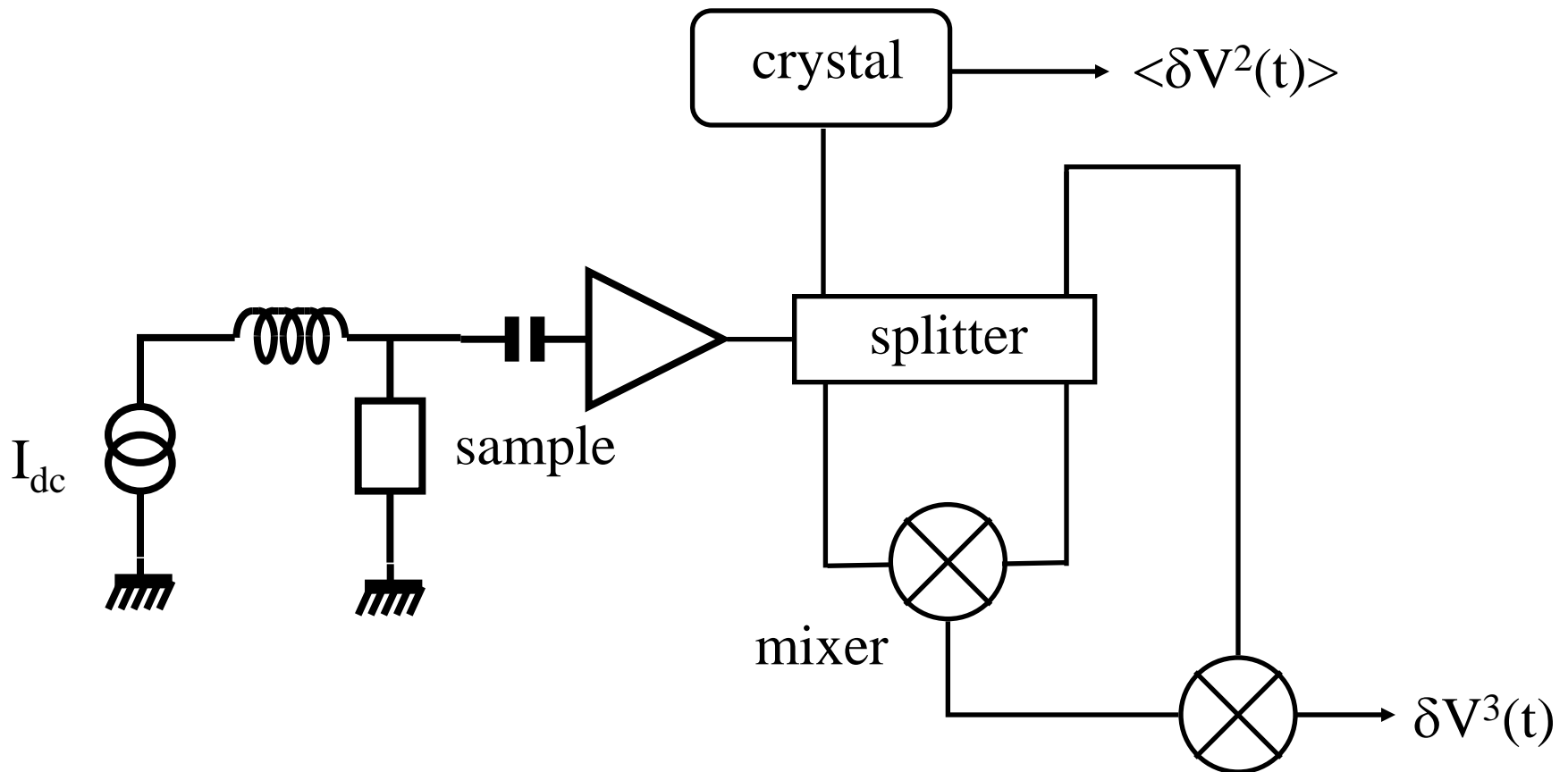
# Current in a tunnel junction

$S_3 = e^2 I$  independent of temperature !



