



# Mesoscopic Quantum electrodynamics: from atomic-like physics to quantum transport

*Audrey COTTET - LPA – Ecole Normale Supérieure – Paris*

## Theory:

M. C. Dartiailh<sup>1</sup>  
B. Douçot  
T. Kontos  
Z. Legthas

## Experiments:

L.E Bruhat  
T. Cubaynes<sup>2</sup>  
M. C. Dartiailh  
M. Delbecq  
M.M. Desjardins  
T. Kontos  
J.J. Viennot<sup>3</sup>

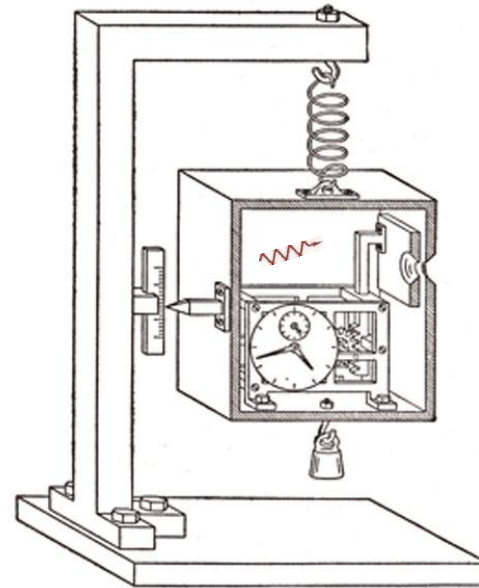
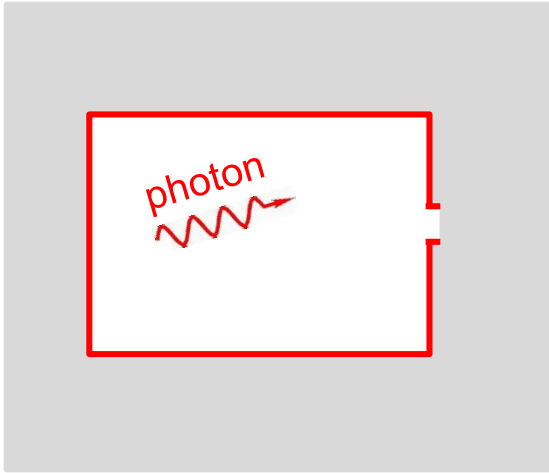
<sup>1</sup>now in New York

<sup>2</sup>now in Karlsruhe

<sup>3</sup>now in Grenoble

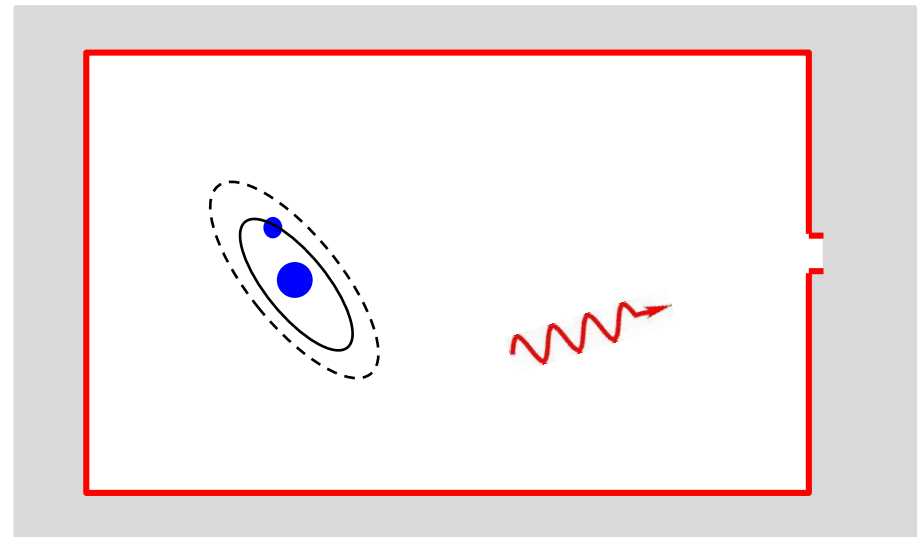
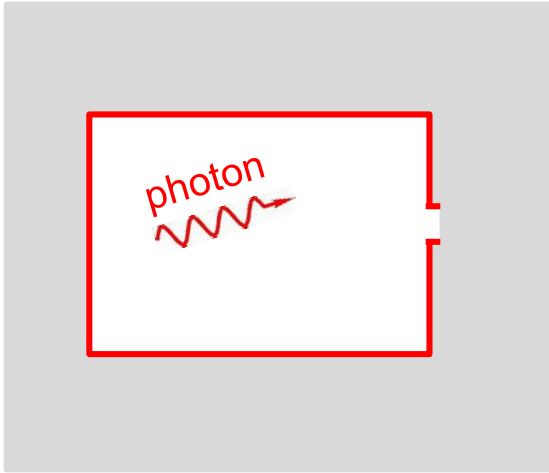


# Photons and atoms in a box



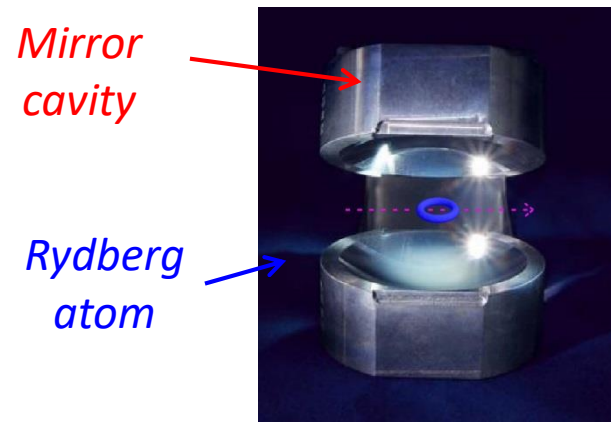
*Einstein's photon box, Solvay 1930  
(scheme by Bohr, 1949)*

# Photons and atoms in a box



# Cavity Quantum ElectroDynamics (QED)

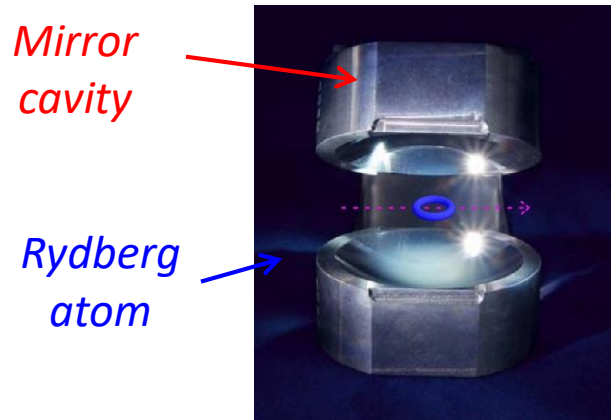
## Atomic cavity QED



*M. Brune et al. PRL (1996)*

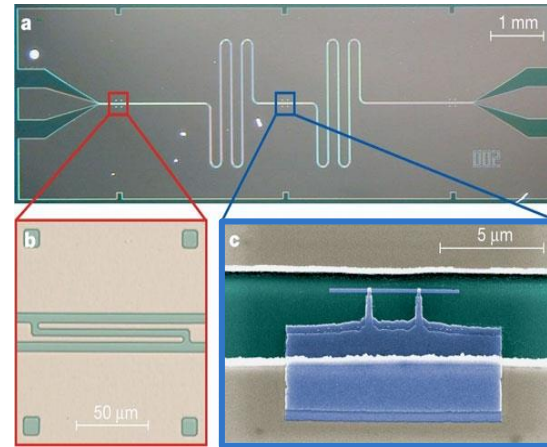
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## Atomic cavity QED



*M. Brune et al. PRL (1996)*

## Circuit QED

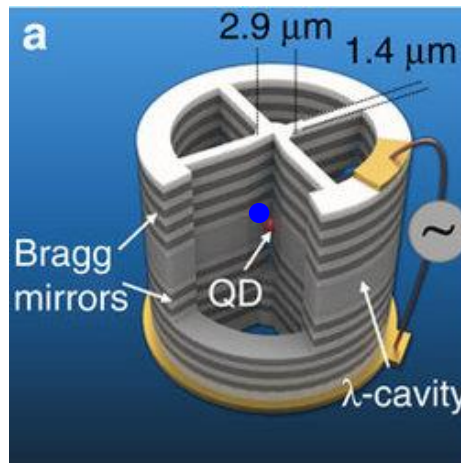


Coplanar waveguide cavity

Superconducting qubit

*A. Wallraff et al, Nature (2004)*

## Optical cavity QED



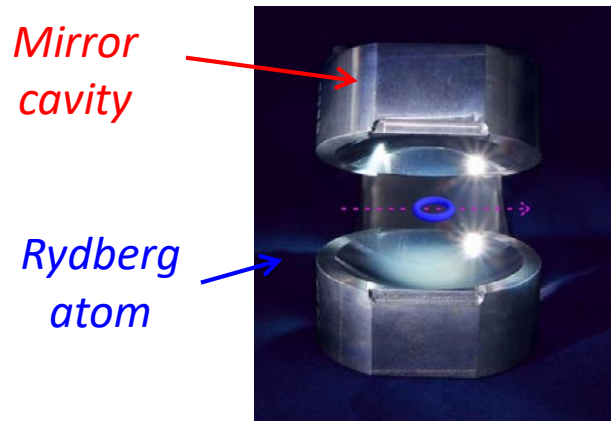
Optical cavity

Self-assembled Quantum dot

*Giesz et al., Nature Comm. (2016)*

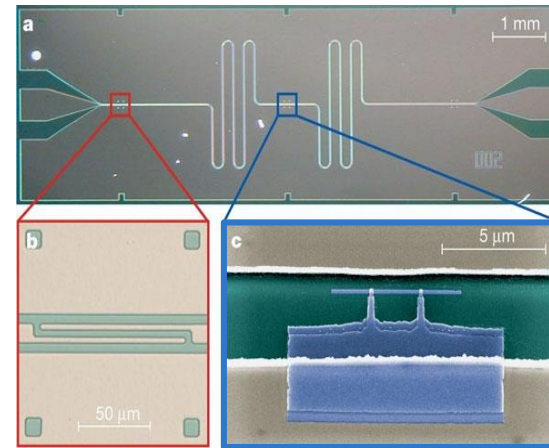
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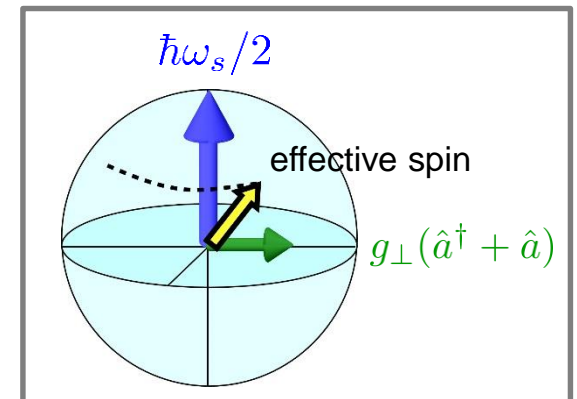
## Jaynes-Cummings Hamiltonian

$$\hat{H} = -\frac{\hbar\omega_s}{2}\hat{\sigma}_z + \hbar\omega_0 \hat{a}^\dagger \hat{a} + g_\perp \hat{\sigma}_x (\hat{a}^\dagger + \hat{a})$$

Two level system

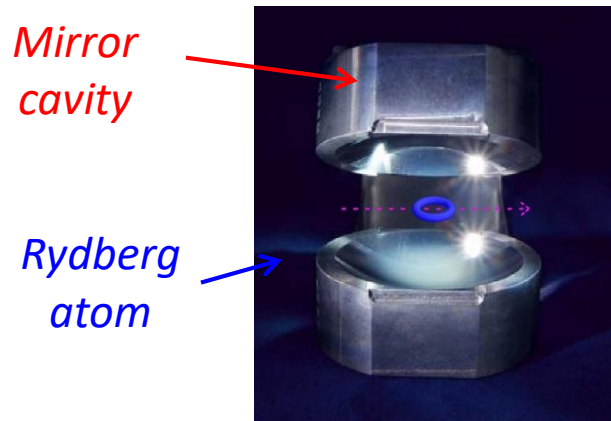
Cavity photons

Transverse coupling



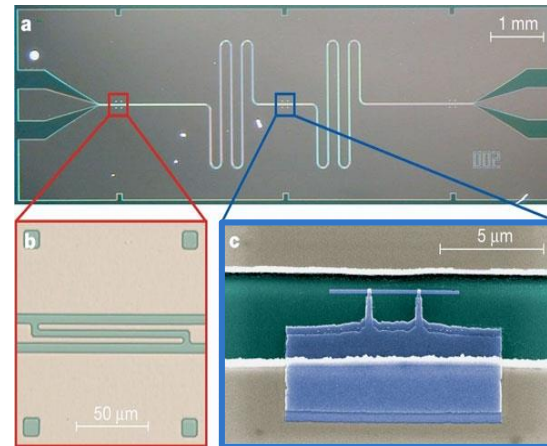
# Cavity Quantum ElectroDynamics (QED)

## Atomic cavity QED



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## Circuit QED



Coplanar waveguide cavity

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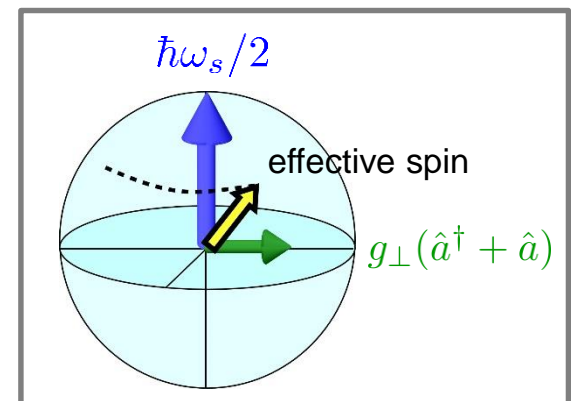
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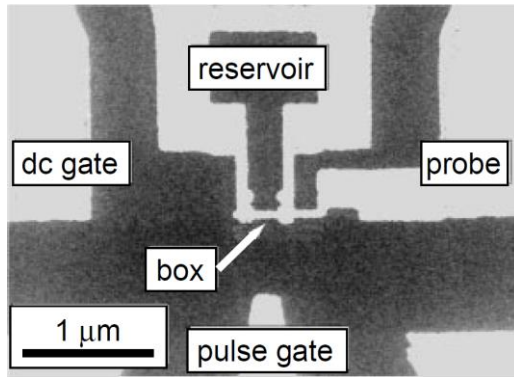
Two level system

