



COLLÈGE
DE FRANCE
—1530—



Chaire de Physique Mésoscopique
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Année 2007, Cours des 7 et 14 juin

INTRODUCTION À LA PHYSIQUE MÉSOSCOPIQUE: ÉLECTRONS ET PHOTONS

INTRODUCTION TO MESOSCOPIC PHYSICS: ELECTRONS AND PHOTONS

Première Leçon / First Lecture

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AIM OF THESE TWO LECTURES

Discuss a selection of basic concepts of mesoscopic physics and contrast their treatment with that of atomic physics

OUTLINE

1. General remarks on mesoscopic physics
 - a) fundamental constants
 - b) survival of quantum effects in macrosystems
2. What are "electrons"?
 - a) example of mesoscopic resistor
 - b) screening
 - c) finite lifetime
3. What are "photons"?
 - a) example of transmission line
 - b) longitudinal and transverse modes
 - c) 1-D Planck's law
4. Conclusion: quantum transport v.s. quantum optics



part 1

part 2

General remarks

Purpose: explain the novelty of mesoscopic phenomena by discussing status of variables and parameters

MACRO versus micro

MACRO SYSTEM

classical
mechanics



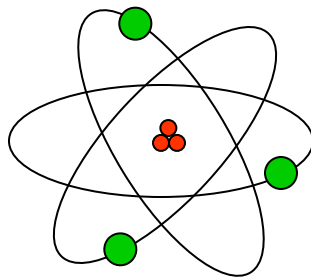
many
particles,
strongly
coupled to
environment

$$H = \frac{p_\theta^2}{2M} + \frac{k\theta^2}{2} + \dots$$

can have
engineered
structural
parameters

micro system

quantum
mechanics



few
particles,
almost
isolated
from
environment

$$\hat{H} = \frac{\hat{p}^2}{2m_e} - \frac{e^2}{4\pi\epsilon_0\hat{r}} + \dots$$

always
"God-given"
structural
parameters

MESOSCOPIC PHENOMENA:

Number of particles: macroscopic
Collective degrees of freedom: quantum

High-Q nanomechanical resonators @ low temperatures:
clearly mesoscopic

Superconducting magnets for MRI: clearly not mesoscopic

Characteristic energies in mesoscopic
phenomena involve not only
fundamental constants, but also
geometric dimensions adjusted by
design and fabrication

REVIEW OF UNIVERSAL CONSTANTS

S. I. Units : s, m, kg, A, K

Classical constants

speed of light c
impedance of vacuum $Z_{vac} \approx 377\Omega$

electrical permittivity
of vacuum $\epsilon_0 = \frac{1}{cZ_{vac}}$

magnetic permeability
of vacuum $\mu_0 = \frac{Z_{vac}}{c}$

Boltzmann's constant k_B

gravitational constant G

Quantum constants

charge quantum e
action quantum \hbar

flux quantum $\Phi_0 = \frac{h}{2e}$

resistance quantum $R_K = \frac{h}{e^2} \approx 26k\Omega$

fine structure cst. $\alpha = \frac{Z_{vac}}{2R_K} \approx \frac{1}{137}$

UNIVERSAL AND MICROSCOPIC QUANTUM CONSTANTS

Universal constants

charge of electron: e

Planck's constant: \hbar

flux quantum: $\Phi_0 = \frac{h}{2e}$

resistance quantum: $R_K = \frac{h}{e^2} \approx 26\text{k}\Omega$

fine structure cst.: $\alpha = \frac{Z_{vac}}{2R_K} \approx \frac{1}{137}$

Microscopic constants

mass of electron: m_e

magn. mt. of electron: μ_B

mass of proton: m_p

mass of neutron: m_n

magn. mt. of proton: μ_p

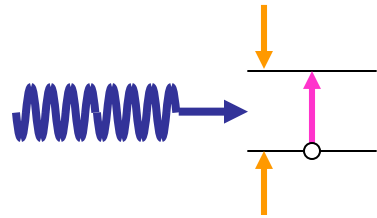
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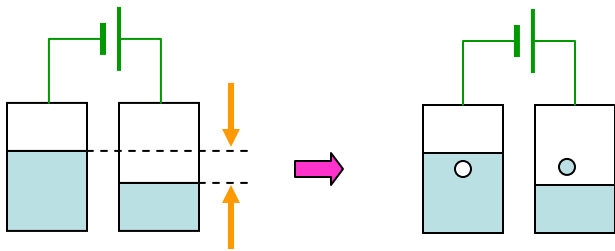
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FORMS OF CHARACTERISTIC ENERGIES