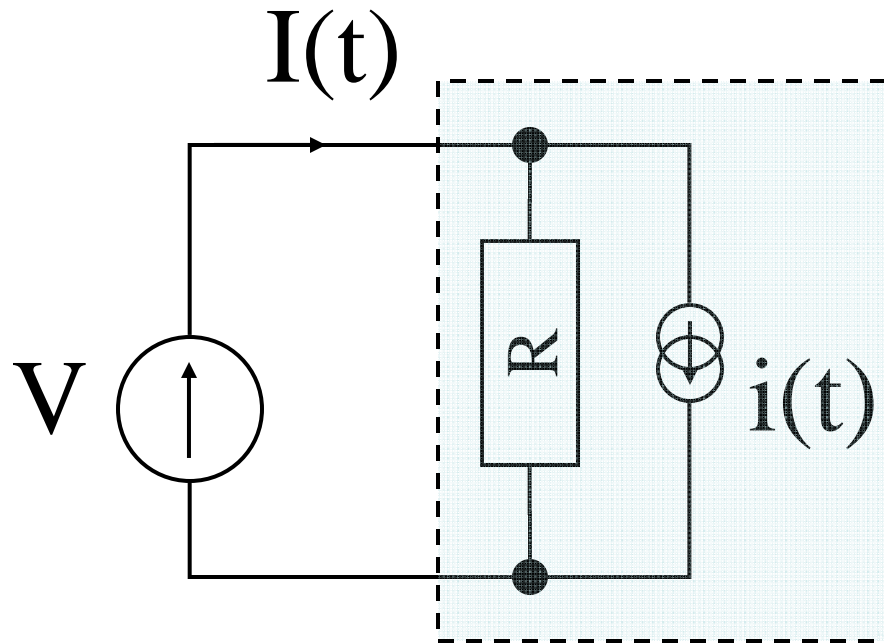


# Measurement: method

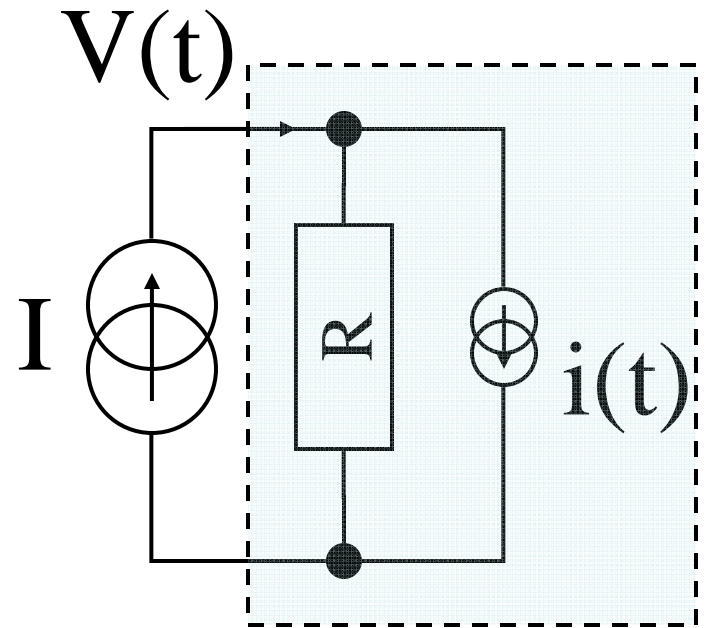


Bandwidth: 10-1200 MHz

# Voltage bias vs. Current bias



$$\delta I(t) = i(t)$$



$$\delta V(t) = -Ri(t)$$

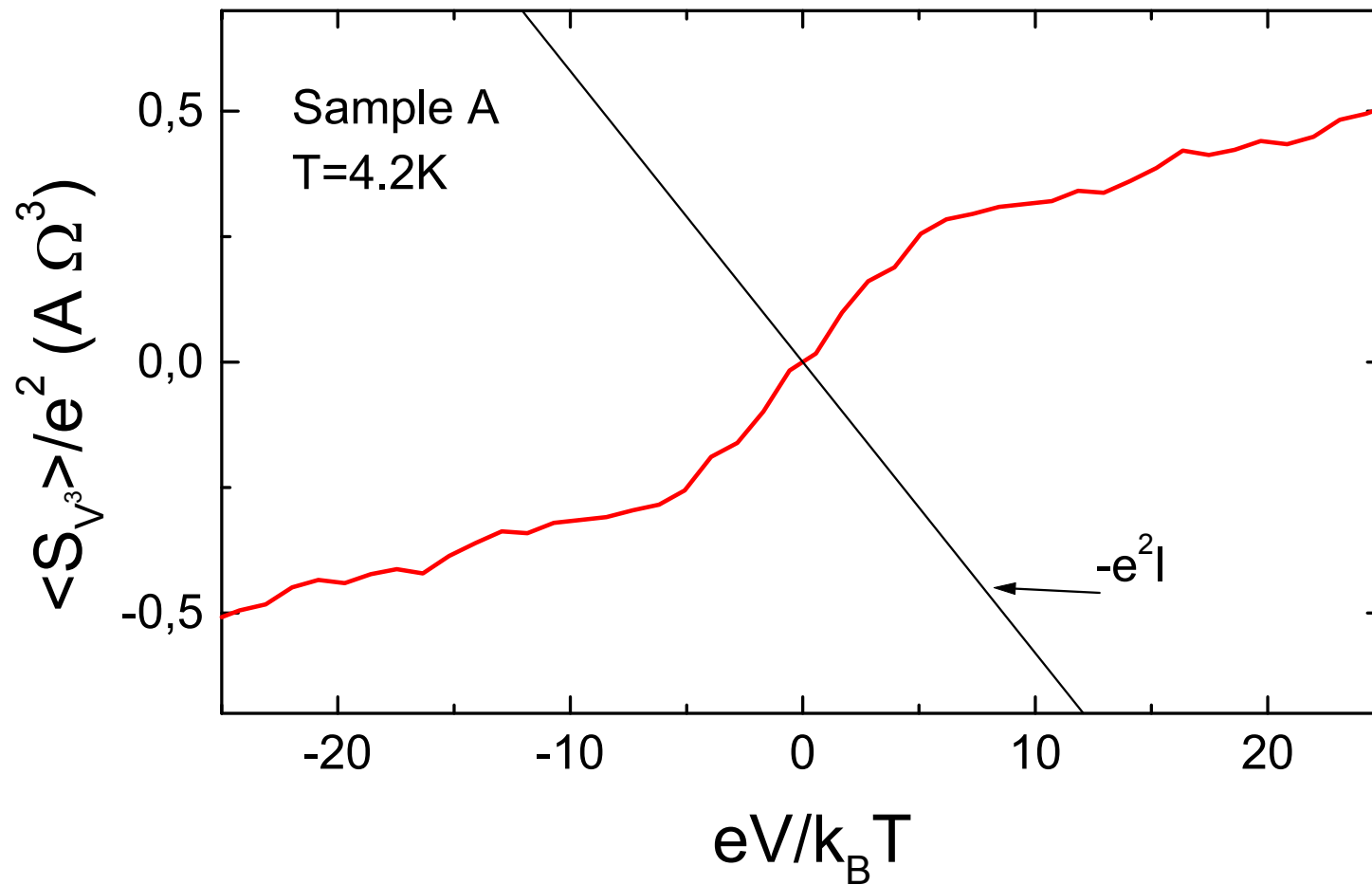
$$V = RI$$

$$\langle \delta V^2 \rangle = R^2 \langle \delta I^2 \rangle$$

$$\langle \delta V^3 \rangle = -R^3 \langle \delta I^3 \rangle ?$$

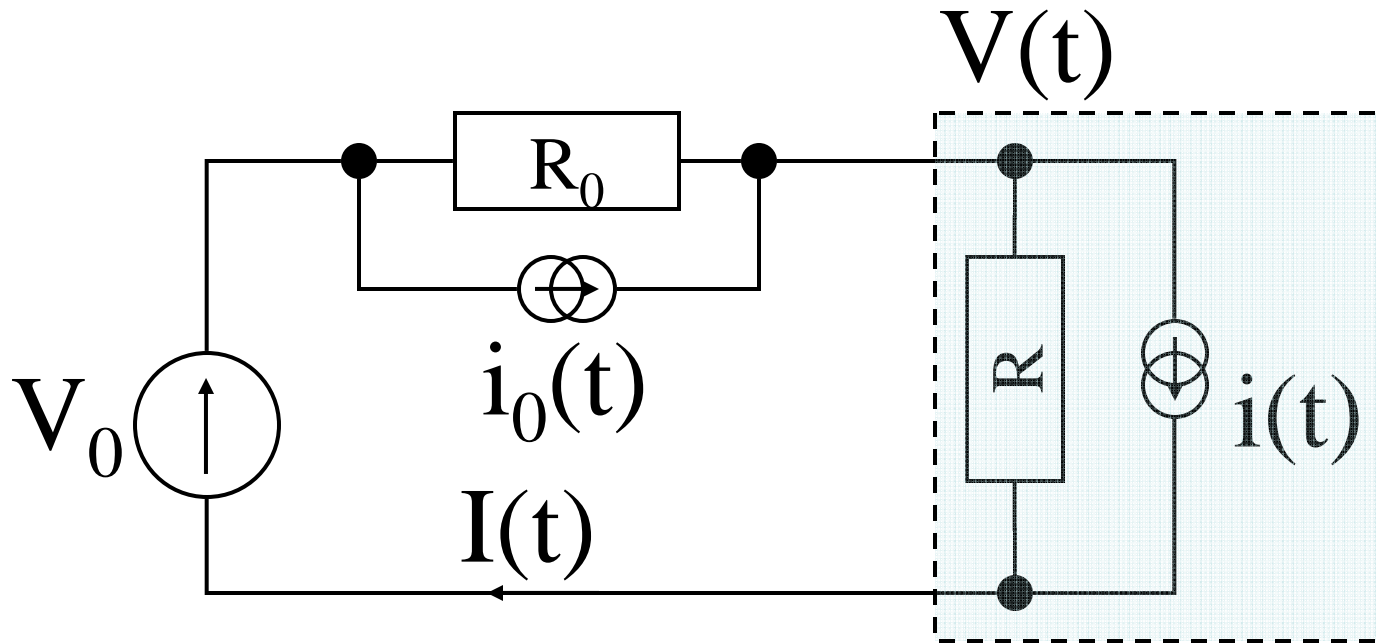
# Result (T=4.2K)

(averaged over 12 days !)



# Environmental effects

# Imperfect voltage bias



$$\delta V(t) = (R // R_0) (i_0(t) - i(t))$$

External noise    Feedback of the environment