Cavity photons

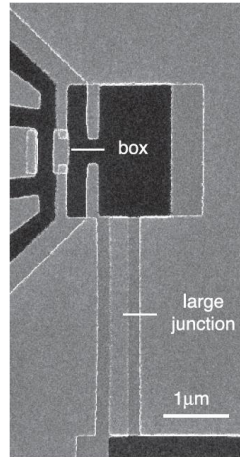
Transverse coupling



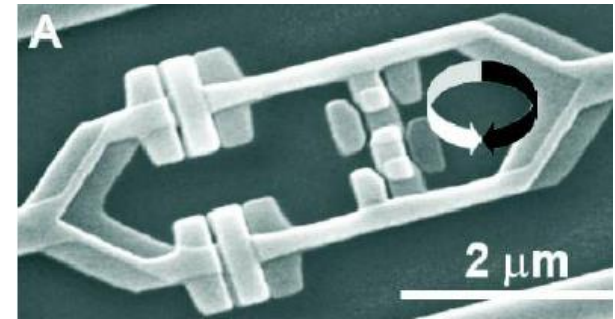
# Different superconducting qubit designs



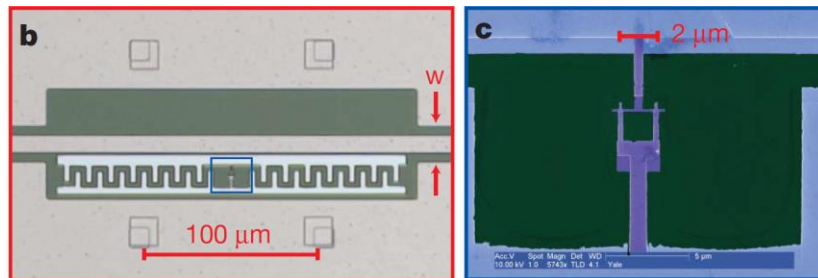
Cooper pair box  
Nakamura et al., Science 1999



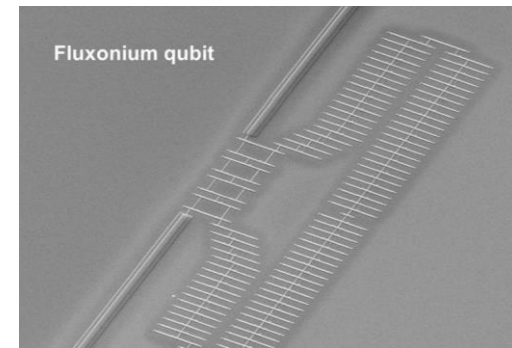
Quantronium  
Vion et al., Science 2002



Flux qubit, Chiorescu et al., Science 2003



Transmon, Schuster et al., Nature 2007

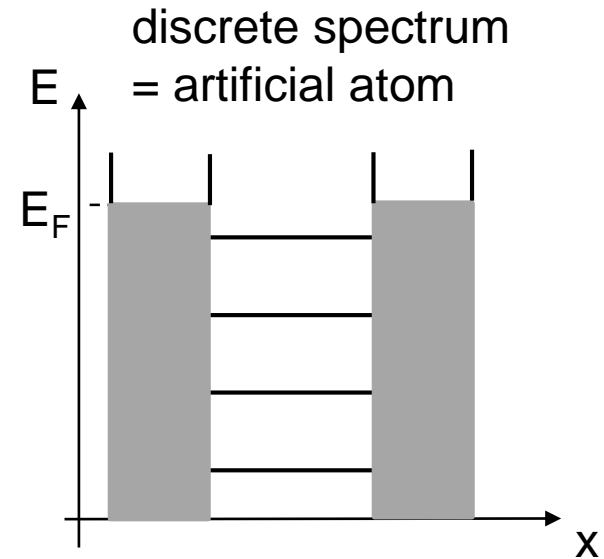
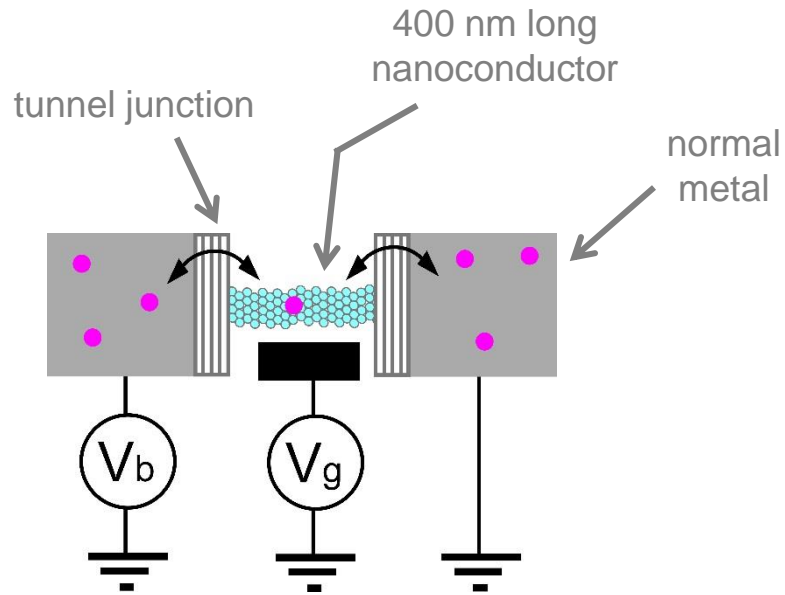


Fluxonium, Manucharyan et al., Science 2009  
Photo: Pop et al., Nature 2014

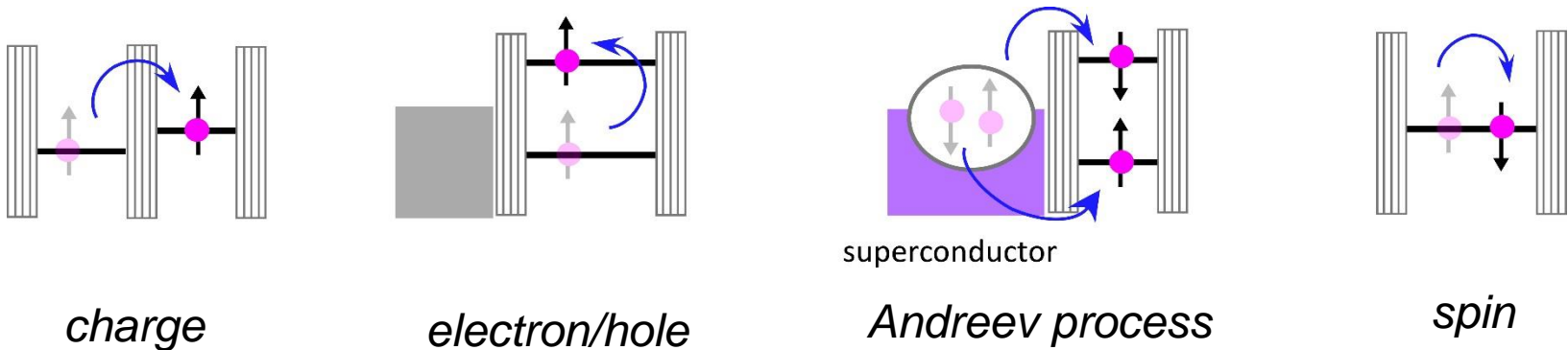
Can we push further the versatility of Circuit QED ?



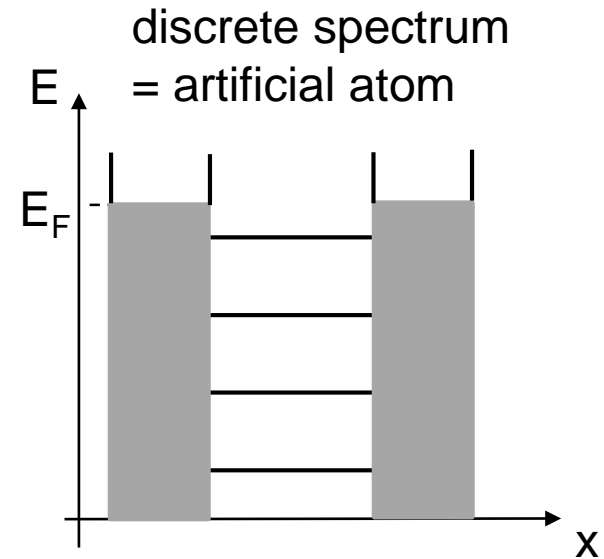
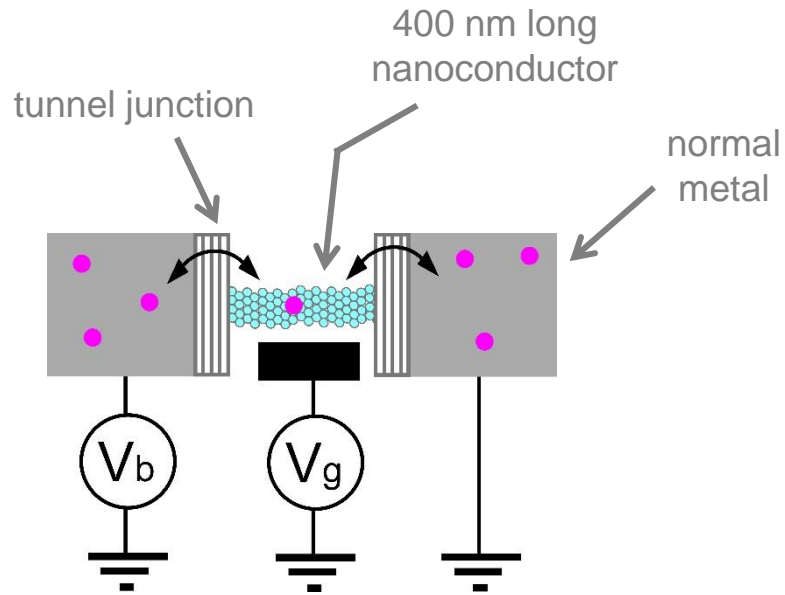
# Hybrid nanocircuits as artificial atoms



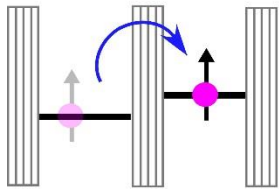
Various degrees of freedom available:



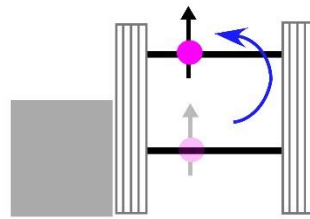
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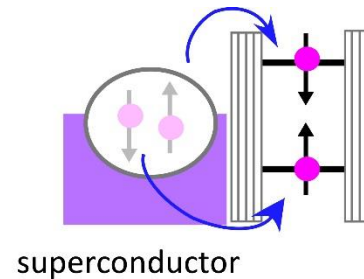
Various degrees of freedom available:



*charge*

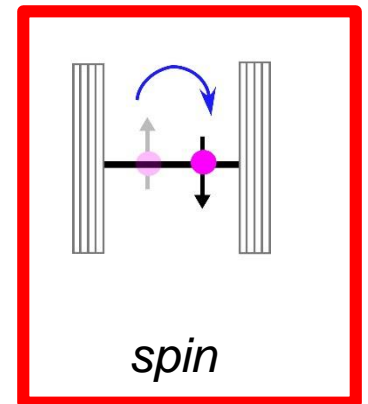


*electron/hole*



superconductor

*Andreev process*



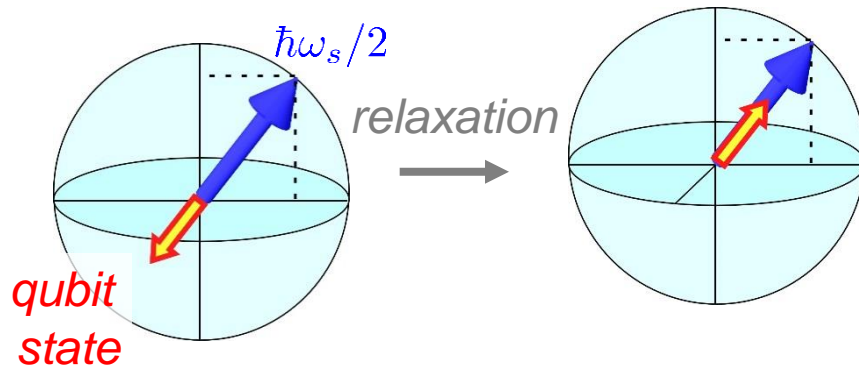
*spin*

# Quantum dot circuits potentialities?

## Superconducting qubits :

Coherence limited by relaxation  $T_1 < 100 \mu\text{s}$

*Review: Devoret and Schoelkopf, Science (2013)*

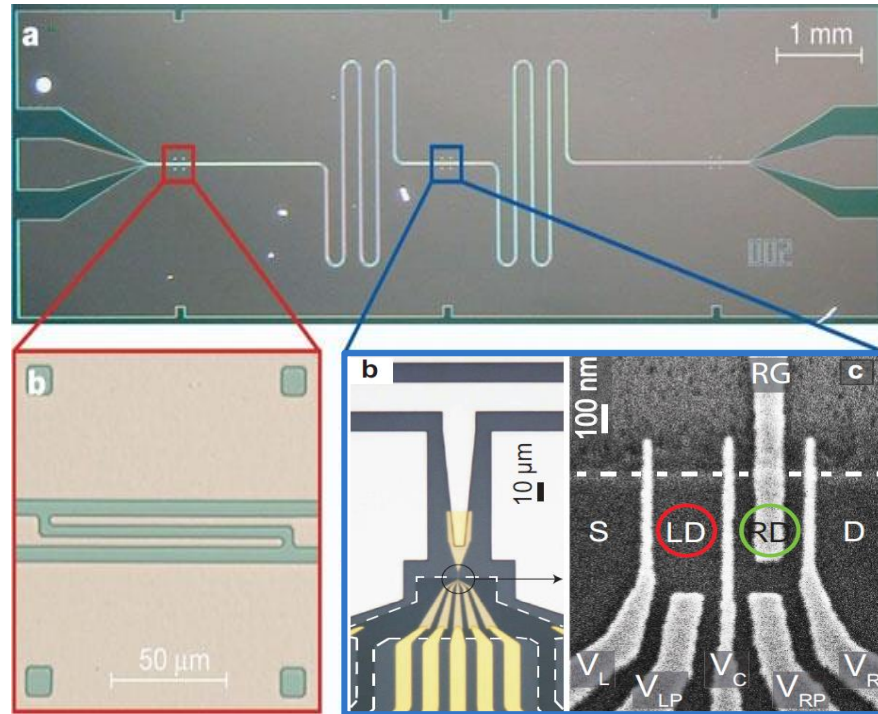


## Spin in a quantum dot :

- GaAs (single dot) @ 130mK:  
 $T_1 \sim 80 \text{ ms}$  *Scarolino et al., PRL 2014*
- Si/SiGe (double dot) @ 15mK:  
 $T_1 \sim 3 \text{ s}$  *Prance et al., PRL 2012*
- Carbon nanotube (bulk) @ 4K:  
 $T_1 \sim 170 \mu\text{s}$  *Rice et al. PRB 2013*