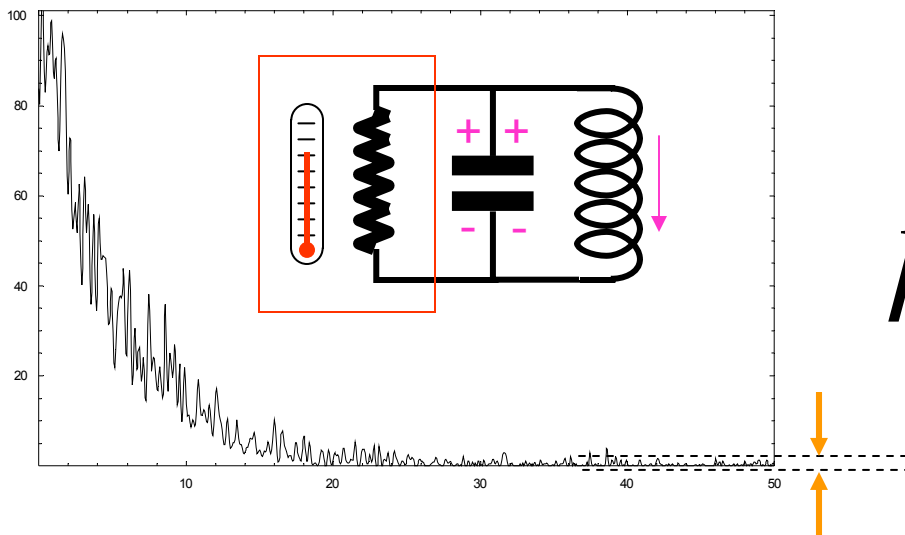
$$\hbar\omega$$

20 GHz



$$eV$$

80 μ V

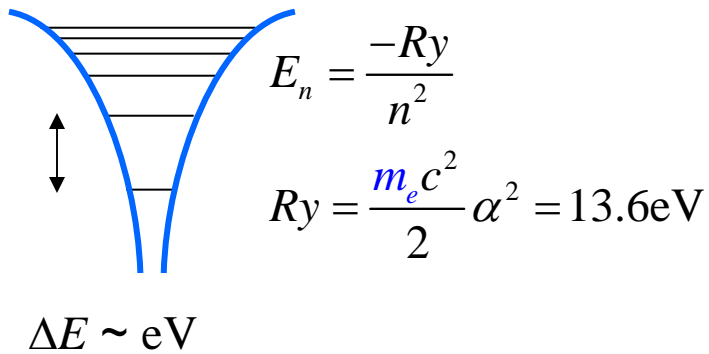


$$k_B T$$

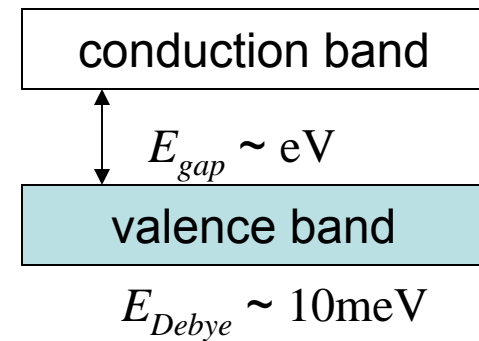
1K

USUAL ENERGY SCALES OF MICROSCOPIC QUANTUM EFFECTS

Atomic matter



Bulk condensed matter



$$\Delta_{sc} = E_{Debye} e^{-\frac{1}{N(0)V}} \sim \text{meV}$$

**IN MESOSCOPIC QUANTUM EFFECTS,
NEW ENERGY SCALES EMERGE!**

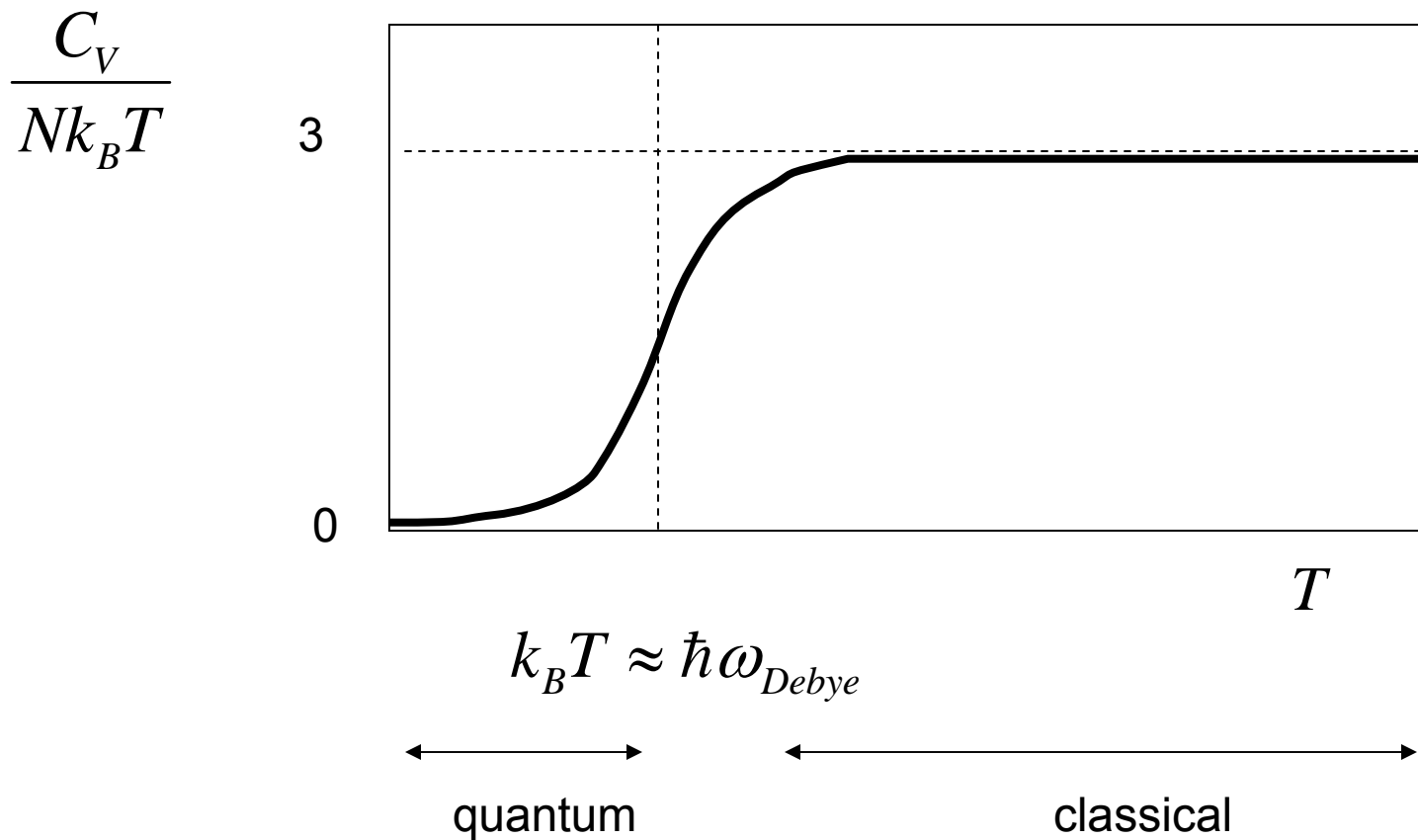
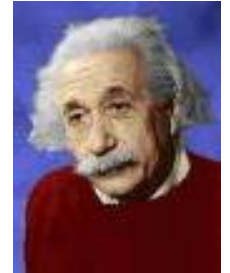
MESOSCOPIC PHYSICS CHALLENGES COMMON WISDOM ABOUT THE SUPPRESSION OF COLLECTIVE QUANTUM EFFECTS IN LARGE AND DIRTY SYSTEMS

- THERMAL ENERGY CAN BE BIGGER THAN SINGLE PARTICLE ENERGY LEVEL SPACINGS
- DISORDER DOES NOT FULLY SUPPRESS INTERFERENCES, IT IS DECOHERENCE WHICH IS IMPORTANT
- A SINGLE ELECTRON CAN COUNT
- LOCAL CIRCUIT THEORY WILL UNEXPECTEDLY FAIL