The probability distribution  $P(i)$  depends on  $V(t)$

# Effect of the environment on the probability distribution $P(i)$

$$P(i, V) = P_V(i, V - R_0 i) \cong P_V(i, V) - R_0 i \frac{\partial P_V(i, V)}{\partial V}$$

Probability when voltage biased

Effect on the moments:

$$\langle i^n \rangle = \langle i^n \rangle_V - R_0 \frac{\partial \langle i^{n+1} \rangle_V}{\partial V}$$

Origin: e-e interactions

# Application: n=3

\* Effect on the third moment:

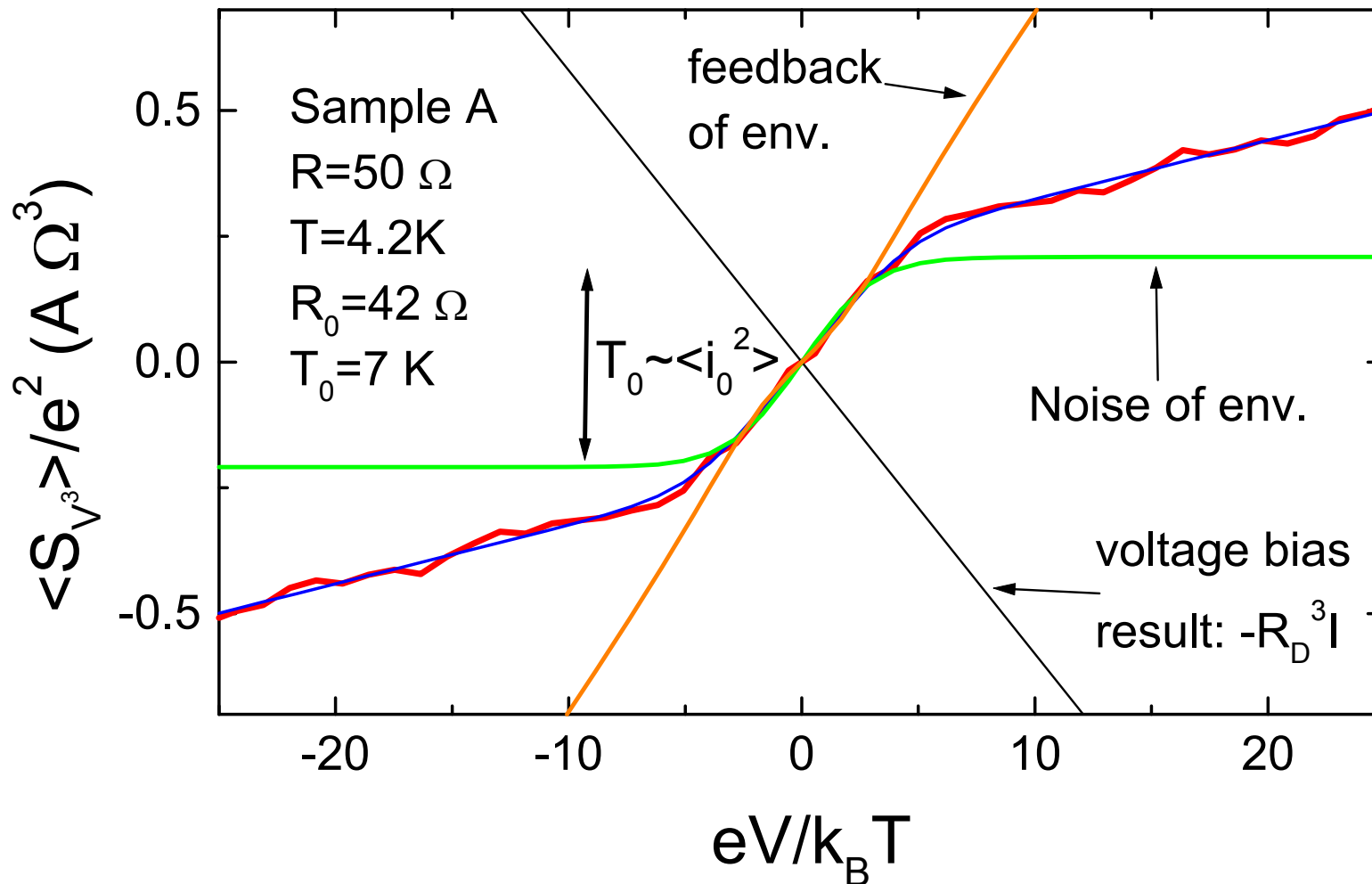
$$\langle i^3 \rangle = e^2 I - R_0 \frac{\partial S_4}{\partial V} \approx e^2 I - 6R_0 S_2 \frac{\partial S_2}{\partial V}$$