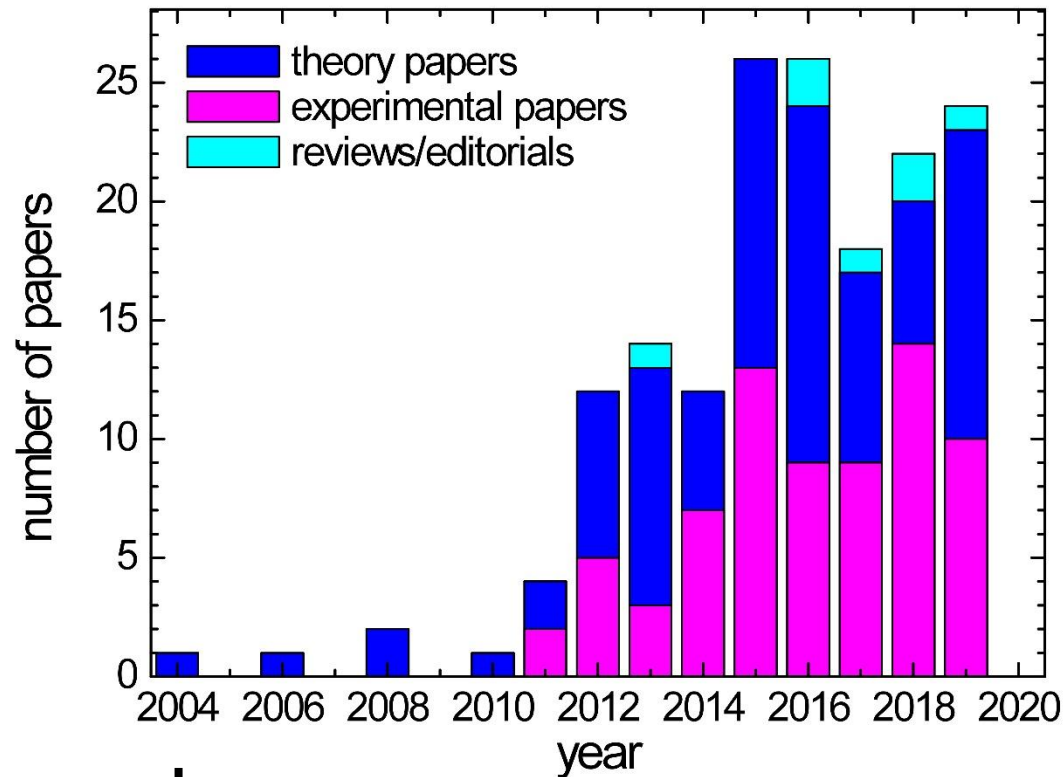
⇒ quantum dot circuit or a hybrid nanocircuit in a microwave cavity?

# The Mesoscopic QED architecture



- Nanoconductors: Carbon nanotubes, 2DEGs, semiconducting nanowires, graphene, atomic contacts...
- Different types fermionic reservoirs (normal, superconducting, ferromagnetic)
- Many circuit geometries possible

# Mesoscopic QED : research activity



## Experiments:

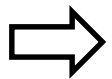
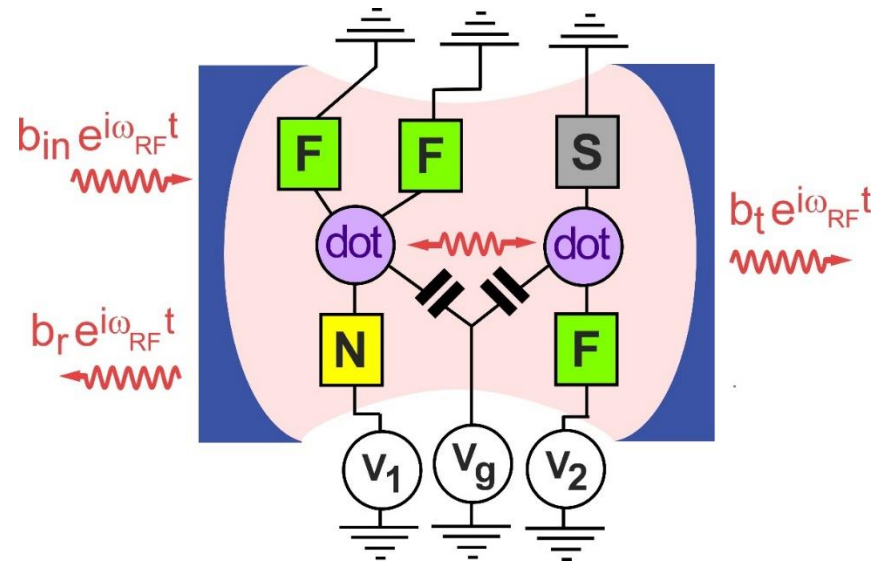
Delft  
Heifei  
Paris  
Princeton  
Yale  
Zurich  
....

Pioneering  
theory paper:  
*Childress et al.,  
PRA (2004)*

Single dot with N contacts:  
*Delbecq et al., PRL (2011)*  
*Frey et al., APL (2011)*

Double dot with N contacts:  
*Frey et al., PRL (2012)*  
*Petersson et al., Nature (2012)*

# Interests of Mesoscopic QED

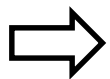
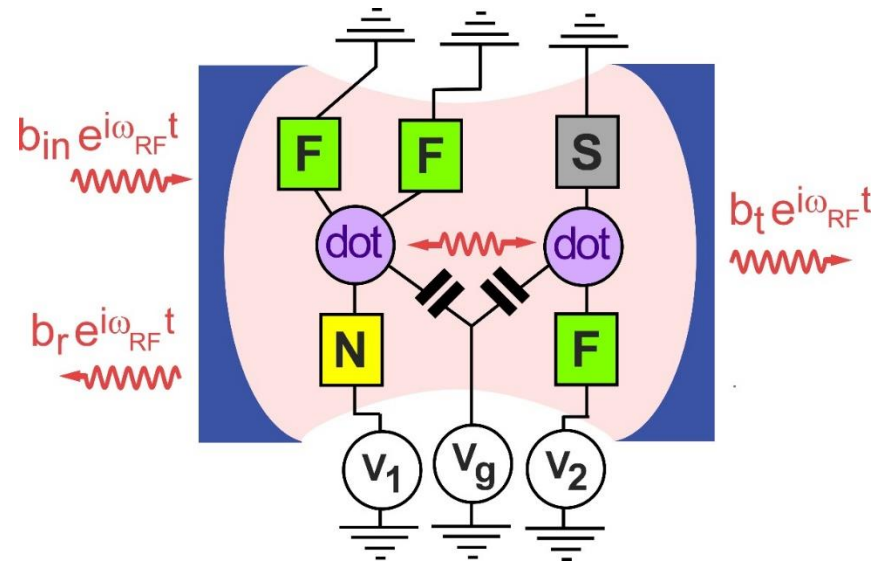


Artificial atom limit:

Transfer ideas of circuit QED to probe/couple/manipulate *new degrees of freedom*

*Dot orbitals, electronic spins...*

# Interests of Mesoscopic QED



## Artificial atom limit:

Transfer ideas of circuit QED to probe/couple/manipulate *new degrees of freedom*

*Dot orbitals, electronic spins...*



## Open system limit:

Study electronic transport and condensed matter problems

*DC currents and cavity response provide different information*

New ways to manipulate cavity state?

# OUTLINE

## **I. Description of a microwave cavity coupled to a nanocircuit**

## **II. A nanocircuit as an artificial atom in a cavity**

- Strong charge/photon coupling
- Strong spin/photon coupling

## **III. Condensed matter problems and electron transport probed with cavity photons**

- Semiclassical cavity response modified by a dissipative nanocircuit
- Photon emission by a dot/superconductor junction
- Majorana bound states in a cavity

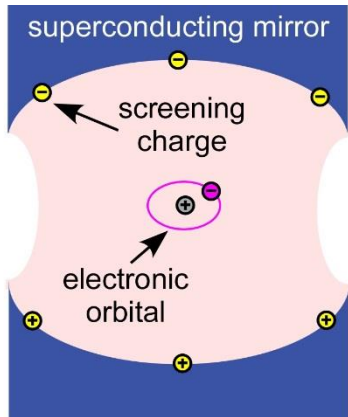
## **III. Using a dissipative nanocircuit to prepare non-classical cavity states**

- Quantum non-linear description of Mesoscopic QED
- Preparation of a photonic Schrödinger cat using dissipative tunneling

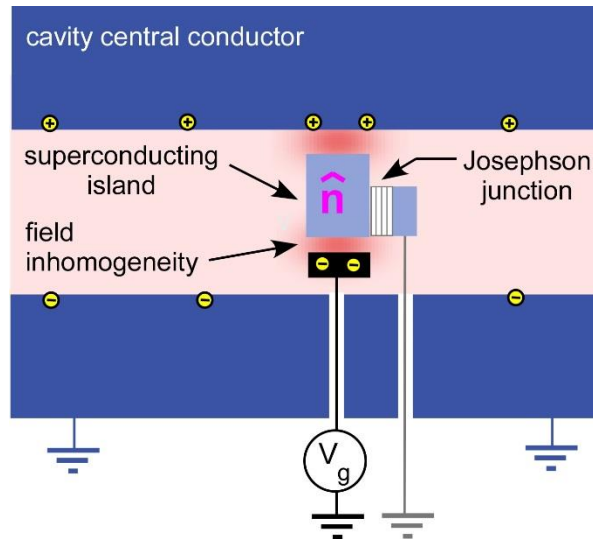


# Comparison between different cavity QED setups

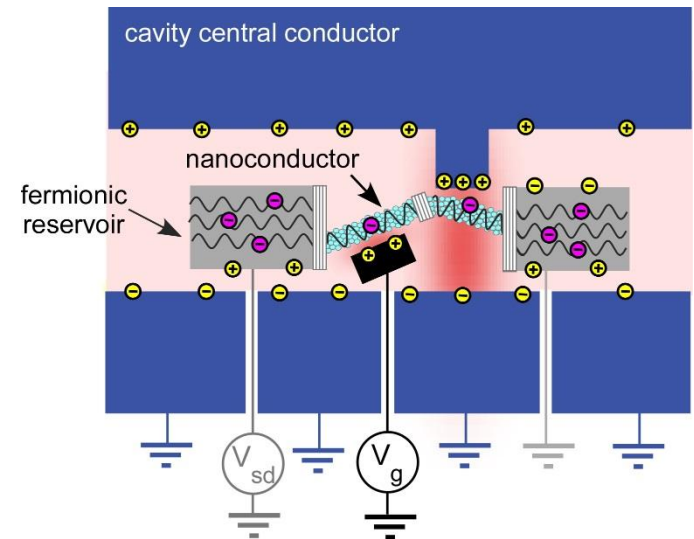
## Cavity QED: *with atoms*



## Circuit QED: *with superconducting quantum bits*



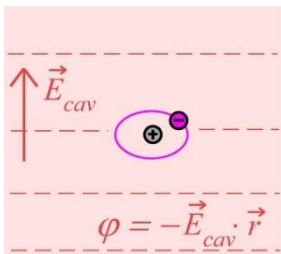
## Mesoscopic QED: *with hybrid nanocircuits*



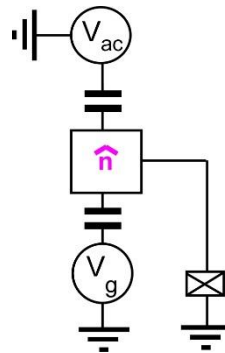
### Models

Electric-dipole approximation

$$H_{int} = -\vec{d} \cdot \vec{E}_{cav}$$



Circuit model  
(electrical nodes)

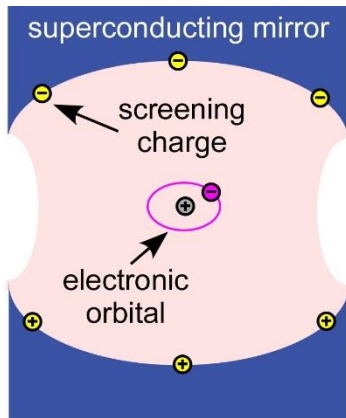


- Orbital fermionic degrees of freedom
- Inhomogeneous cavity field
- Tunneling physics

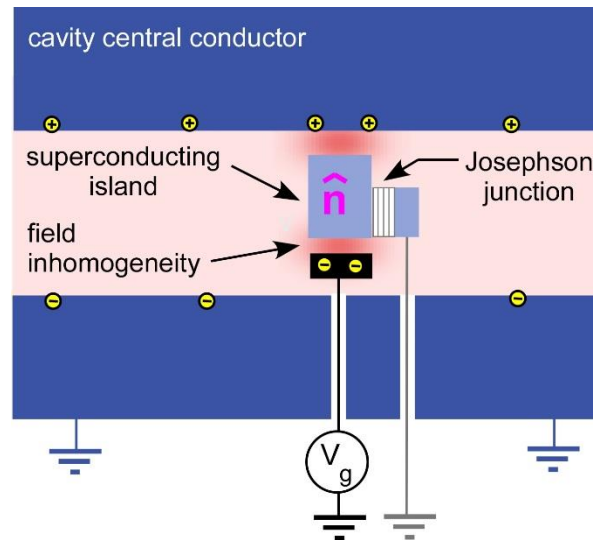
?