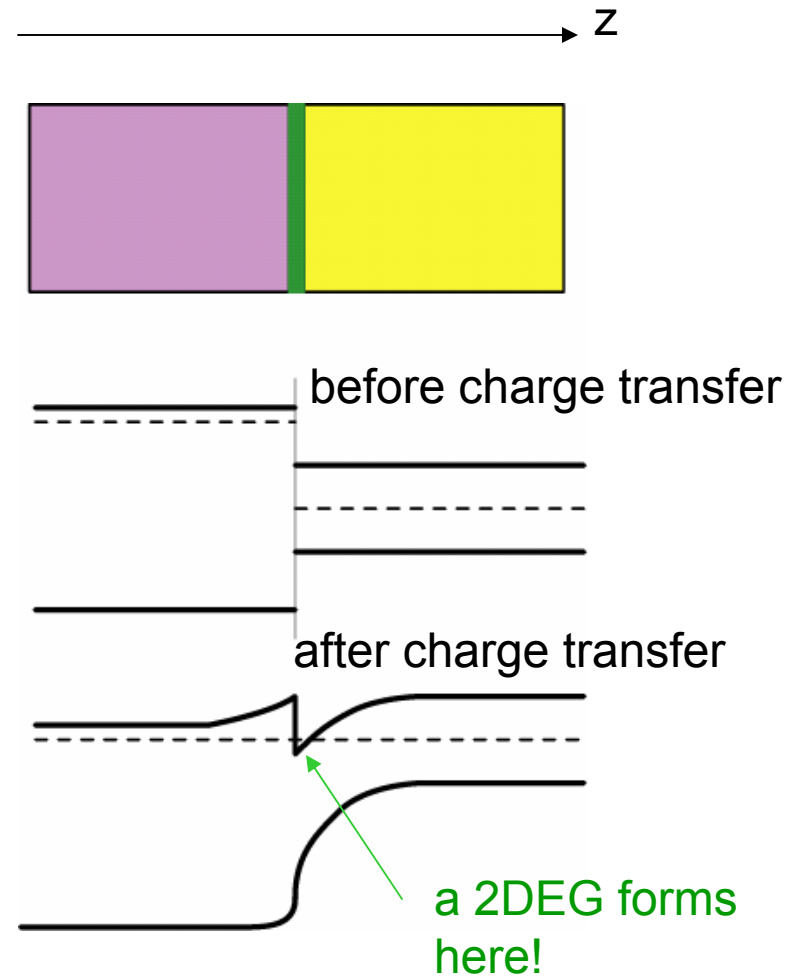
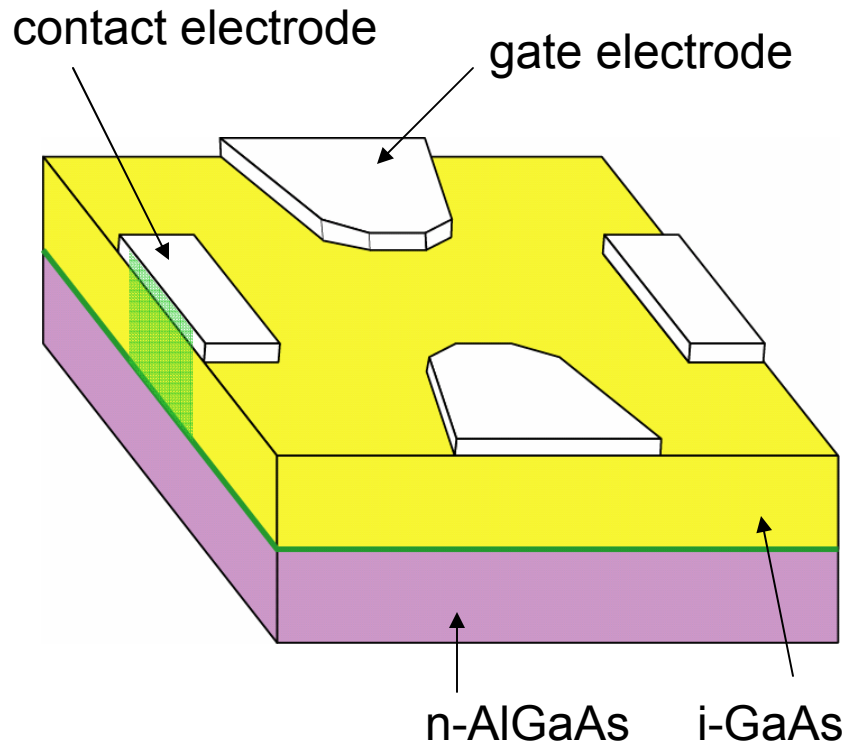
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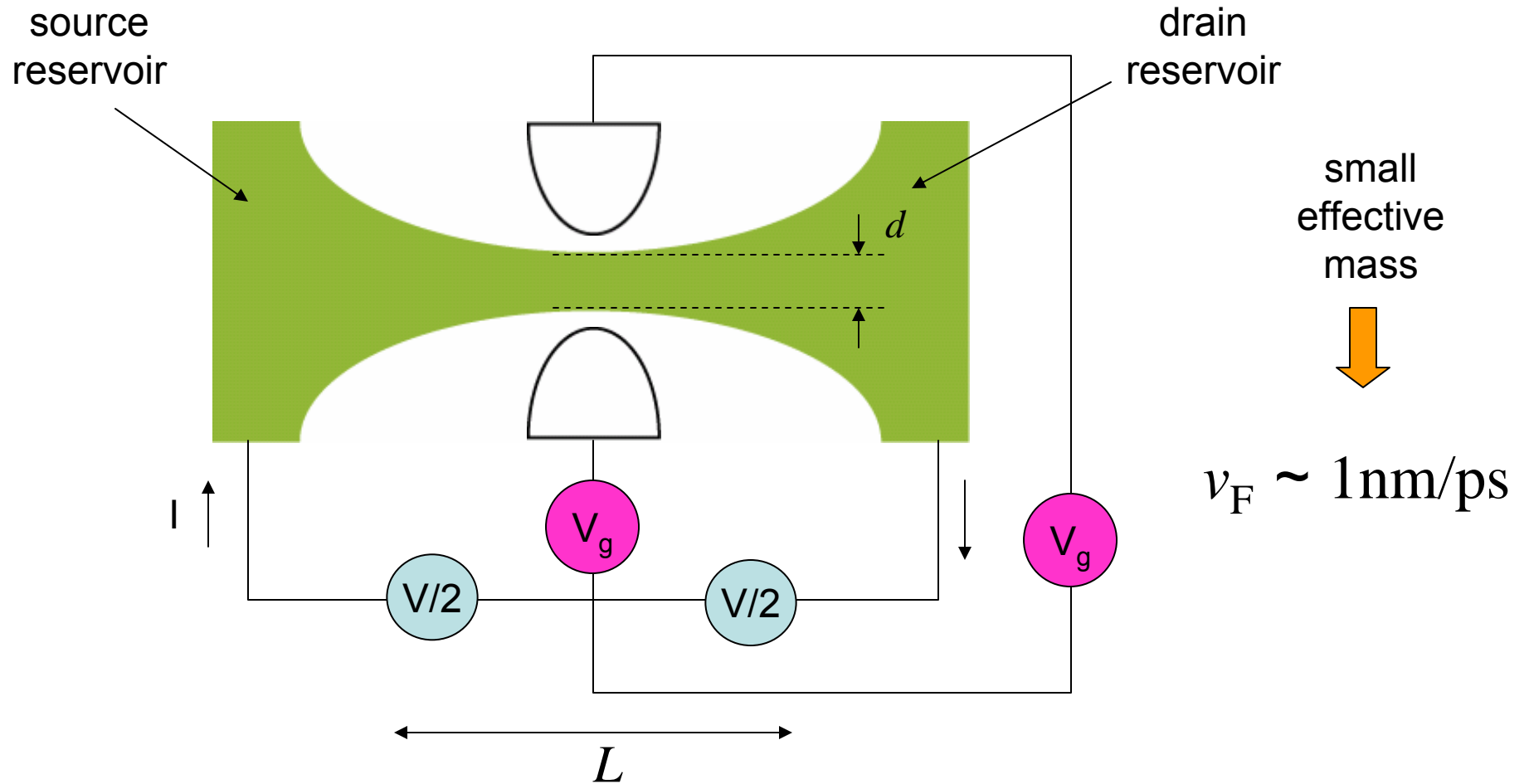
EXAMPLE OF COMMON WISDOM AT WORK: THE CROSSOVER TO DULONG-PETIT LAW