Noise Susceptibility  
at  $\omega_0 \sim 0$

$$\chi_{\omega_0=0}(\omega) = \frac{\partial S_2(\omega)}{\partial V}$$

# Exp. Result + theory





# Effect of the environment on the probability distribution $P(i)$

$$\langle i^n \rangle = \langle i^n \rangle_V - R_0 \frac{\partial \langle i^{n+1} \rangle_V}{\partial V}$$

## Application: $n=1$

\* Effect on the dc current:  
dynamical Coulomb blockade

$$\langle i \rangle = I_{dc} - R_0 \frac{\partial S_2}{\partial V}$$

Noise Susceptibility  
at  $\omega_0 \sim \omega$

$$\chi_\omega(\omega)$$

# Link between Dynamical Coulomb Blockade and Noise Susceptibility

$$\delta I = \frac{\hbar}{e^2} \int K(-\omega) \omega \chi_\omega(\omega) d\omega$$

Correlation function of  
the environment

Noise susceptibility for  $\omega_0 = \omega$

cf Lamb shift...

For small environmental impedance:

$$\delta I = - \int \text{Re} Z_0(\omega) \chi_\omega(\omega) d\omega$$

There is nothing quantum ...

The third cumulant  
of current fluctuations  
in the quantum regime

# $S_3$ and Q mechanics: ordering ???

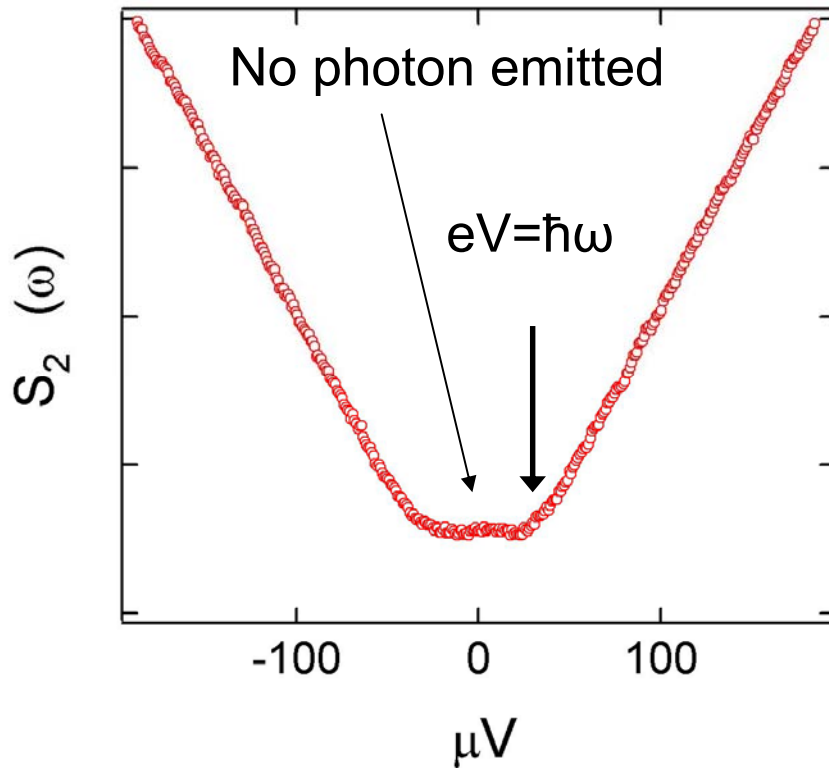
$$S_3(\omega, \omega') = \int dt dt' e^{i(\omega t + \omega' t')} \langle \hat{I}(0, t, t') \hat{I}(0, t, t') \hat{I}(0, t, t') \rangle$$

Two results:

$$S_3(0,0) = \frac{e^2}{h} V \cdot \begin{cases} t(1-t)(1-2t) & \text{Keldysh ordering} \\ t^2(1-t) & \text{Fully symmetrized} \end{cases}$$