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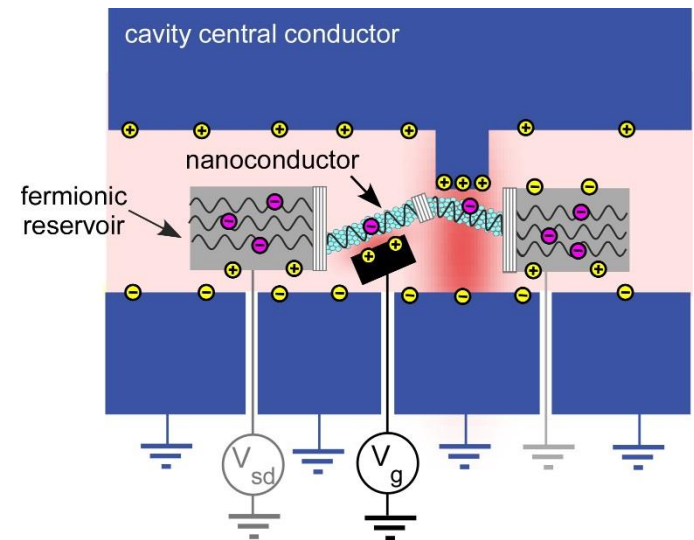
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## Circuit QED: *with superconducting quantum bits*



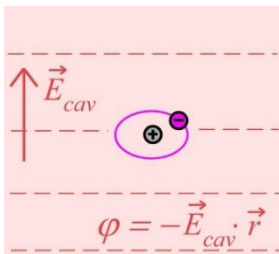
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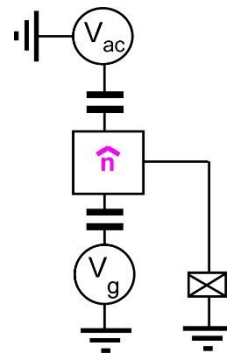
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- Orbital fermionic degrees of freedom
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- Tunneling physics

Approximations:

- plasmonic dynamics ultra-fast
- cavity magnetic field disregarded

# Electron/photon electric coupling scheme

Cottet, Kontos and Douçot, *Phys. Rev. B* 91, 205417 (2015)

$$\begin{aligned} \widetilde{\hat{H}}_{tot} = & \int d^3r \hat{\psi}^\dagger(r) \hat{h}_{\mathcal{T}}(\vec{r}) \hat{\psi}(r) + \hat{H}_{Coul} + \hbar\omega_0 \hat{a}^\dagger \hat{a} \\ & + \hat{\mathcal{V}}(\hat{a} + \hat{a}^\dagger) + (\hat{\mathcal{V}}^2/\hbar\omega_0) + \int d^3r \left( \Delta(\vec{r}) \hat{\psi}_\uparrow^\dagger(\vec{r}) \hat{\psi}_\downarrow^\dagger(\vec{r}) + H.c. \right) \end{aligned}$$

*linear electron/photon coupling (scalar photonic potential)*

$$\begin{cases} \hat{h}_{\mathcal{T}}(\vec{r}) = -\hbar^2 \Delta_{\vec{r}}/2m - e\Phi_{harm}(\vec{r}) - eV_{conf}(\vec{r}) \\ \hat{\mathcal{V}} = -e \int d^3r V_\perp(\vec{r}) \hat{\psi}^\dagger(r) \hat{\psi}(r) \end{cases}$$

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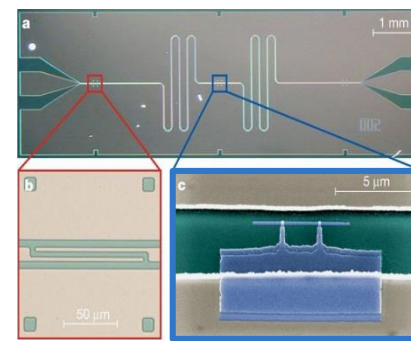
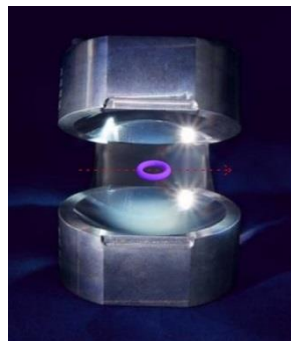
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This description bridges between **cavity QED** and **circuit QED** (two limit cases)

*dipolar approximation* ←



→ *electrical circuit model*

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Cottet, Kontos and Douçot, *Phys. Rev. B* 91, 205417 (2015)

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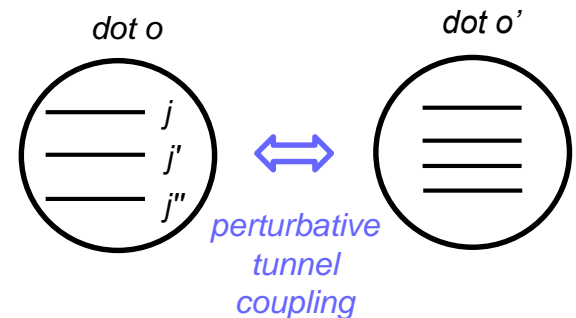
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Tunnel model:

$$\hat{\psi}^\dagger(\vec{r}) = \sum_{o,j} \underbrace{\varphi_{oj}^*(\vec{r})}_{\text{wavefunction for orbital } j \text{ in dot/reservoir } o} \hat{c}_{oj}^\dagger$$

*wavefunction for orbital j in dot/reservoir o*

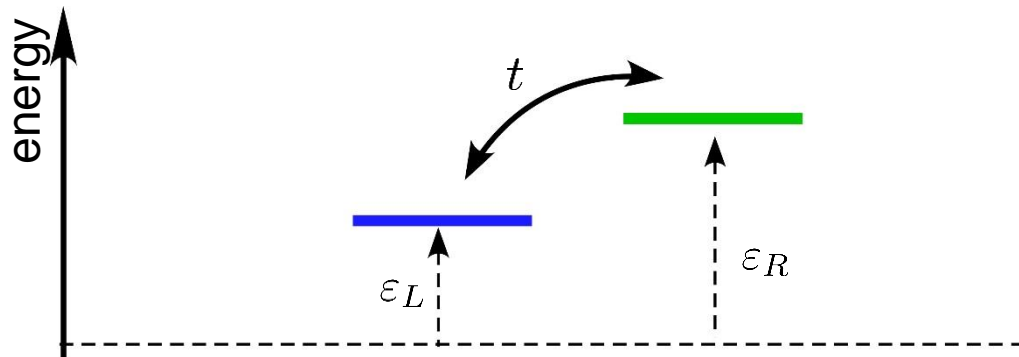
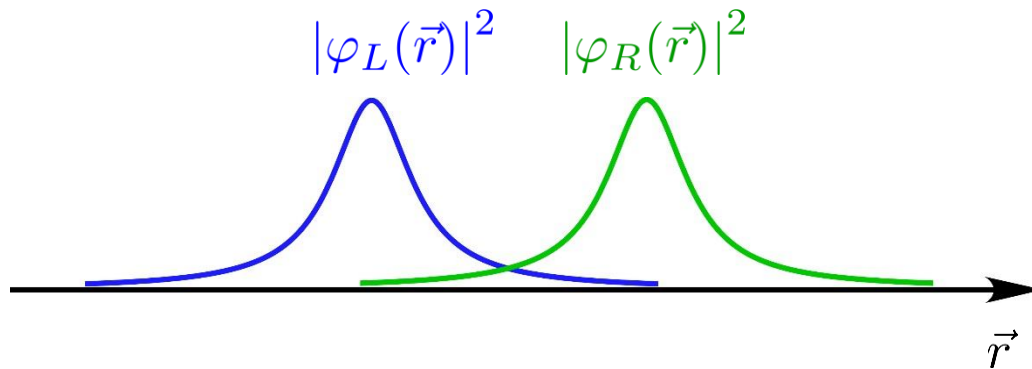


Prange, *Phys. Rev.* 131, 1083 (1963)

# Tunnel model of Mesoscopic QED

Cottet, Kontos and Douçot, *Phys. Rev. B* 91, 205417 (2015)

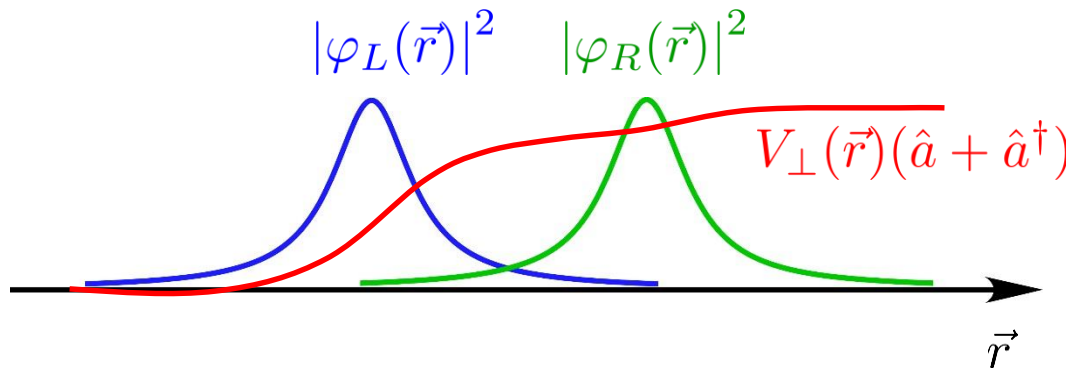
$$\hat{H}_{DC} = \varepsilon_L \hat{c}_L^\dagger \hat{c}_L + \varepsilon_R \hat{c}_R^\dagger \hat{c}_R + t \hat{c}_L^\dagger \hat{c}_R + t^* \hat{c}_R^\dagger \hat{c}_L$$



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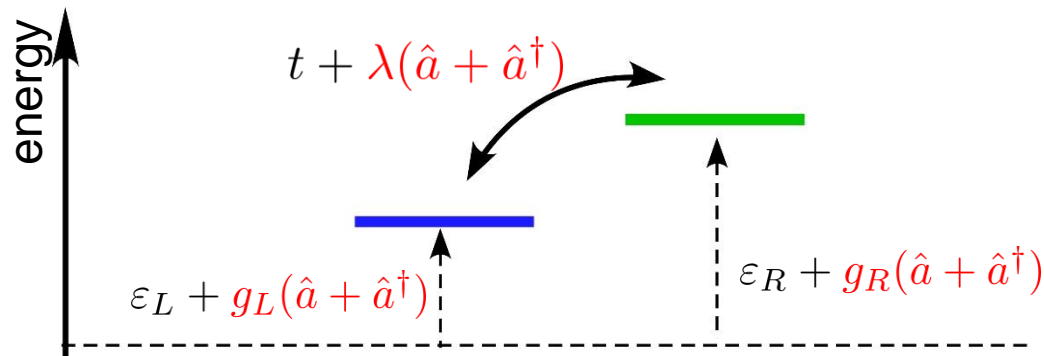
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$$g_L = -e \int d^3r |\varphi_L(\vec{r})|^2 V_\perp(\vec{r})$$

$$g_R = -e \int d^3r |\varphi_R(\vec{r})|^2 V_\perp(\vec{r})$$

$$\lambda = -e \int d^3r \varphi_L^*(\vec{r}) \varphi_R(\vec{r}) V_\perp(\vec{r})$$

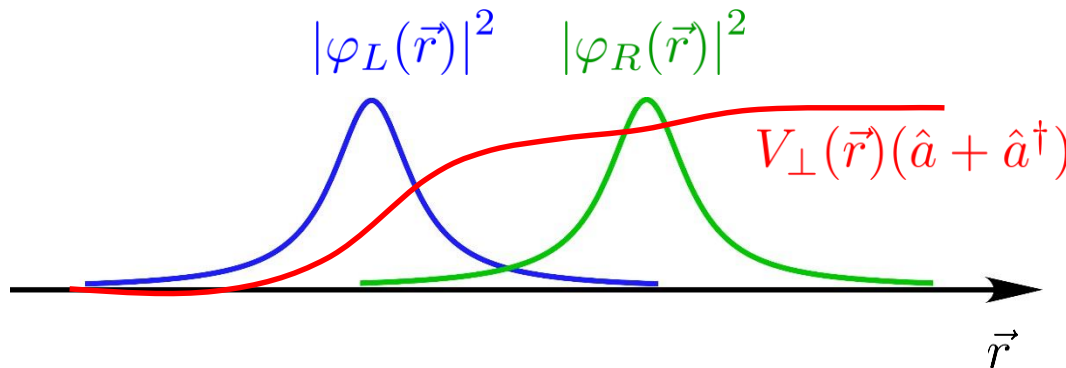


$$\begin{aligned} \hat{H}_{AC} = & g_L(\hat{a} + \hat{a}^\dagger) \hat{c}_L^\dagger \hat{c}_L \\ & + g_R(\hat{a} + \hat{a}^\dagger) \hat{c}_R^\dagger \hat{c}_R \\ & + \lambda(\hat{a} + \hat{a}^\dagger) \hat{c}_L^\dagger \hat{c}_R + H.c. \end{aligned}$$

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Cottet, Kontos and Douçot, *Phys. Rev. B* 91, 205417 (2015)

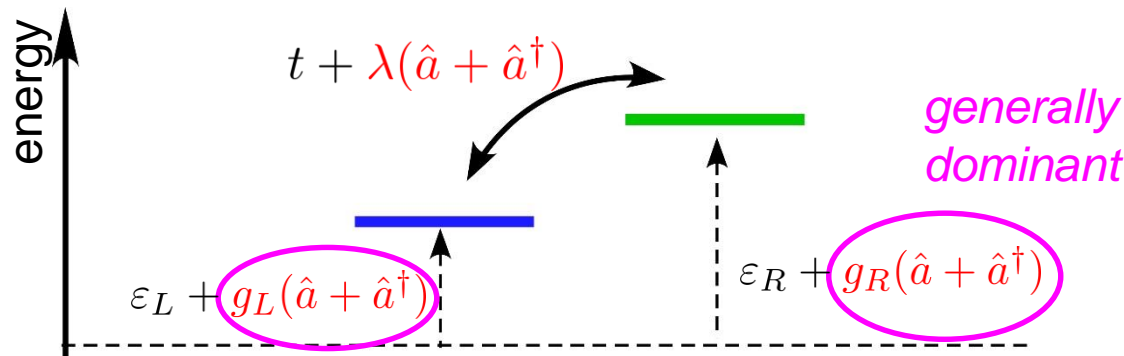
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$$g_R = -e \int d^3r |\varphi_R(\vec{r})|^2 V_\perp(\vec{r})$$

$$\lambda = -e \int d^3r \varphi_L^*(\vec{r}) \varphi_R(\vec{r}) V_\perp(\vec{r})$$



$$\hat{H}_{AC} = g_L (\hat{a} + \hat{a}^\dagger) \hat{c}_L^\dagger \hat{c}_L + g_R (\hat{a} + \hat{a}^\dagger) \hat{c}_R^\dagger \hat{c}_R + \lambda (\hat{a} + \hat{a}^\dagger) \hat{c}_L^\dagger \hat{c}_R + H.c.$$



# OUTLINE

I. Description of a microwave cavity coupled to a nanocircuit

 II. A nanocircuit as an artificial atom in a cavity

- Strong charge/photon coupling
- Strong spin/photon coupling

III. Condensed matter problems and electron transport probed with cavity photons

- Semiclassical cavity response modified by a dissipative nanocircuit
- Photon emission by a dot/superconductor junction
- Majorana bound states in a cavity

III. Using a dissipative nanocircuit to prepare non-classical cavity states

- Quantum non-linear description of Mesoscopic QED
- Preparation of a photonic Schrödinger cat using dissipative tunneling

# OUTLINE

I. Description of a microwave cavity coupled to a nanocircuit

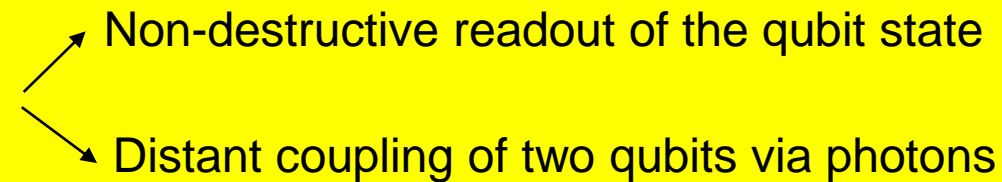
→ II. A nanocircuit as an artificial atom in a cavity

- Strong charge/photon coupling

- Strong spin/photon coupling

III. C

Artificial atom limit: designs to reach the **strong coupling limit** ?