2-DIMENSIONAL ELECTRON GAS SYSTEM



PINCHING THE 2DEG INTO A SMALL WIRE

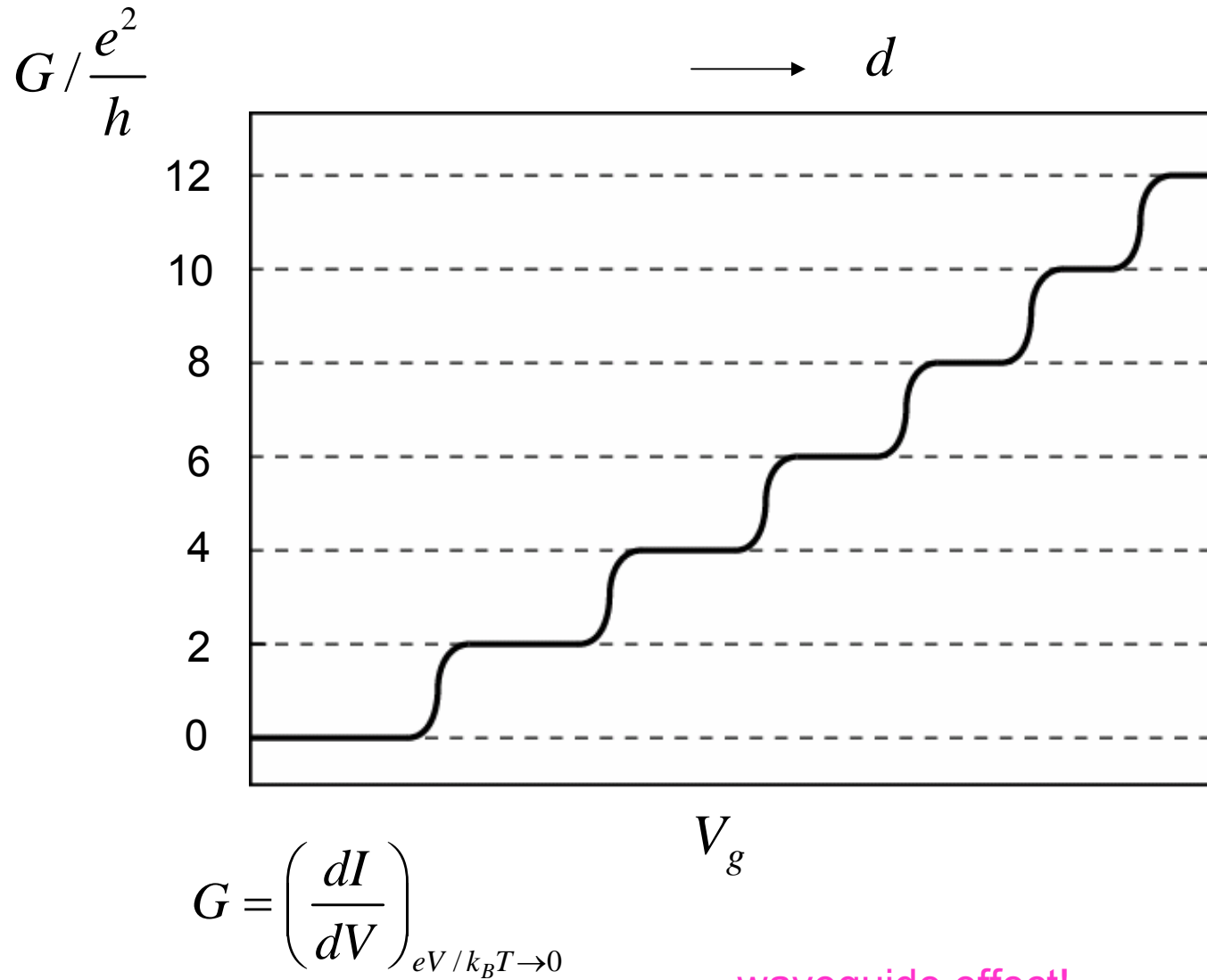


Ballistic transport:

$$L \ll \ell_e$$

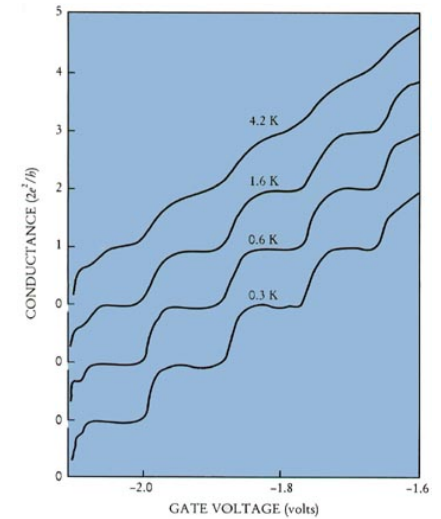
what if $k_B T \ll \frac{\hbar v_F}{d}$?