$t \approx 10^{-5}$  for a tunnel junction

# The third cumulant in the quantum regime: $S_3(0, \omega)$ with $\hbar\omega > k_B T, eV$

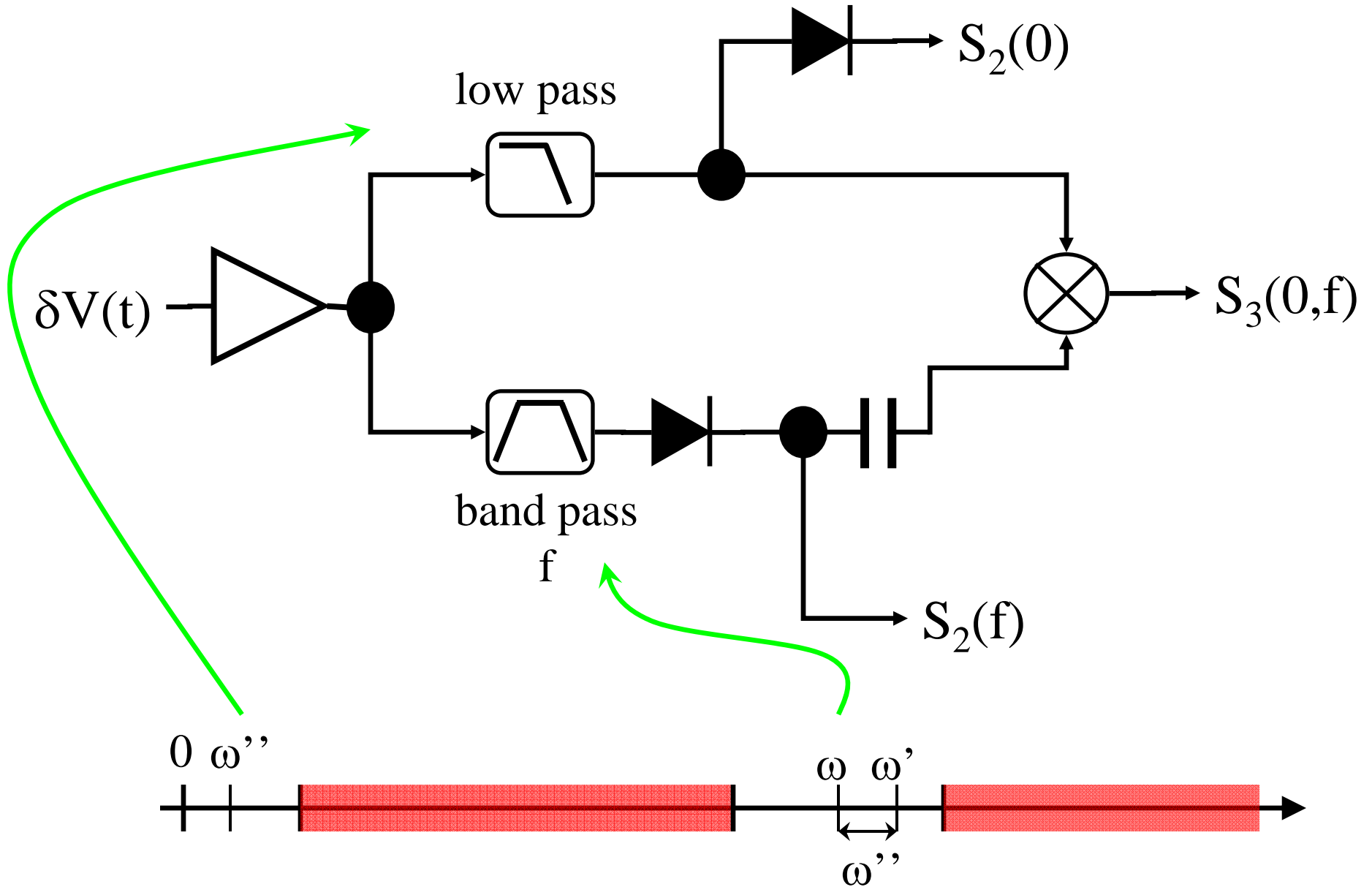


Prediction:

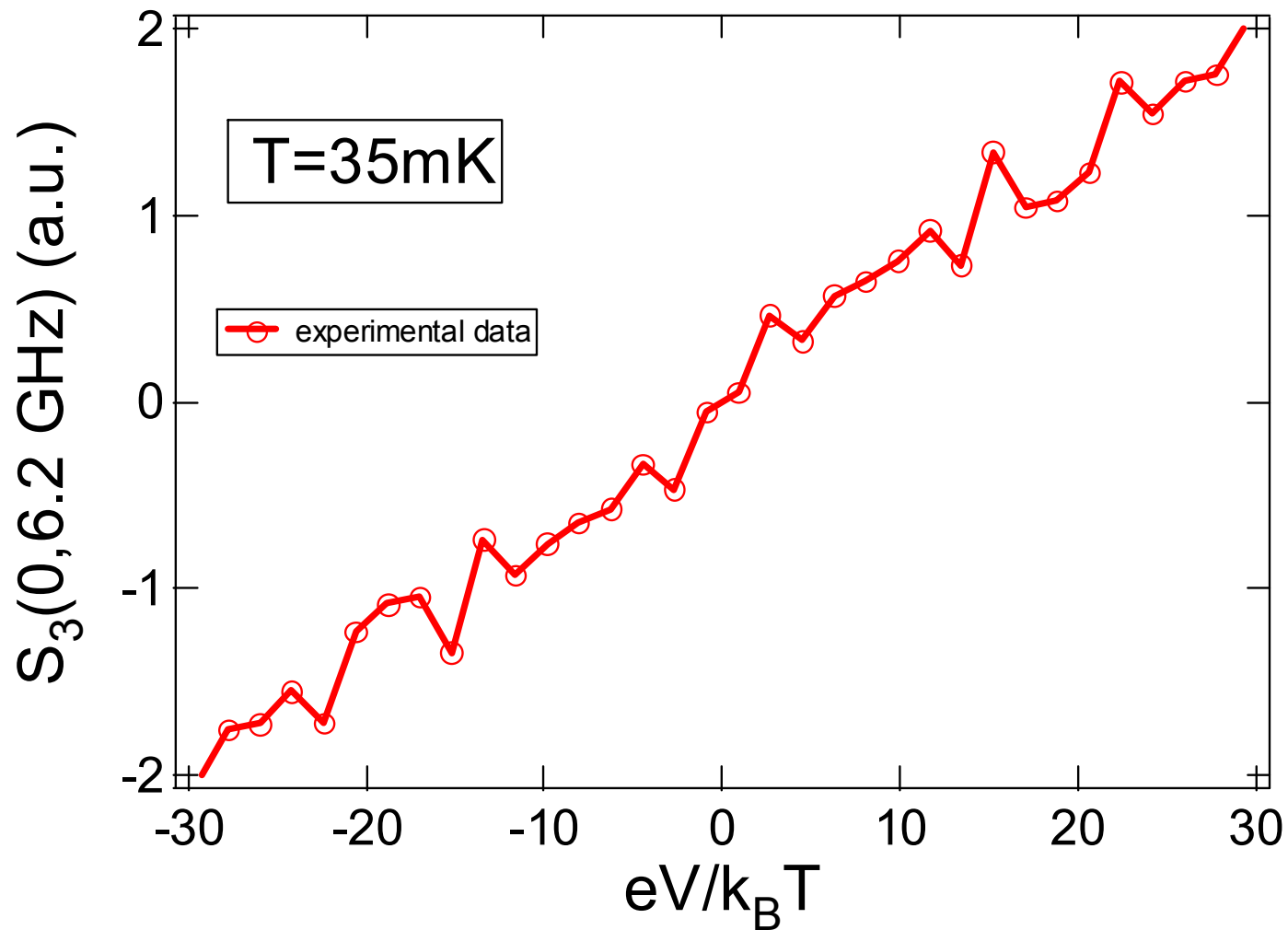
$$S_3(\omega, \omega') = e^2 I$$

indep of frequency !!