$g_{\perp} \gg \Gamma_s, \Lambda_0$  

- Majorana bound states in a cavity

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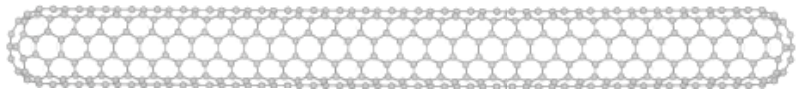
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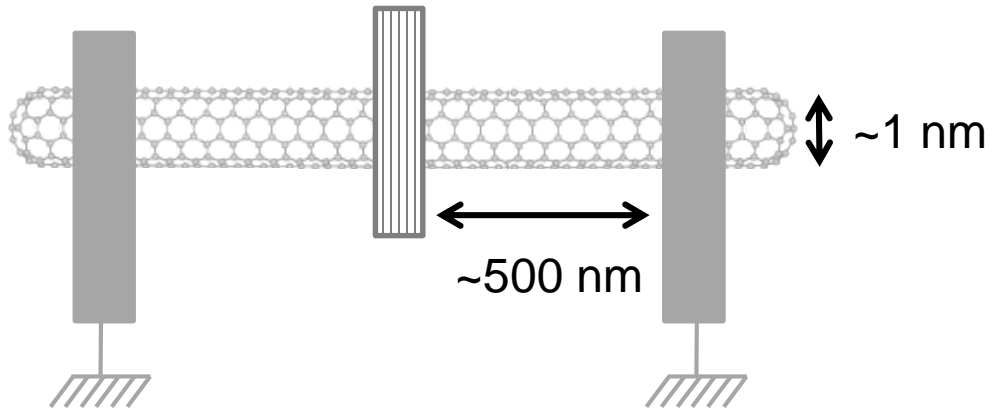
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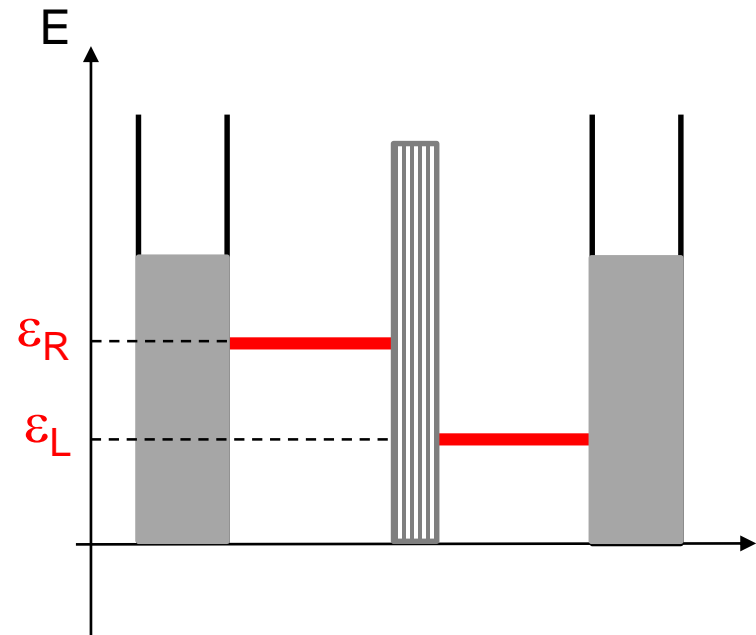
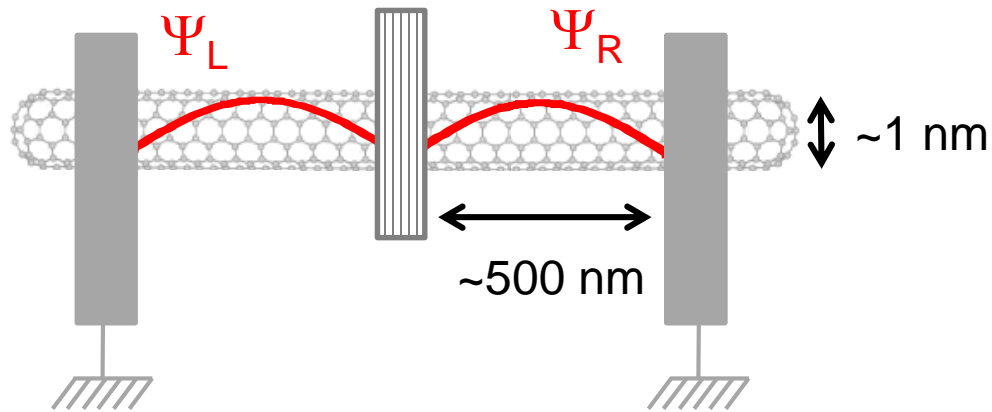
# Double quantum dot (DQD) in a carbon nanotube



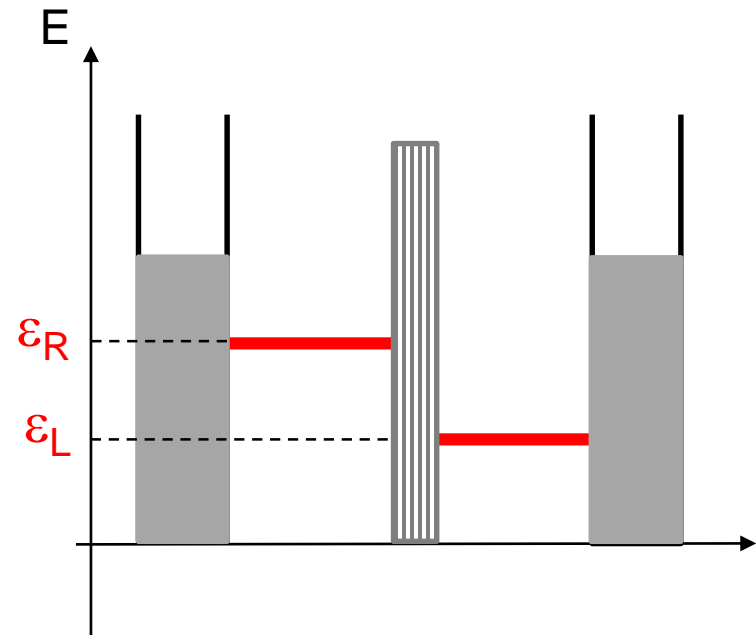
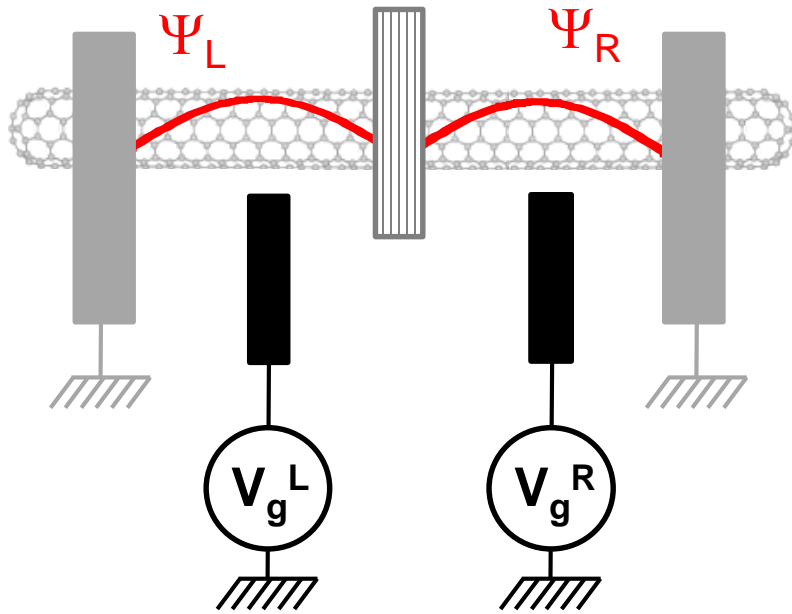
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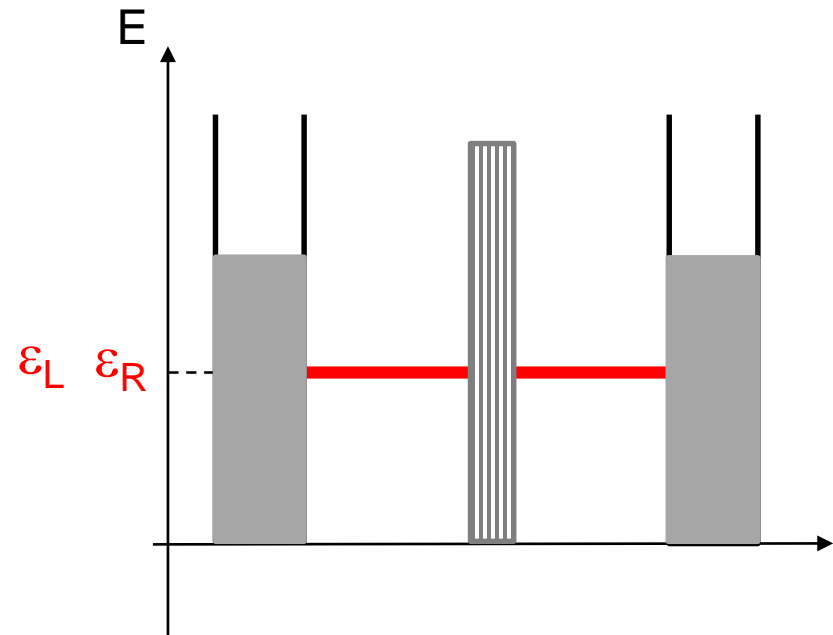
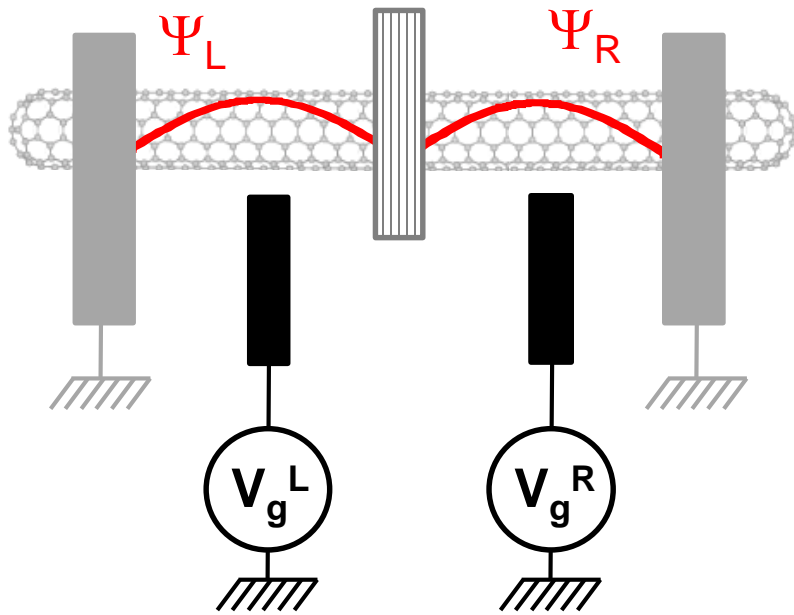
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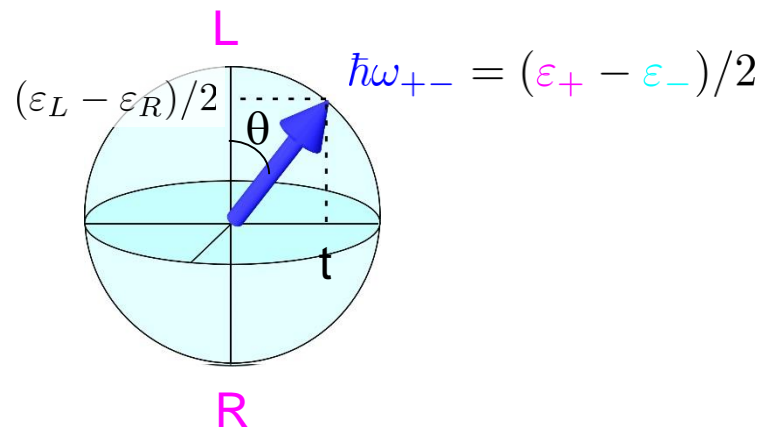
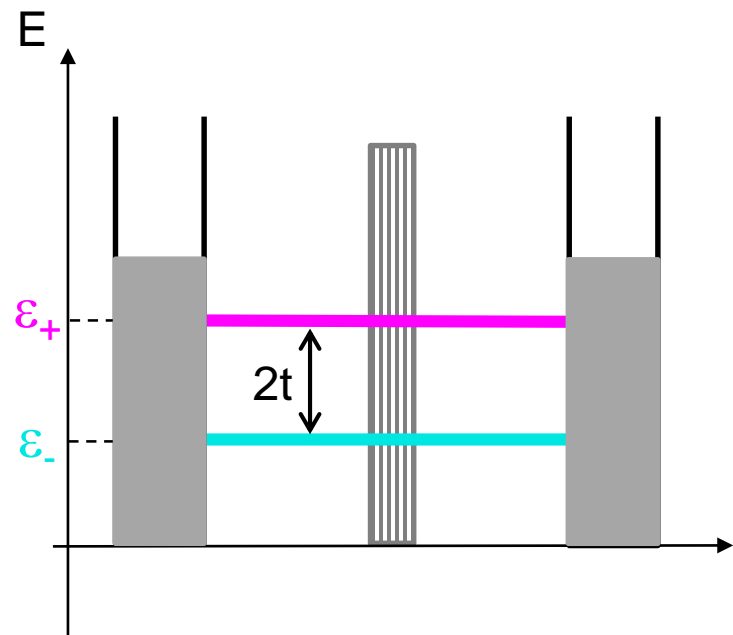
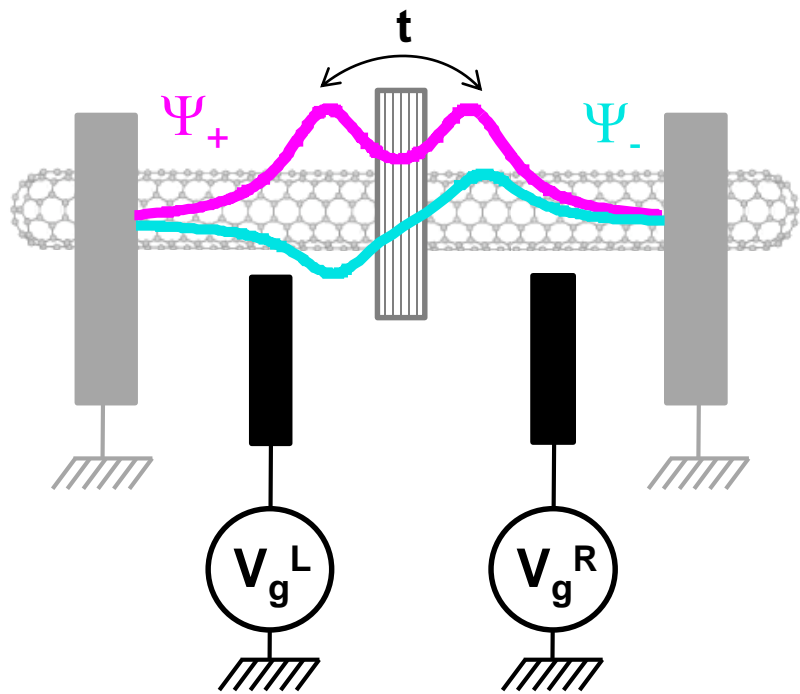