CONDUCTANCE QUANTIZATION



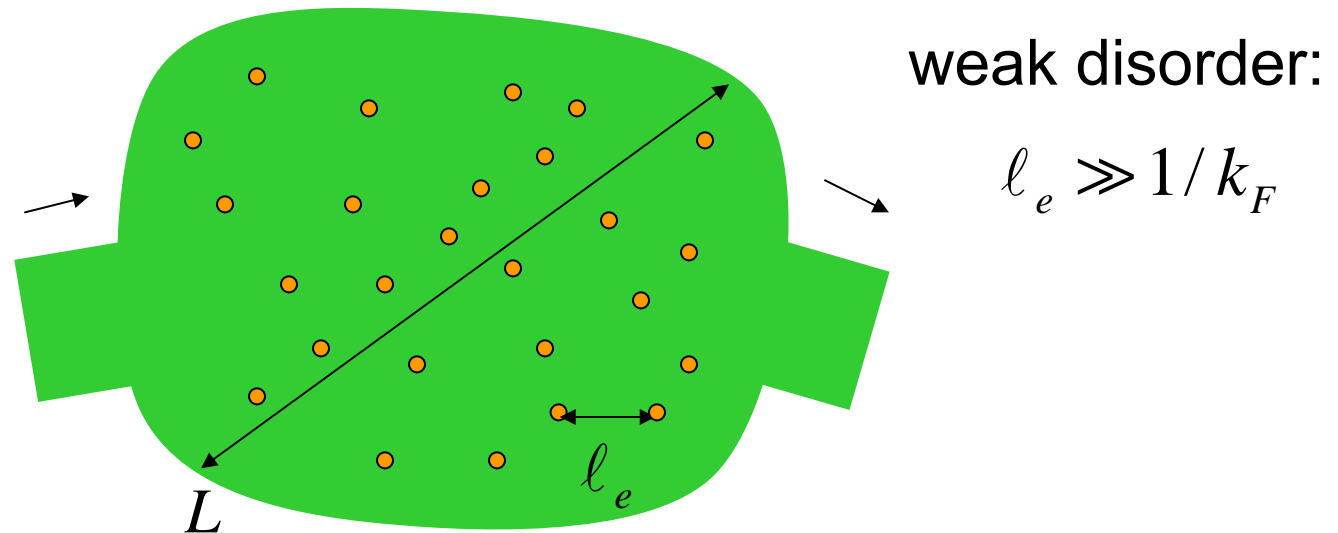
$$k_B T \ll \frac{\hbar v_F}{d}$$

van Wees *et al.*
1988:



waveguide effect!

UP TO WHAT SIZE CAN QUANTUM EFFECTS PERSIST IN PRESENCE OF DISORDER?



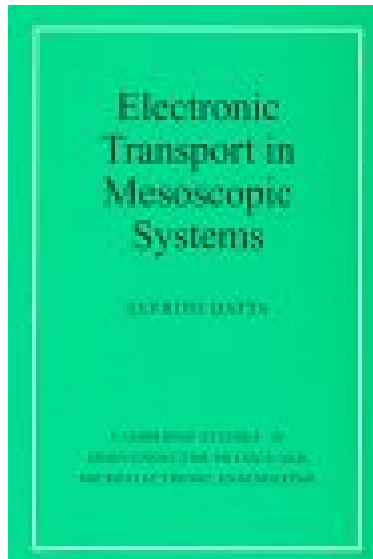
NON-TRIVIAL
RESULT

$$k_B T, \frac{\hbar}{\tau_\varphi} \ll \frac{\hbar v_F}{L} \frac{l_e}{L} \longleftarrow \text{Thouless energy}$$

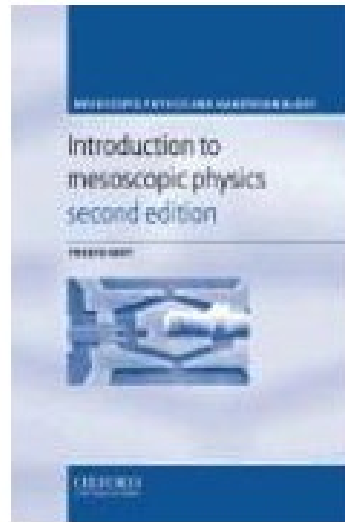
Can be much larger
than level spacing in
box of size L !

$$\frac{\hbar v_F k_F}{(L k_F)^2} \ll \ll \frac{\hbar v_F l_e}{L^2}$$

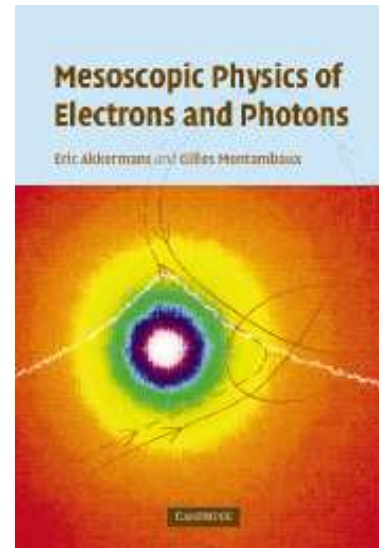
BOOKS ON MESOSCOPIC PHYSICS



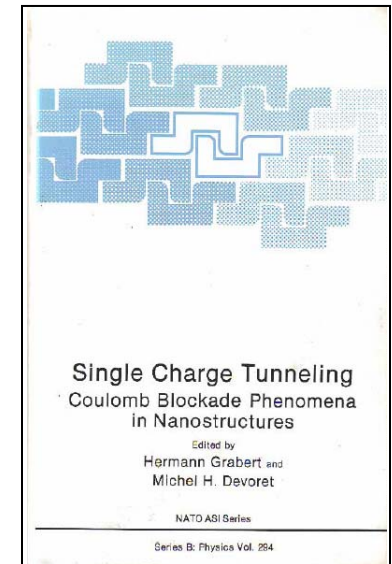
S. Datta



Y. Imry



E. Akkermans and
G. Montambaux



H. Grabert
and M. Devoret

REQUISITES:

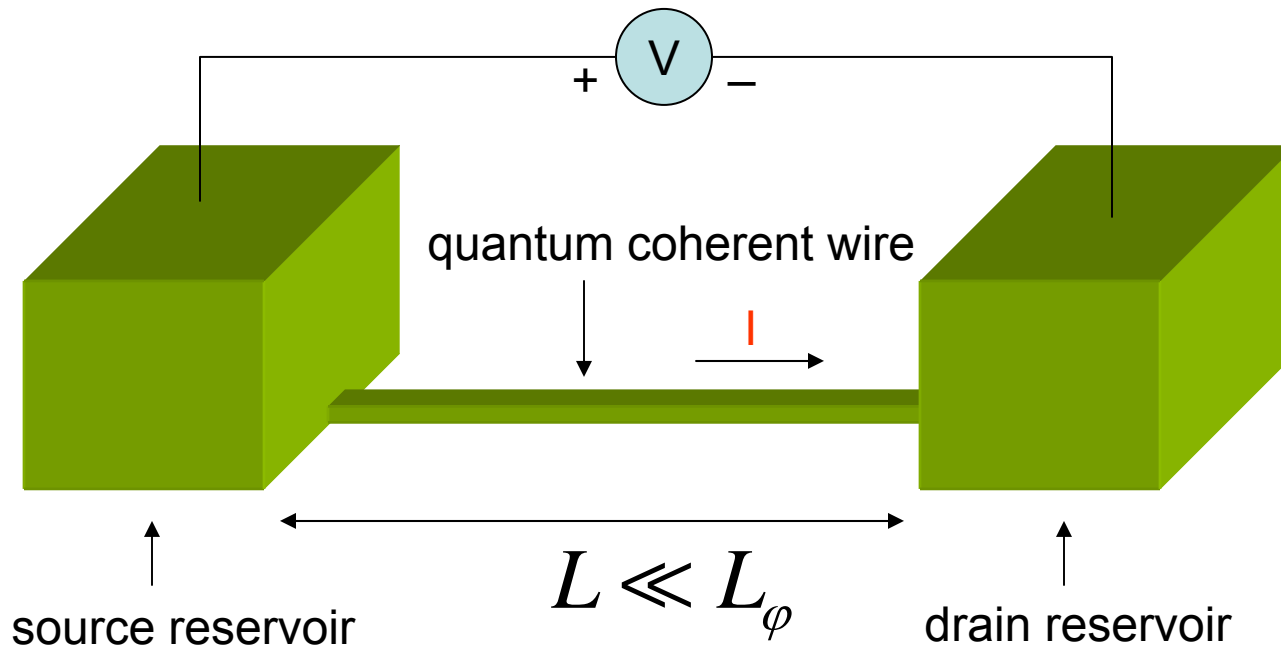
Quantum mechanics
Solid state physics
Statistical mechanics

Useful reference books on quantum statistical physics: Feynman; Schrieffer; Pines and Nozières

2. How mesoscopic physics models the "electron"

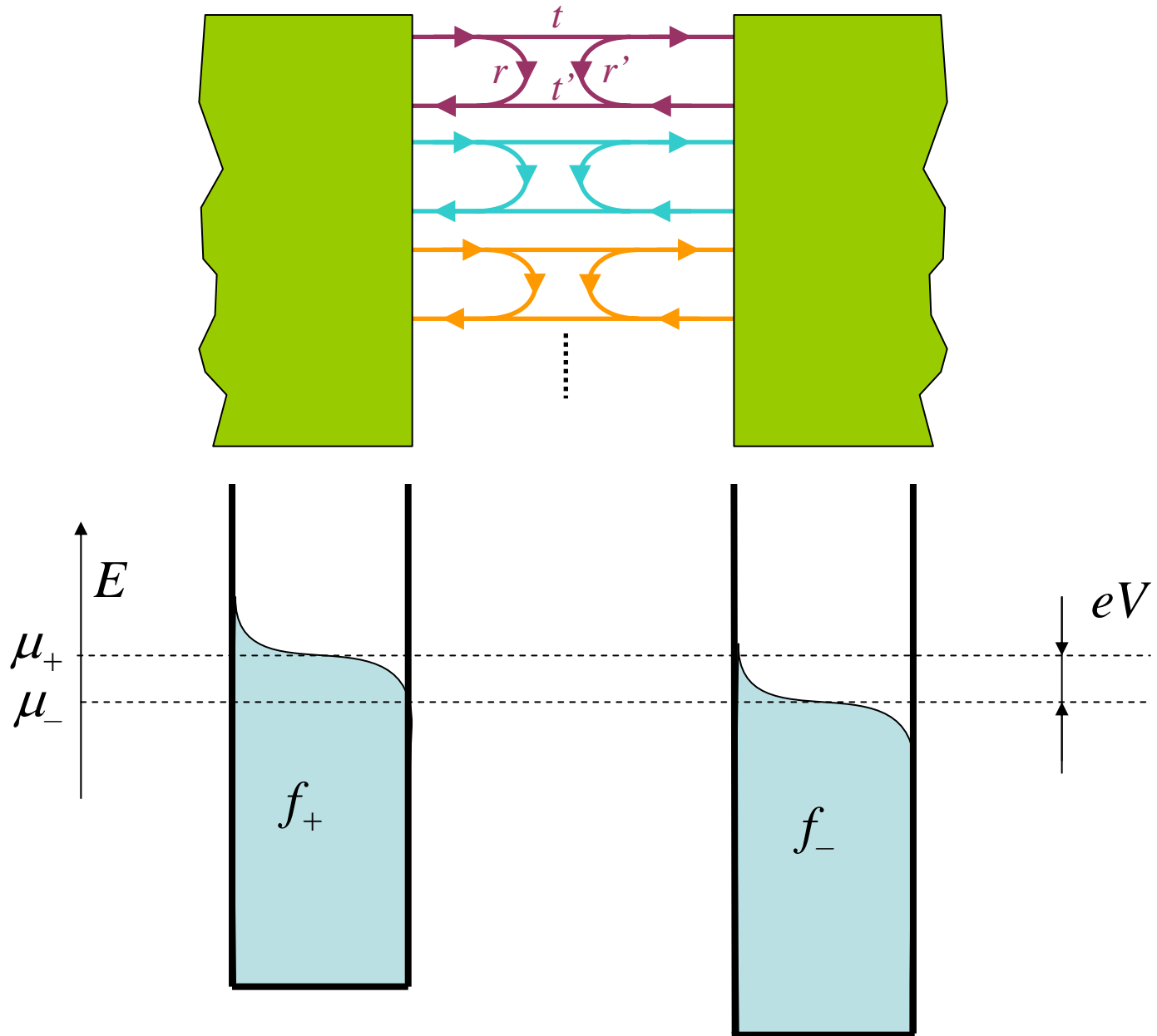
Purpose: provide groundwork for Landauer's approach of transport phenomena

THE MESOSCOPIC RESISTOR



Landauer reservoir: metallic electrode which injects into the quantum coherent region quasi-electron waves with well-defined chemical potential and temperature. It also accepts without reflection any quasi-electron waves and thoroughly recycles them. The Landauer reservoir is to Fermi waves what a black-body is to Bose waves.