# How to measure $S_3(0, f)$ ?



# $S_3(0, \omega)$ : experiment - minus environment in the Q regime



# The third cumulant in the quantum regime: $S_3(0, \omega)$ with $\hbar\omega > k_B T, eV$

$$S_2(\omega) = \langle \delta I(\omega) \delta I(-\omega) \rangle \quad \sim \text{power of light emitted at frequency } \omega$$

=0 for  $\hbar\omega > eV$

$$S_3(0, \omega) = \langle \delta I(0) \delta I(\omega) \delta I(-\omega) \rangle \quad \sim \text{correlations between low frequency current fluctuations and power of light at frequency } \omega.$$

Prediction:

$$S_3(\omega, \omega') = e^2 I \text{ independent of frequency !?}$$

**WARNING:** zero point motion of electrons !!



# Another way to measure $S_3(0,f)$ ?

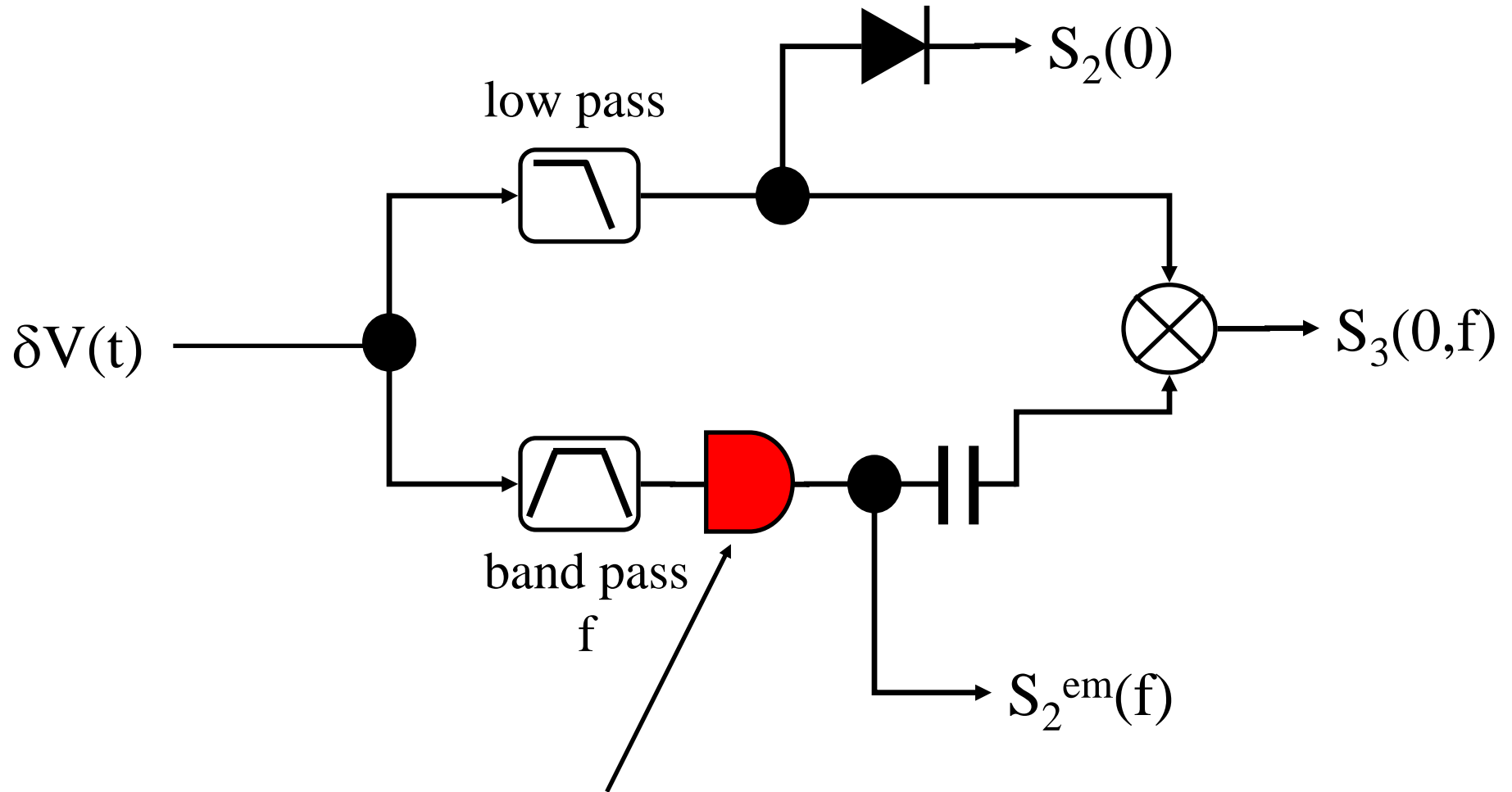
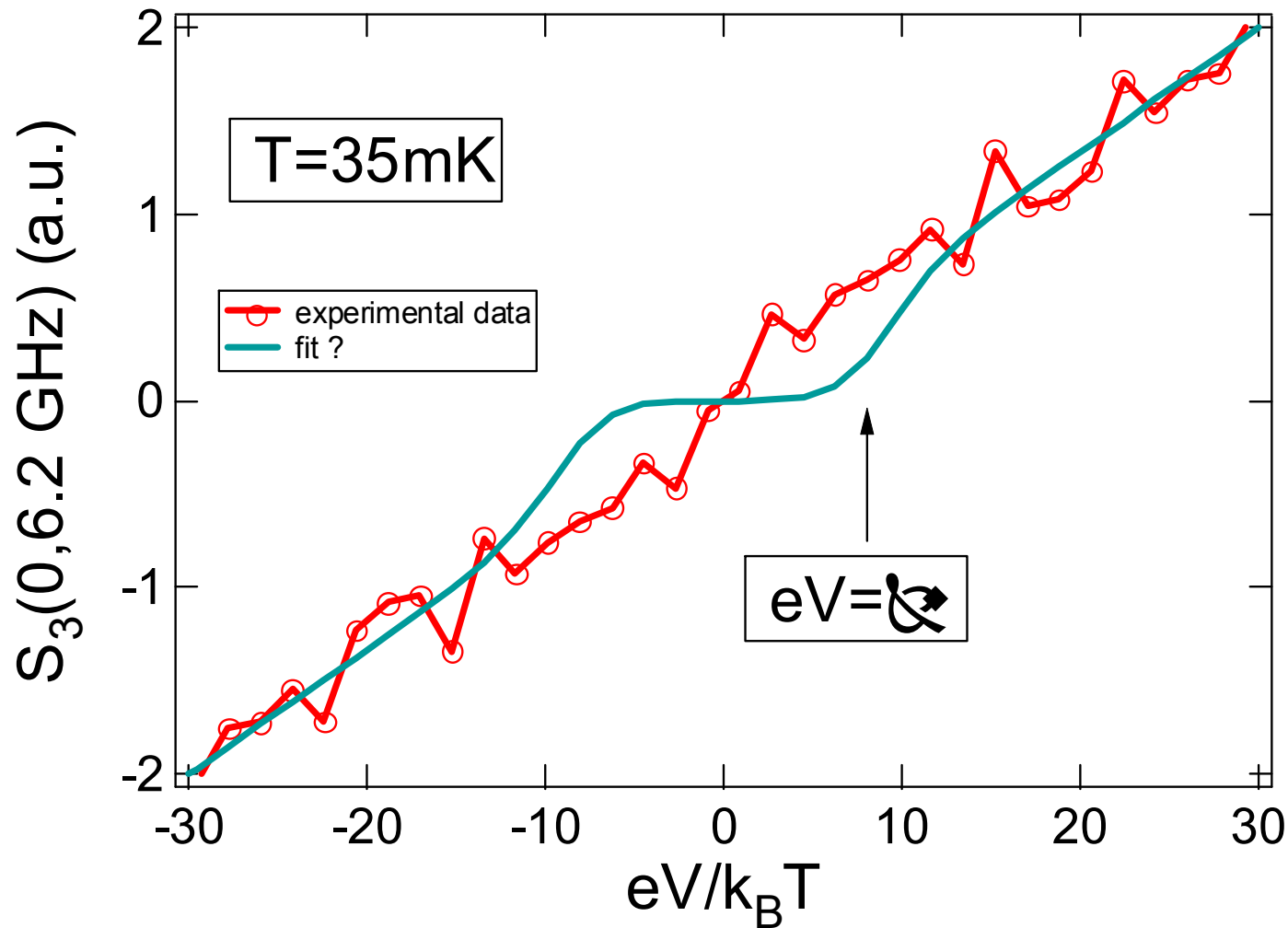