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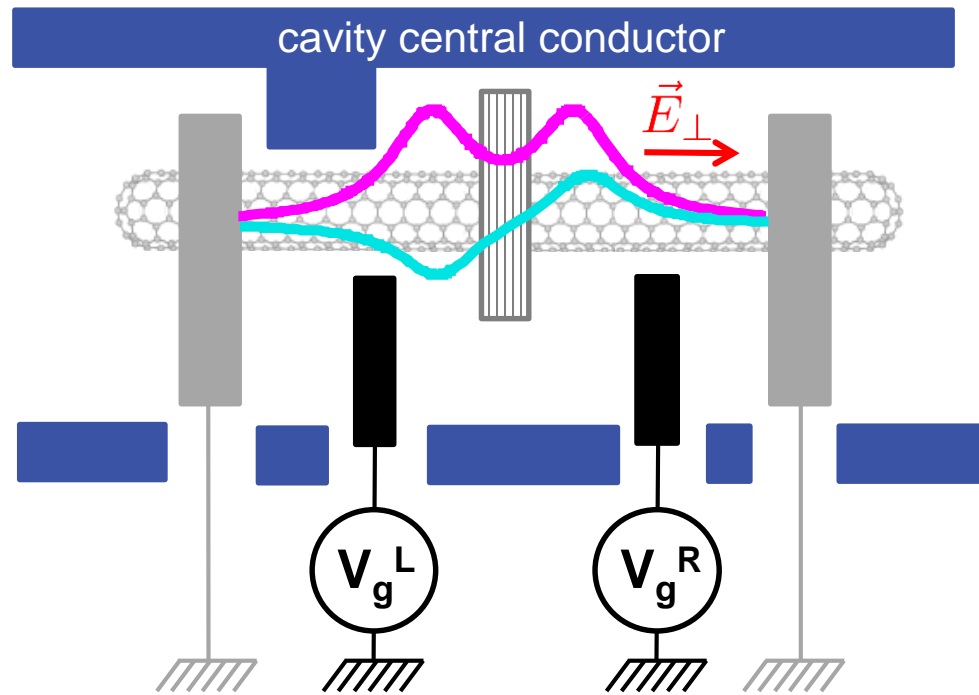




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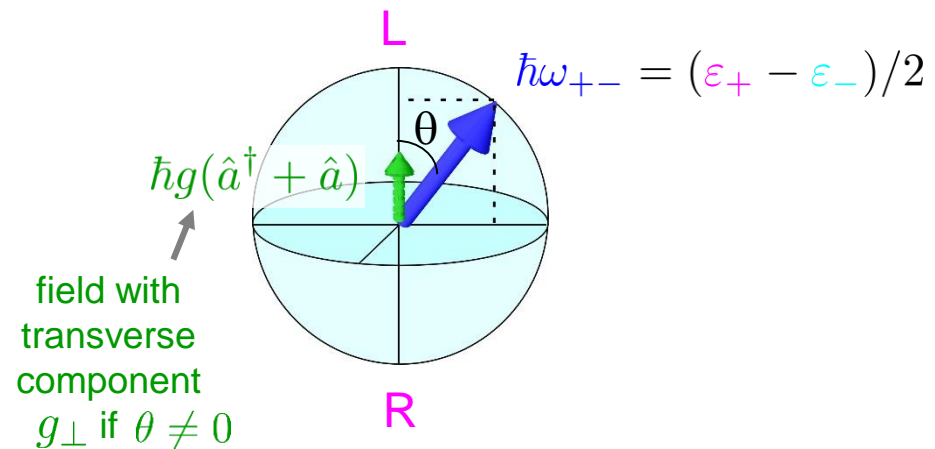
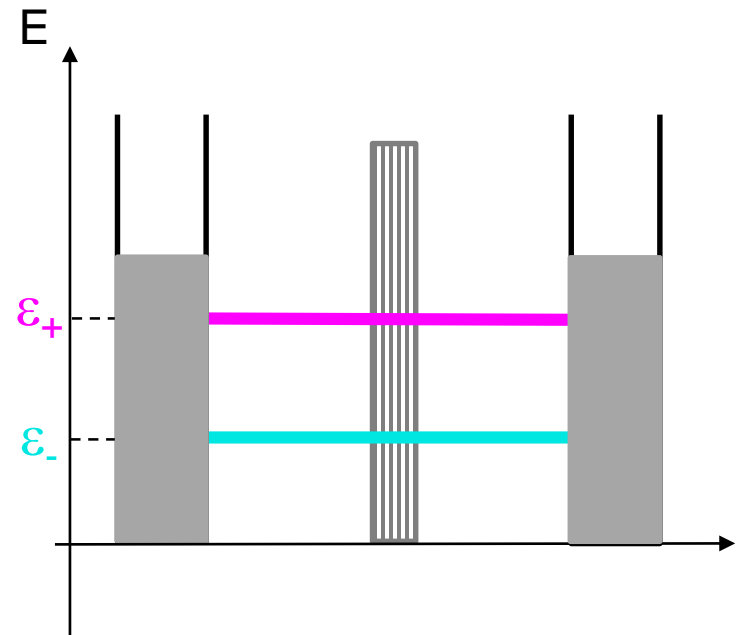
# DQD coupled to a microwave cavity



$$E \sim (\hat{a} + \hat{a}^\dagger) \times 1 \text{V.m}^{-1}$$

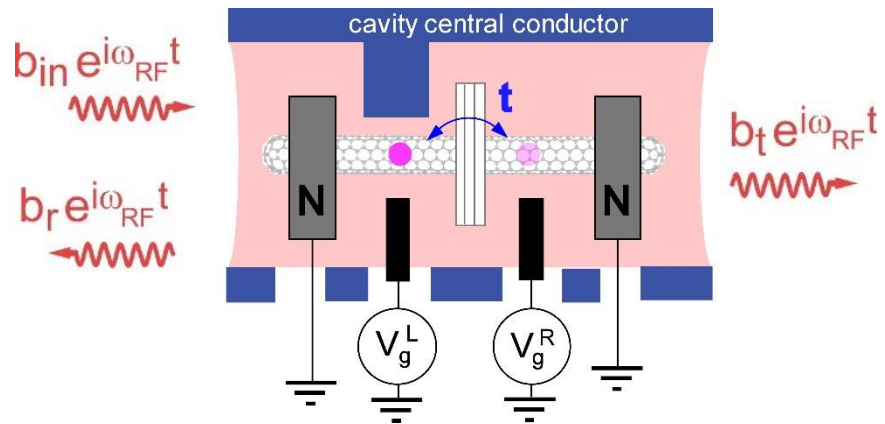
$$g_\perp = \frac{g_R - g_L}{2} \frac{2t}{\sqrt{4t^2 + (\epsilon_L - \epsilon_R)^2}}$$

$\sim 100 \text{ MHz}$



# Cavity response in presence of the DQD

*Bruhat, Cubaynes, Viennot, Dartiailh, Desjardins, Cottet, and Kontos, PRB 2018*



$$\frac{b_t}{b_{in}} = ?$$

# Cavity+DQD, semiclassical linear response model

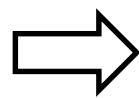
$$\hat{H}_{tot} = \omega_0 \hat{a}^\dagger \hat{a} + \hat{H}_{DD} + (g_L \hat{c}_L^\dagger \hat{c}_L + g_R \hat{c}_R^\dagger \hat{c}_R)(\hat{a} + \hat{a}^\dagger) + \hat{H}_{bath} + (\hat{a} \varepsilon_{in} e^{-i\omega_{RF}t} + H.c.)$$

⇒  $\frac{d\hat{a}}{dt} = -i\omega_0 \hat{a} - \Lambda_0 \hat{a} - i(g_L \hat{c}_L^\dagger \hat{c}_L + g_R \hat{c}_R^\dagger \hat{c}_R) - i\varepsilon_{in} e^{-i\omega_{RF}t}$

Semi-classical approximation:  $\hat{a} = \langle \hat{a} \rangle \simeq \bar{a} e^{-i\omega_{RF}t}$

Linear response:  $\hat{c}_j^\dagger \hat{c}_j = a_0 + \sum_j \chi_{j,j'}(\omega_{RF}) e^{-i\omega_{RF}t} + \sum_j \chi_{j,j'}(\omega_{RF}) e^{i\omega_{RF}t}$  for  $j \in \{L, R\}$

resonant approximation

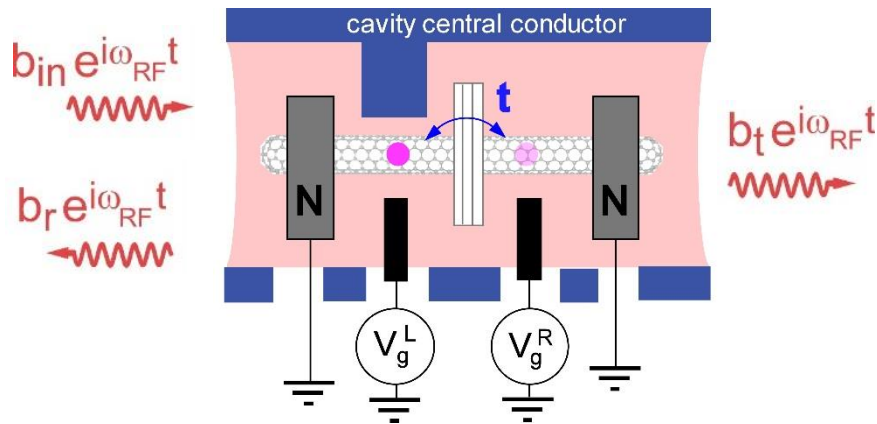


$$\bar{a} = \frac{a_0}{\omega_{RF} - \omega_0 - i\Lambda_0 - \chi(\omega_{RF})}$$

Second order in  $g$

# Cavity response in presence of the DQD

Bruhat, Cubaynes, Viennot, Dartiailh, Desjardins, Cottet, and Kontos, PRB 2018



semiclassical cavity transmission  
linear response treatment

$$\frac{b_t}{b_{in}} = (A_0 + \Delta A) e^{i(\varphi_0 + \Delta\varphi)}$$

$$= \frac{t_0}{\omega_{RF} - \omega_0 - i\Lambda_0 - \frac{\chi(\omega_{RF})}{2}}$$

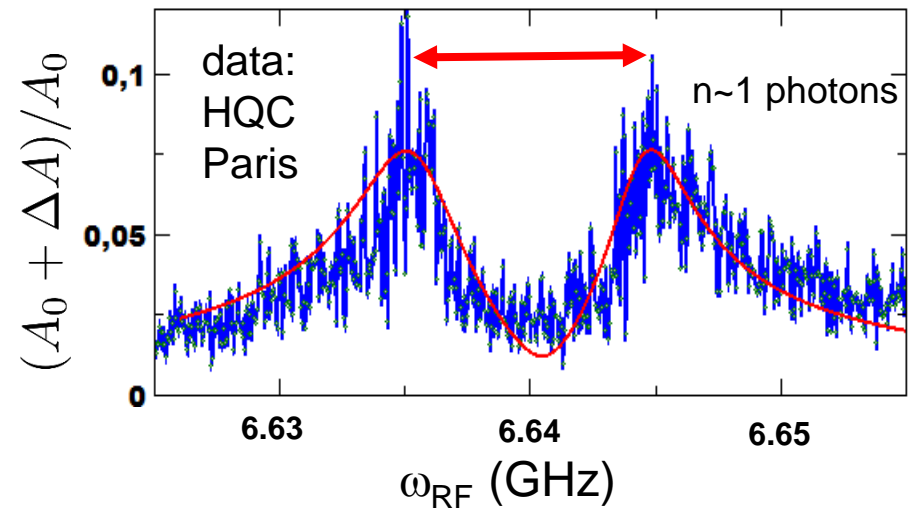
Nanocircuit charge susceptibility:

$$\frac{\chi(\omega_{RF})}{2} = \frac{(g_{\perp})^2}{\omega_{RF} - \omega_{+-} + i\Gamma_2}$$

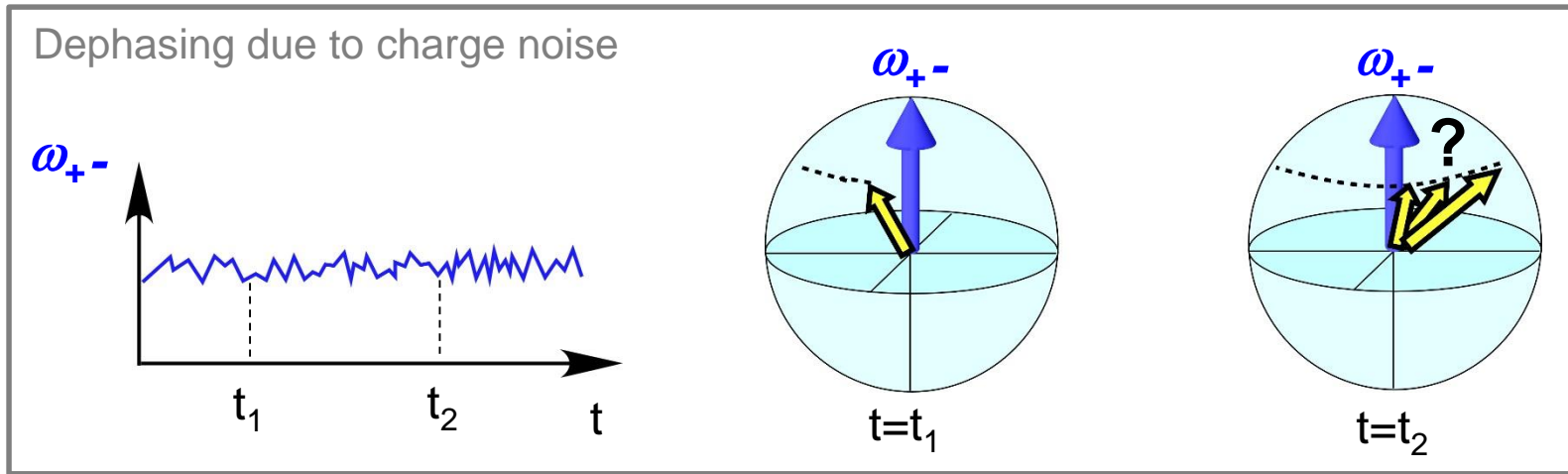
*transverse coupling crucial!*

$$g_{\perp} \gg \Gamma_2, \Lambda_0$$

$2\tilde{g}_{\perp} \sim 10$  MHz



# Reaching the strong coupling regime



Ideas:

$$g_{\perp} \gg \Gamma_s, \Lambda_0$$

increase  $g_{\perp}$

$$\Gamma_s \sim \frac{1}{\pi} \frac{\partial^2 \omega_{+-}}{\partial \epsilon^2} E_c^2 (\delta n)^2$$

decrease

$$E_c = \frac{e^2}{2C_{dot}}$$

decrease intrinsic  
charge noise

# Charge transition strongly coupled to cavity

	charge decoherence rate:	reduced charge/photon coupling:
● <b>Carbon nanotube:</b> <i>Bruhat et al., PRB 2018 (Paris)</i>	$\Gamma_2 = 5 \text{ MHz}$	$\tilde{g}_\perp / (\Lambda_0 + \Gamma_2) = 2.3$
● <b>2DEG: GaAs/AlGaAs</b> <i>Scarlino et al., PRL 2019 (Zurich)</i>	$\Gamma_2 = 3,3 \text{ MHz}$	$g_\perp / (\Lambda_0 + \Gamma_2) = 3.1$
● <b>2DEG Si/SiGe</b> <i>Mi et al., Science 2017 (Princeton)</i>	$\Gamma_2 = 6.7 \text{ MHz}$	$g_\perp / (\Lambda_0 + \Gamma_2) = 2.2$