Collection of independent channels



THE LANDAUER-BÜTTIKER FORMULA

$$I = I_+ - I_-$$

$$I_{\pm} = \frac{e}{h} \sum_m \int_{-\infty}^{+\infty} f_{\pm}(E) |t_m(E)|^2 dE$$

$$f_{\pm}(E) = \frac{1}{1 + \exp\left(\frac{E - \mu_{\pm}}{k_B T}\right)} \quad \mu_+ - \mu_- = eV$$

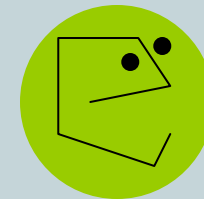
schizophrenic! why does it work?

**RESISTANCE DETERMINED
BY ELASTIC TRANSMISSION
COEFFICIENT!**

WHAT ABOUT JOULE HEATING?

QUESTION: What are the differences between a quasi-electron (a.k.a. dressed electron) and the usual (bare) electron?

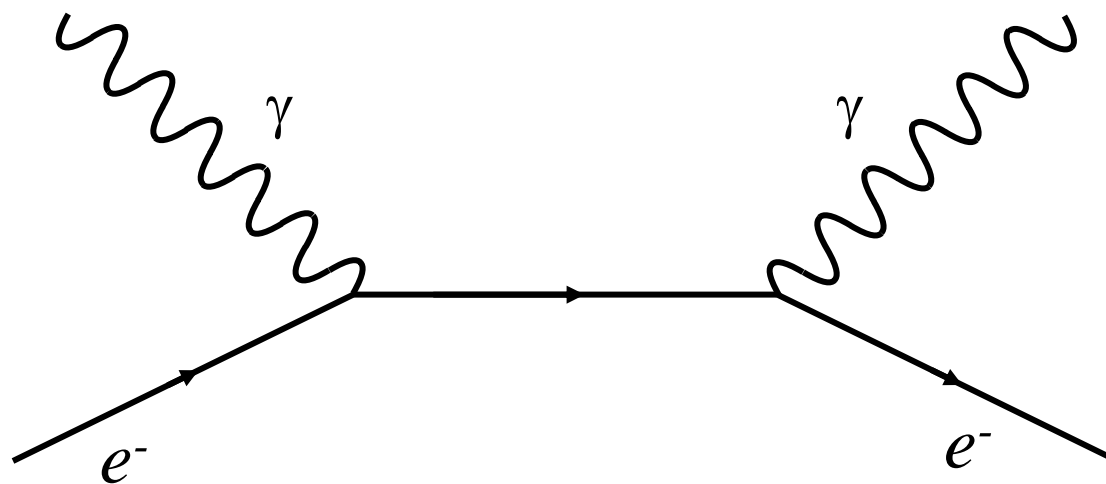
PARTICLE
IDENTIFICATION
CARD



Last Name: Electron	First name: Bare
Address: Vacuum	Genre: Fermion
Occupation: Wave packet	Lifetime: infinite
Average energy: $\hbar\omega$	Average momentum: $\hbar k$
Velocity: $v=d\omega/dk$	Mass: $\hbar dk/dv=m_e$
Charge: $-e$	
Spin: $S=1/2$	Magnetic moment: $g_e S \mu_B$

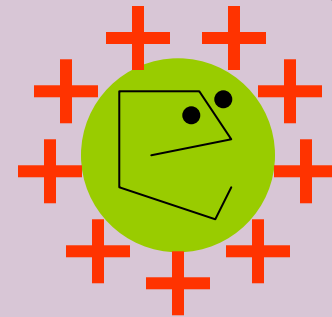
$$g_e = 2 \left[1 + \frac{\alpha}{2\pi} + O(\alpha^2) \right]$$

An example of a Feynman diagram
involving the usual electron and photon
of atomic physics
which propagate in vacuum



THE "ELECTRON" OF MESOSCOPICS

PARTICLE IDENTIFICATION CARD



Last Name: Electron

Address: Metal

Occupation: Wave packet

Average energy: $\hbar\omega$

Velocity: $v=d\omega/dk$

Transverse charge: $-e$

Spin: $S=1/2$

First name: Quasi

Genre: Fermion

Lifetime: finite, except @ k_F

Average momentum: $\hbar k$

Mass: $\hbar dk/dv=m_{\text{eff}}(k)$

Longitudinal charge: 0 ($q \rightarrow 0$)

Magnetic moment: $g_{\text{eff}} S\mu_B$

Definition of the longitudinal and transverse part of a field:

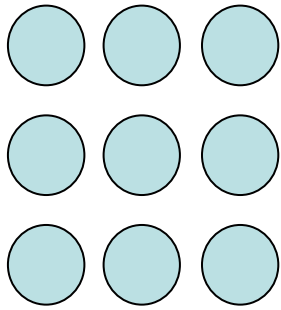
$$\vec{F} = \vec{F}_l + \vec{F}_t$$

$$\vec{\nabla} \cdot \vec{F}_t = 0$$

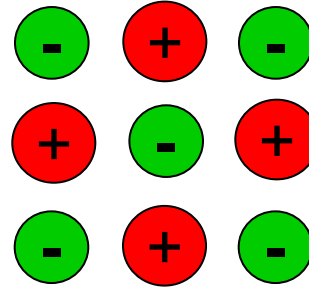
$$\vec{\nabla} \times \vec{F}_l = 0$$

The longitudinal and transverse charges are the sources of the longitudinal and transverse parts of the electrical field, respectively.