Photo-multiplier: absorbs photons

# $S_3(0, \omega)$ : experiment - minus environment in the Q regime



# Another way to measure $S_3(0,f)$ ?

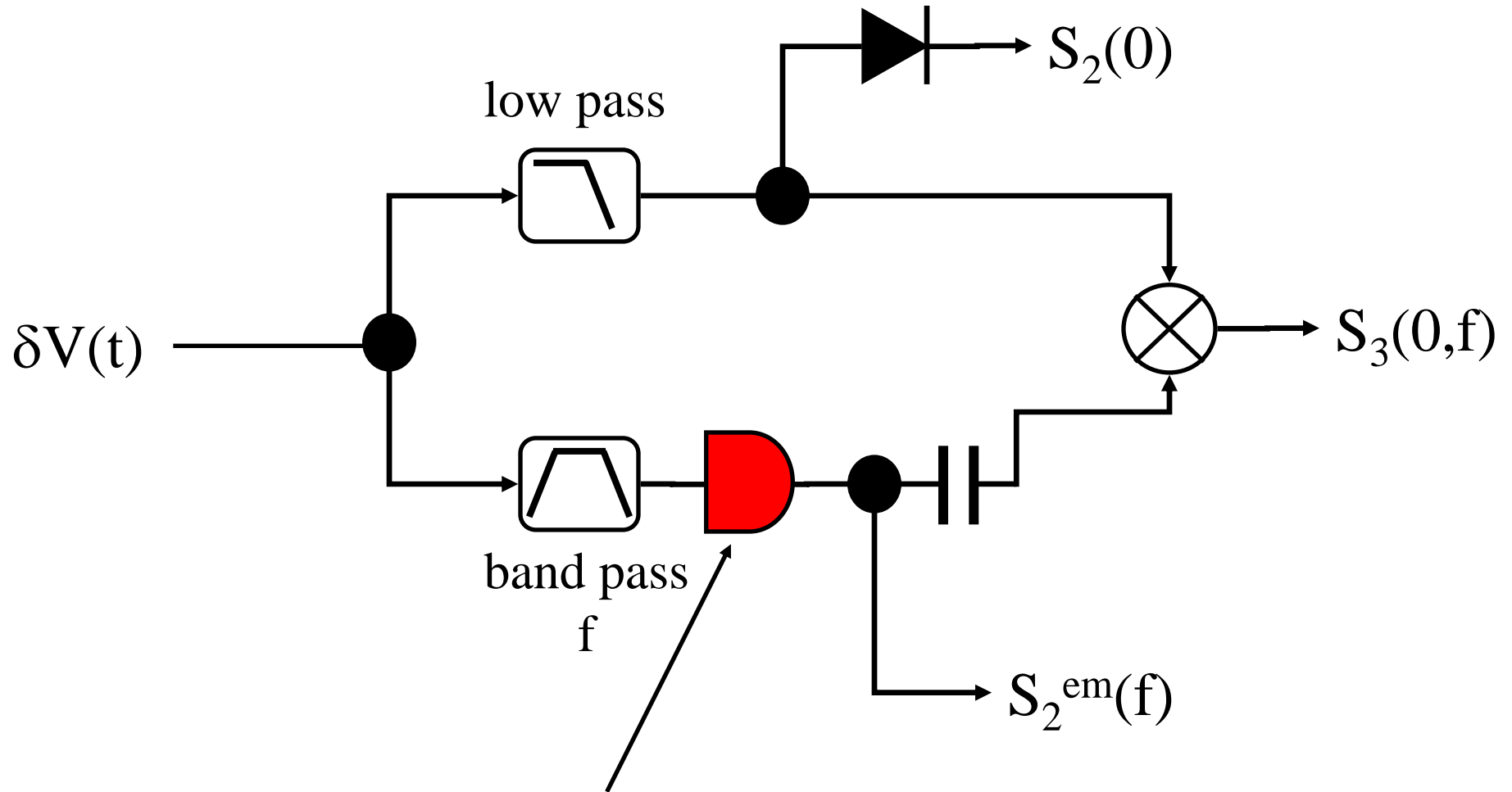
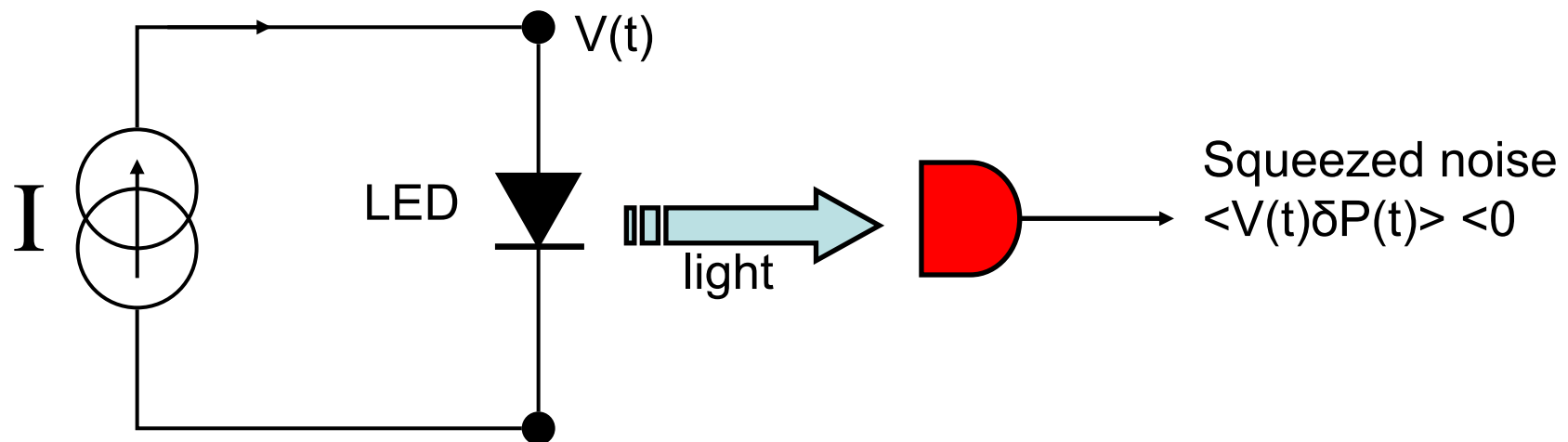
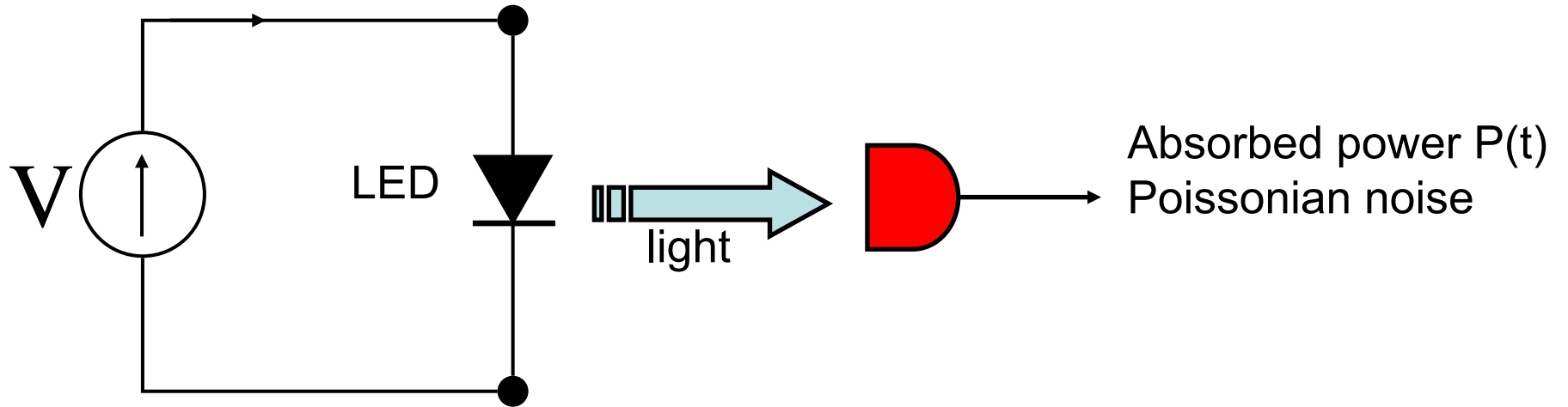


Photo-multiplier: absorbs photons

# Similar Q optics experiments...



# Conclusion

## NOISE SUSCEPTIBILITY

Probe for inelastic mechanisms:

- \* electron-phonon time vs  $T$  in normal metals at very low  $T$
- \* electron-electron time
- \* in more subtle systems (carbon nanotubes, quasi-crystals, etc.)

Link with dynamical Coulomb blockade:

- \* Van der Waals interaction between noise sources ?
- \* Coulomb blockade at room  $T$  ?

## THIRD CUMULANT OF NOISE

Test for deep understanding of transport mechanisms ?

Measurement theory: how to order three current operators in quantum mechanics ?

# Which system exhibits gaussian noise ?

- \* Not a radiator, not a light bulb run through by a current (more current, more heat or light)...
- \* Not the cosmologic background (?)
- \* Not a chemical reaction (cf tunnel junction)
- \* Not a finite system, even at equilibrium (there is  $S_4$ ...)

ONLY a perfectly transmitting channel...