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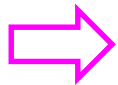


# Coupling spins to photons

**Readout of a single spin and distant spin/spin coupling  
by using circuit QED techniques?**

# Coupling spins to photons

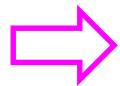
**Readout of a single spin and distant spin/spin coupling  
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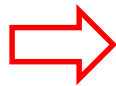
**Boost weak natural magnetic coupling** *Haikka et al. PRA 95 022306 (2017)*

# Coupling spins to photons

**Readout of a single spin and distant spin/spin coupling by using circuit QED techniques?**



**Boost weak natural magnetic coupling** *Haikka et al. PRA 95 022306 (2017)*



**Electric coupling?**



- Inhomogeneous nuclear fields: *Burkard et al., PRB (2006)*
- Spin-orbit interaction: *Trif et al., PRB (2008)*

# Coupling spins to photons

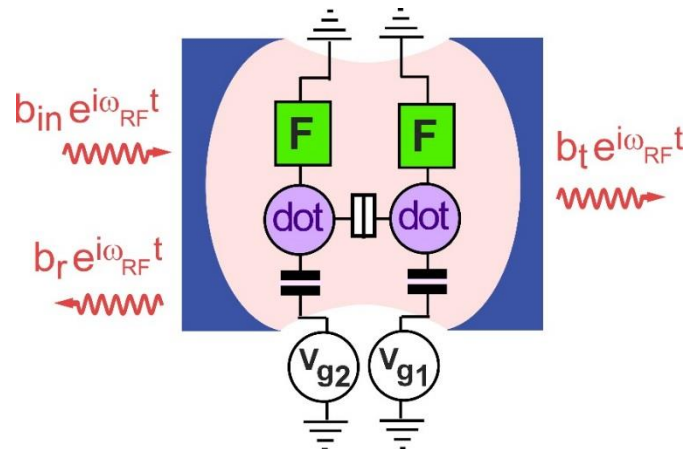
Readout of a single spin and distant spin/spin coupling by using circuit QED techniques?

➡ Boost weak natural magnetic coupling *Haikka et al. PRA 95 022306 (2017)*

➡ Electric coupling?

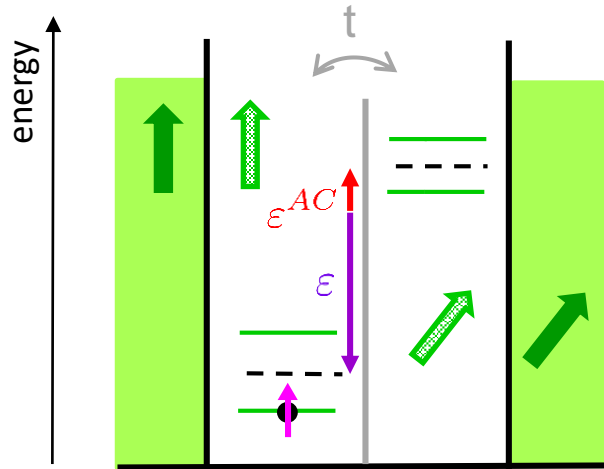


- Inhomogeneous nuclear fields: *Burkard et al., PRB (2006)*
- Spin-orbit interaction: *Trif et al., PRB (2008)*
- Artificial spin-orbit coupling from ferromagnetic contacts: *Cottet/Kontos, PRL (2010)*

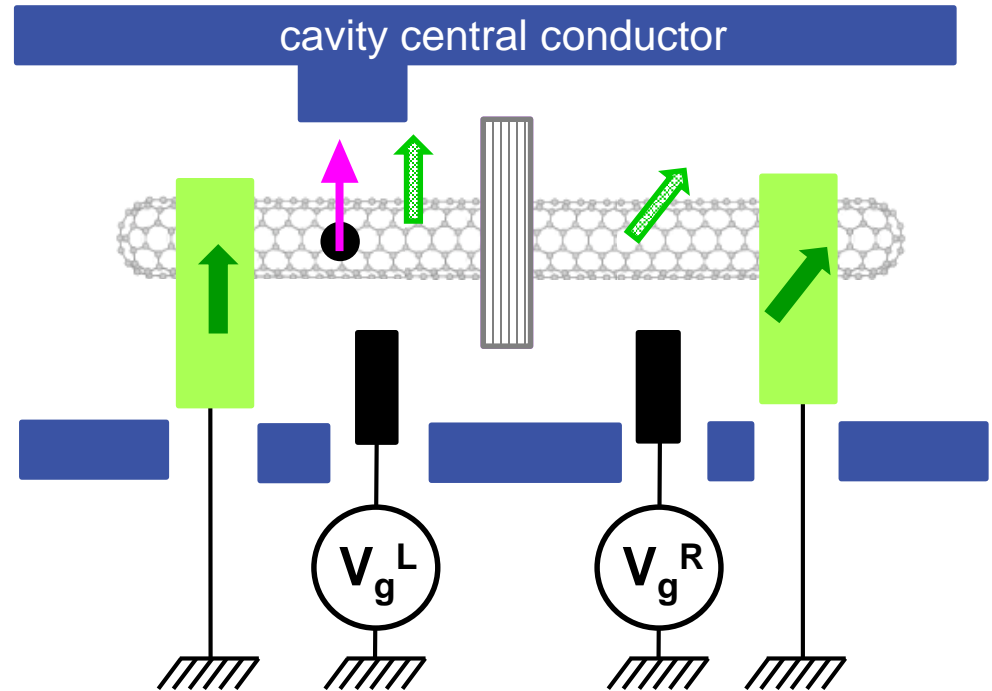


# DQD with non-collinear ferromagnetic contacts

Theory proposal: Cottet & Kontos, PRL **105**, 160502 (2010)



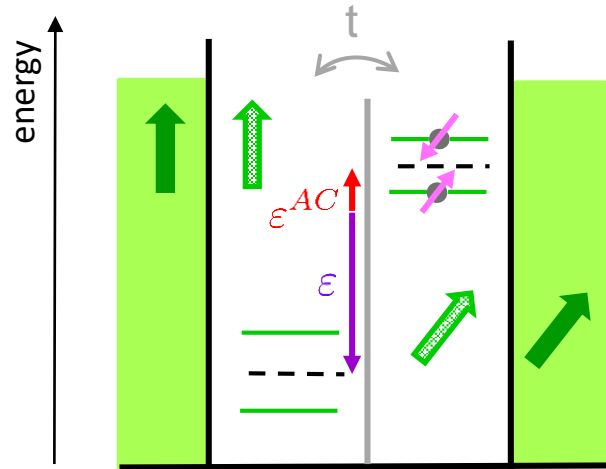
$$\varepsilon^{AC} = \hbar(g_L - g_R)(\hat{a} + \hat{a}^\dagger)$$



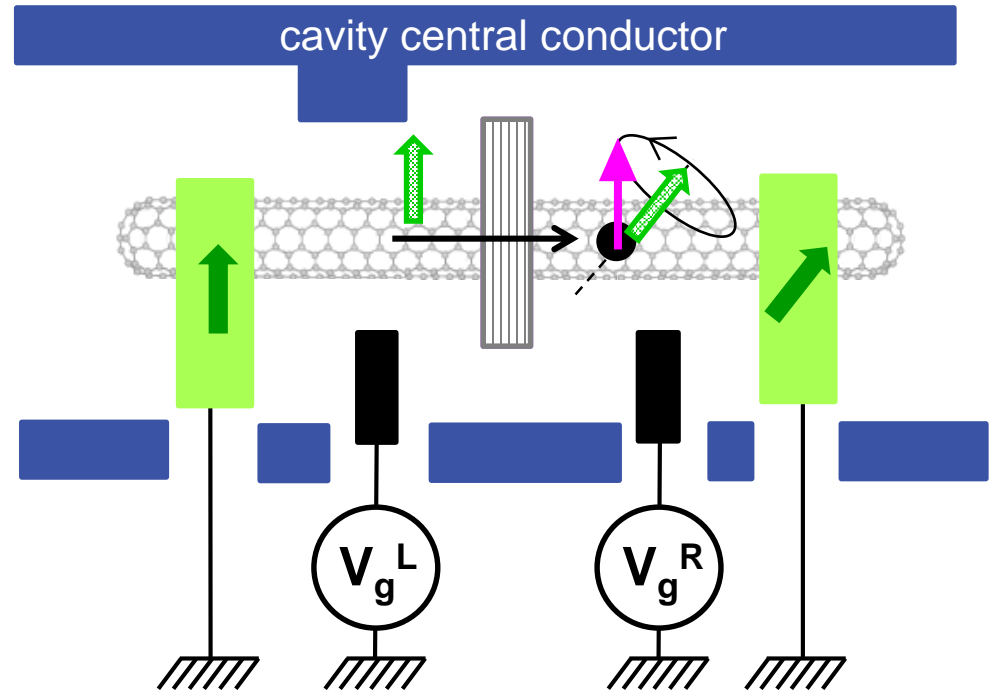
- Coupled orbital and spin variables  $\implies$  « artificial spin-orbit interaction »
- Photons couple to orbital part and therefore to spin  $\hat{h}_{int}(\hat{a} + \hat{a}^\dagger) \simeq \hbar g(\hat{a}^\dagger \hat{c}_\uparrow^\dagger \hat{c}_\downarrow + \hat{a} \hat{c}_\downarrow^\dagger \hat{c}_\uparrow)$

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$$\varepsilon^{AC} = \hbar(g_L - g_R)(\hat{a} + \hat{a}^\dagger)$$



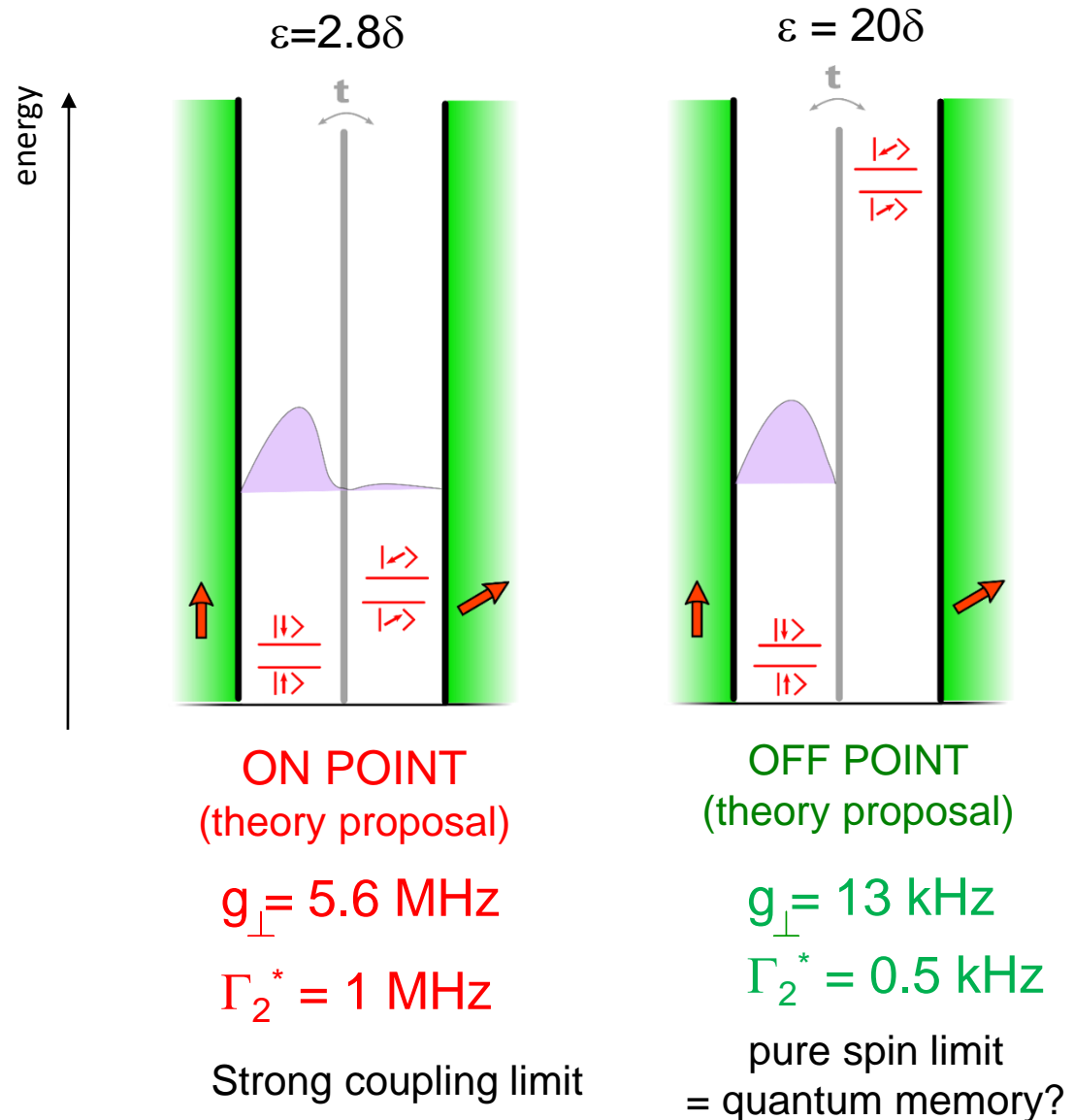
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# Optimization of the spin/photon coupling

Theory proposal:  
Cottet & Kontos,  
*PRL* **105**, 160502 (2010)