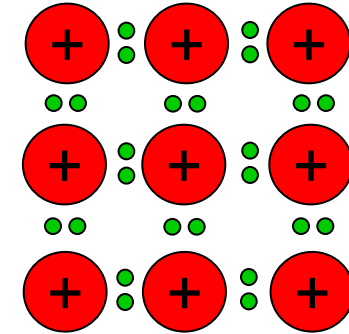
SOLIDS



van der
Waals

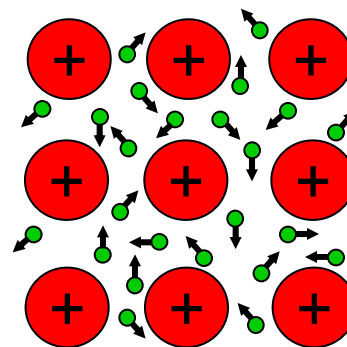


ionic



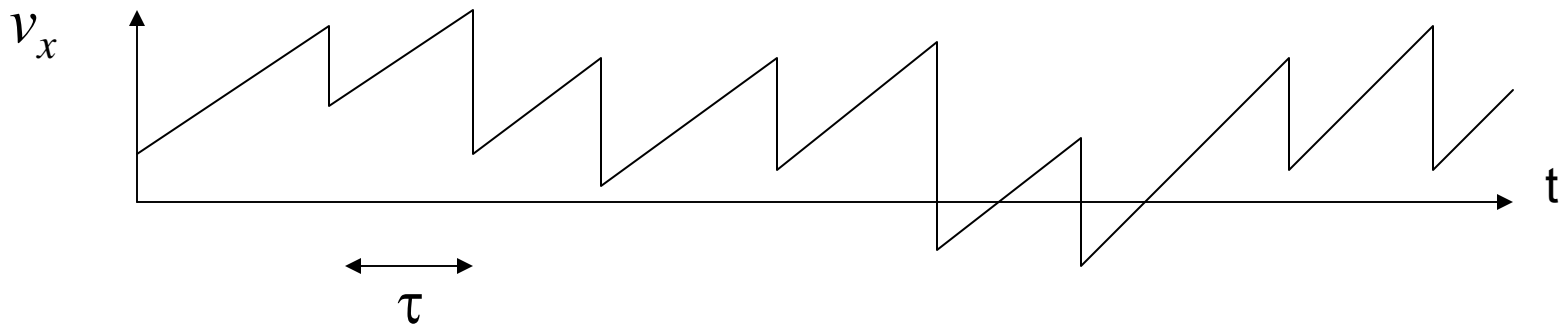
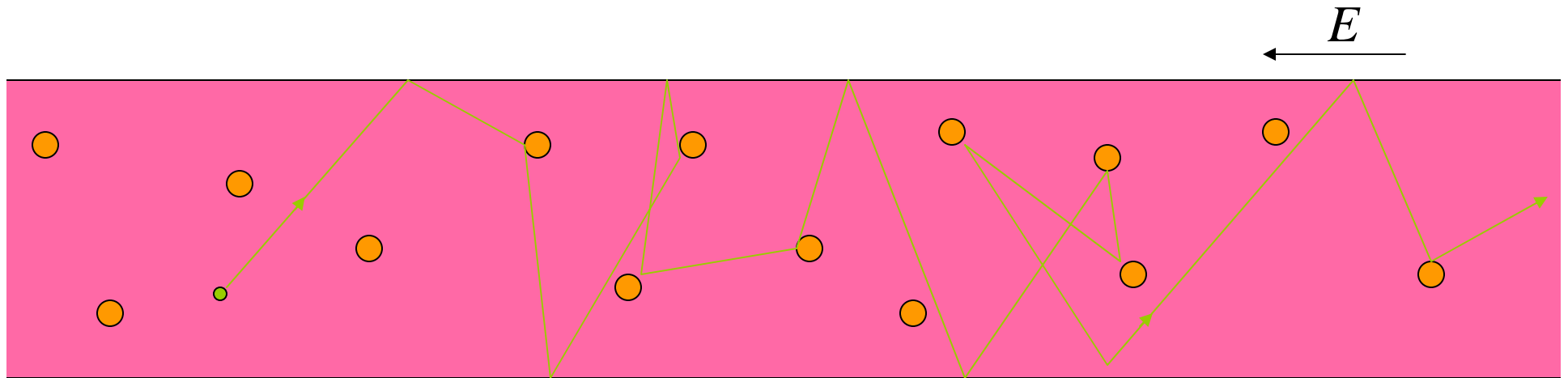
covalent

insulators



metal

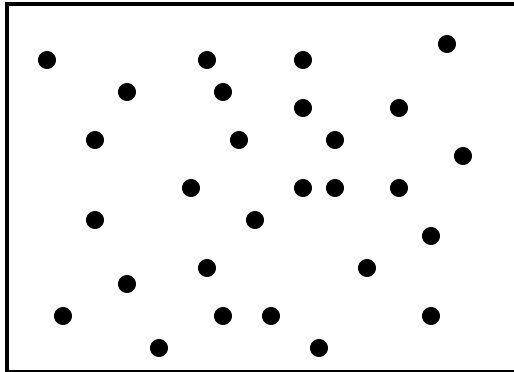
REVIEW OF DRUDE MODEL



$$\langle v_x \rangle = \frac{(-e)E}{m} \tau \quad j = n(-e)\langle v_x \rangle \quad \sigma = \frac{j}{E} = \frac{ne^2}{m} \tau$$

**PROBLEM: ELECTRONS ARE
FERMIONIC WAVES!**

REVIEW OF FREE FERMION GAS MODEL



Box volume V

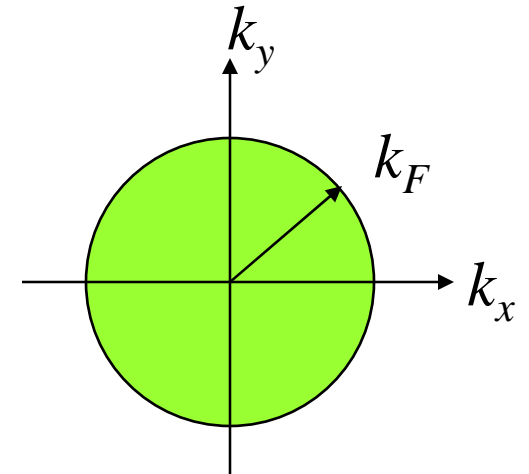
Nb of fermions N

Total energy E_K

Length scale a_0

Energy scale Ry

$$a_0 = 4\pi\epsilon_0 \frac{\hbar^2}{m_e e^2}; Ry = \frac{m_e e^4}{2\hbar^2 \cdot 4\pi\epsilon_0}$$



$$\frac{N}{V} = n = \frac{4\pi}{3} \frac{k_F^2}{(2\pi)^2} = \frac{1}{\frac{4\pi}{3} a^3}; \quad r_s = \frac{a}{a_0}; \quad k_F = \frac{1.92}{r_s a_0}$$

$$\frac{E_K}{N} = \frac{3}{5} \frac{(\hbar k_F)^2}{2m_e} = \frac{3}{5} E_F = \frac{2.22}{r_s^2} Ry$$

$$v_F = \frac{\hbar}{m_e} k_F = v_g \Big|_{E_F}$$

$$\frac{E_C}{N} = \frac{1}{r_s} Ry$$

**OTHER PROBLEM: ELECTRONS
INTERACT STRONGLY!**

**... BUT TOO STRONG INTERACTION
KILLS INTERACTION.....**

Acknowledgements : D. Esteve, H. Pothier, D. Stone
and C. Urbina