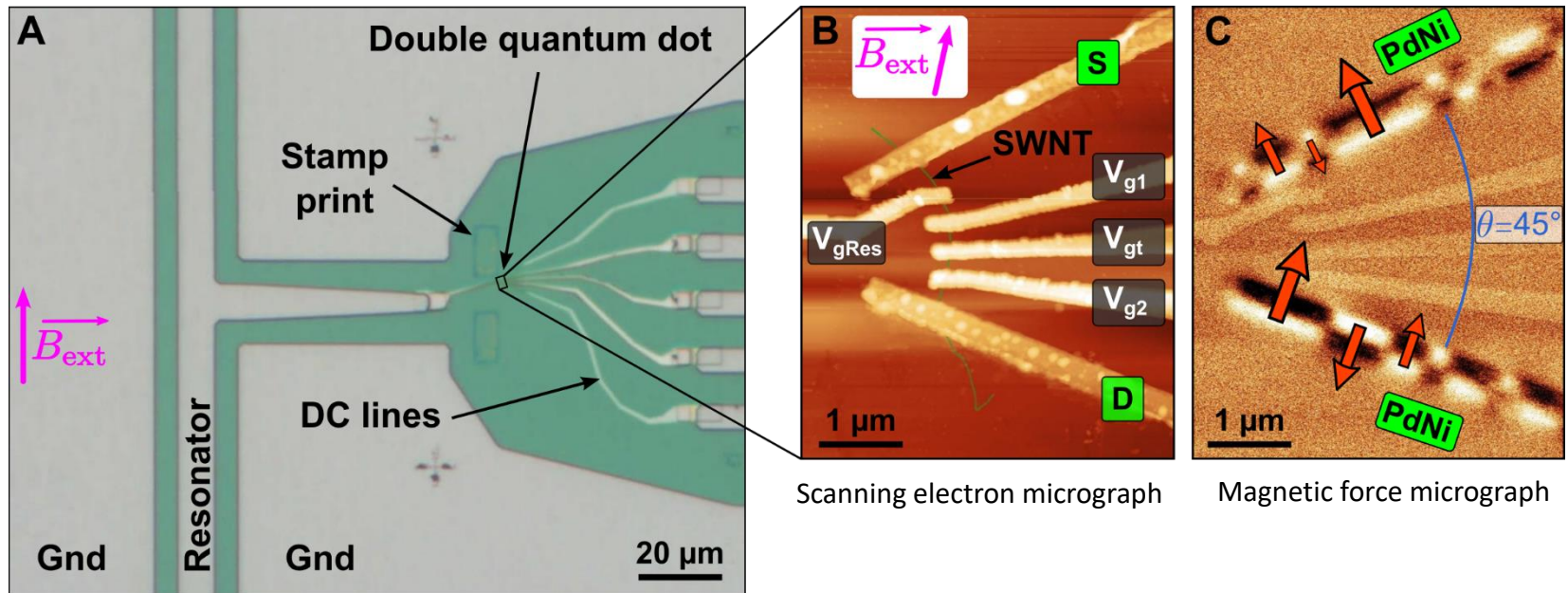
## Main sources of decoherence:

- charge noise => dephasing
- phonons => relaxation



# Experiment: carbon nanotube and PdNi contacts

Viennot, Dartailh, Cottet, & Kontos, *Science* **349**, 6246 (2015)

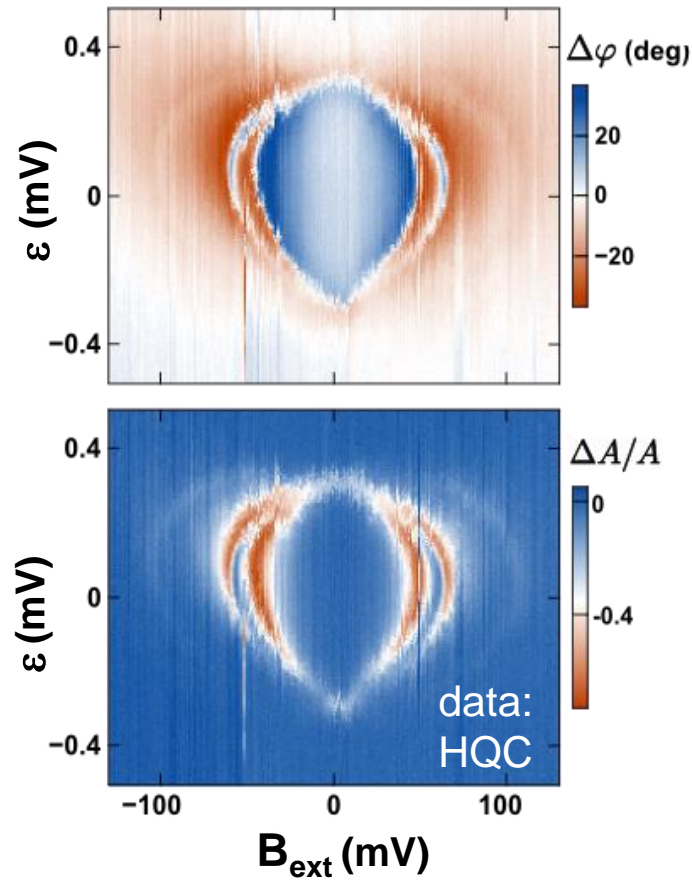


- $V_{g1}, V_{g2}$  are DC electrostatic gates controlling  $\epsilon^{DC}$
- Selective coupling of the right dot to the cavity by the gate  $V_{gRes}$
- Non-collinear magnetizations imposed by contacts shape
- Resonator  $\omega_0/2\pi = 6.72\text{GHz}$ , quality factor  $Q \approx 10^4$  up to  $B_{ext} = 100\text{mT}$



# Magneto-spectroscopy

*Viennot, Dartailh, Cottet, & Kontos, Science 349, 6246 (2015)*

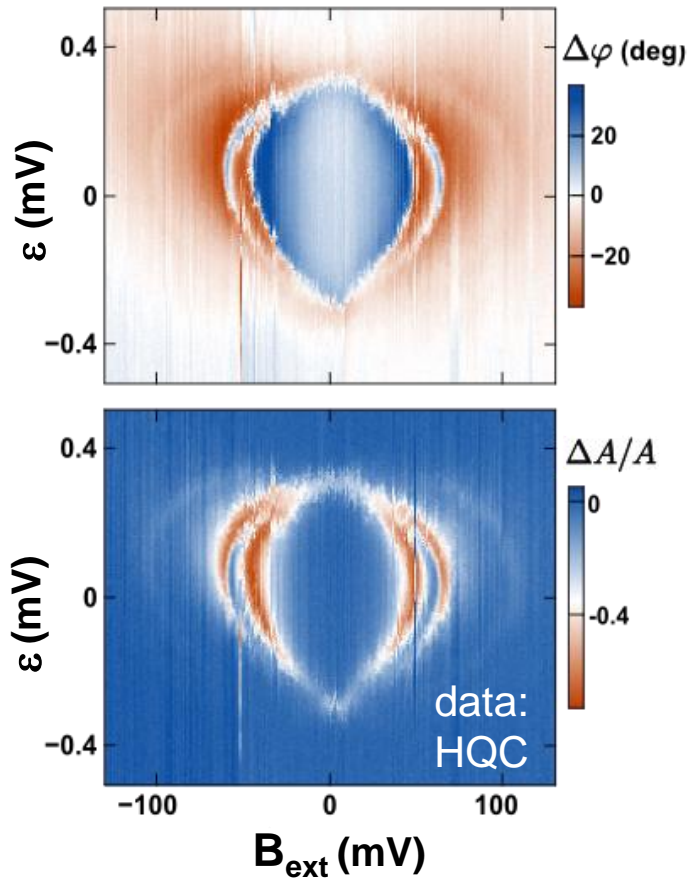


These resonances strongly  
move with  $B_{\text{ext}}$  !

⇒ Spin degree of freedom involved

# Magneto-spectroscopy

Viennot, Dartiailh, Cottet, & Kontos, *Science* **349**, 6246 (2015)



Cavity transmission at  $\omega_{RF} = \omega_0$

$$\begin{aligned} \frac{b_t}{b_{in}} &= (A_0 + \Delta A) e^{i(\varphi_0 + \Delta\varphi)} \\ &= \frac{t_0}{i\Lambda_0 - \chi(\omega_0)} \end{aligned}$$

Nanocircuit charge susceptibility:

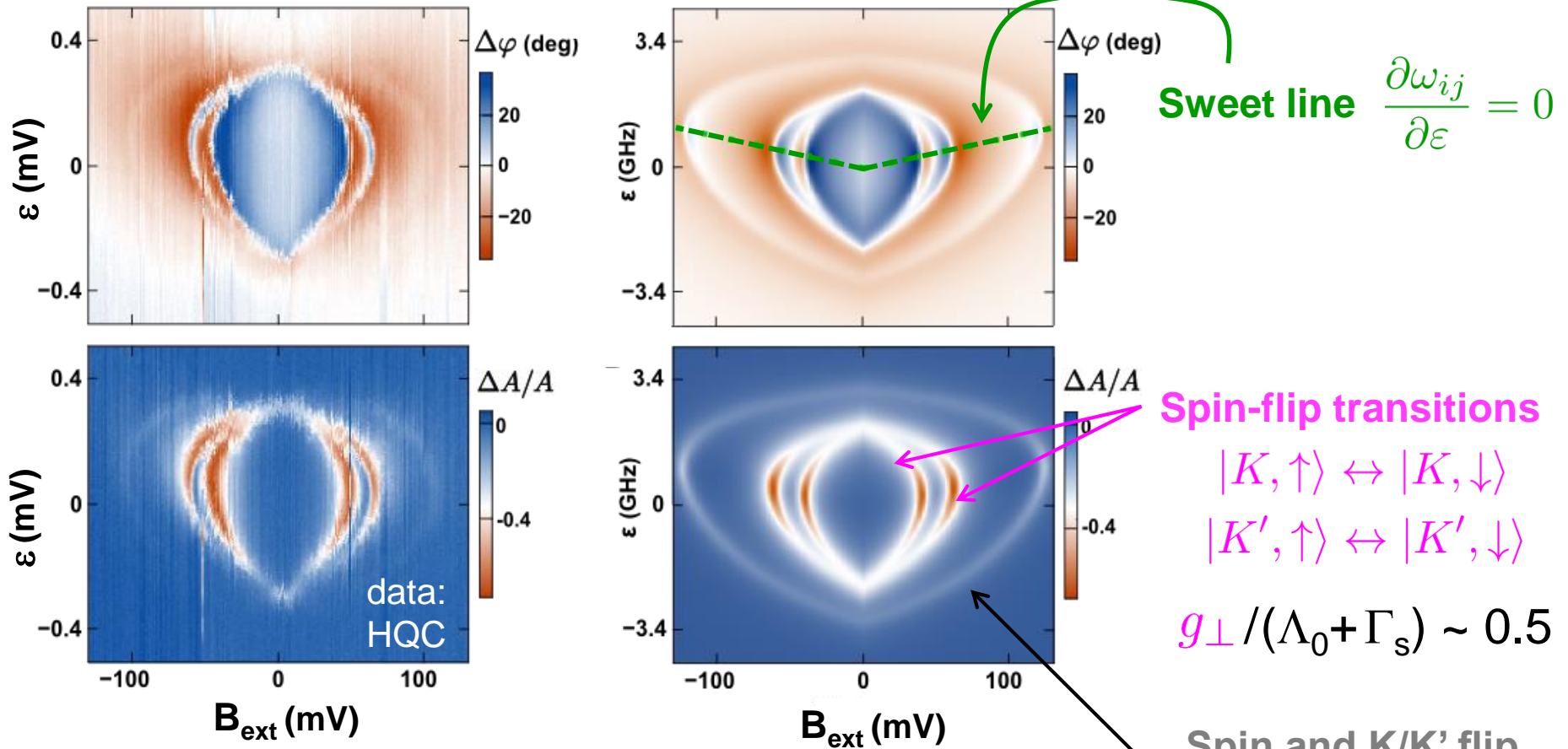
$$\chi_2(\omega_0) = \sum_{ij} \frac{g_{ij}^2}{\omega_0 - \omega_{ij} + i\Gamma_{ij}}$$

Multiple transitions due to  
L/R, spin and K/K' degrees of freedom

The cavity provides a cut of the DQD spectrum at frequency  $\omega_0$

# Magneto-spectroscopy

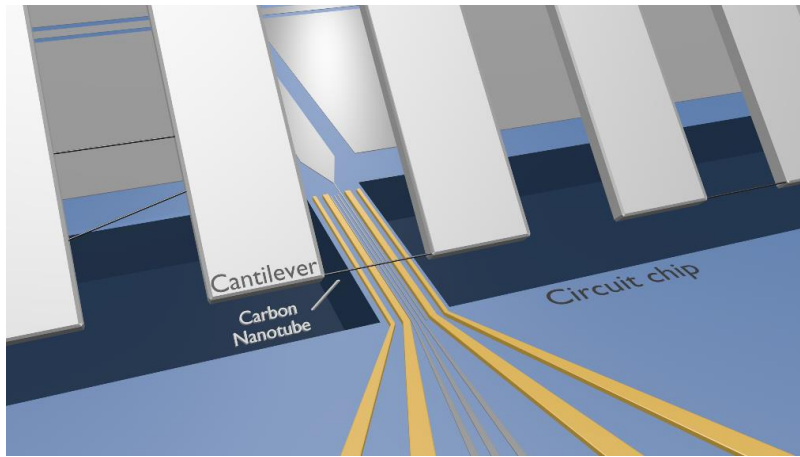
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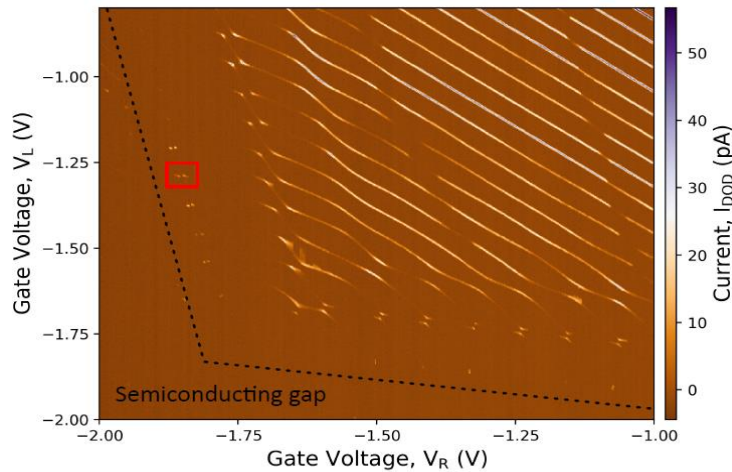
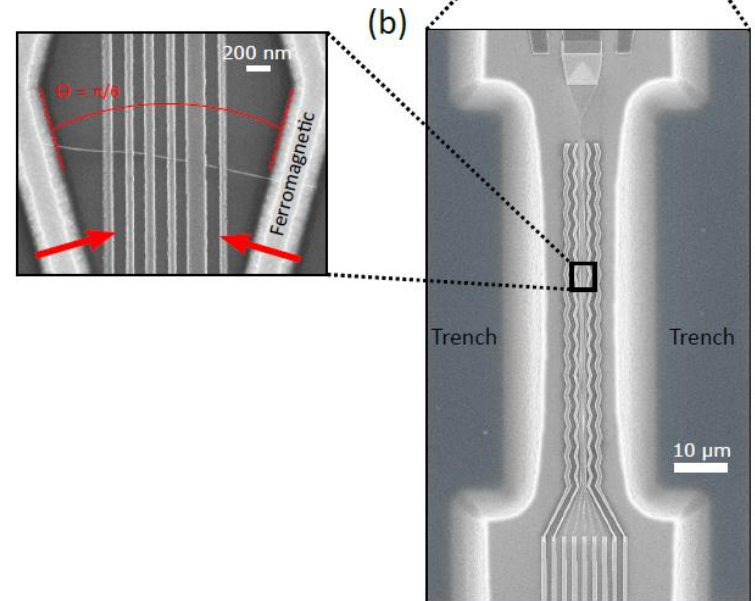
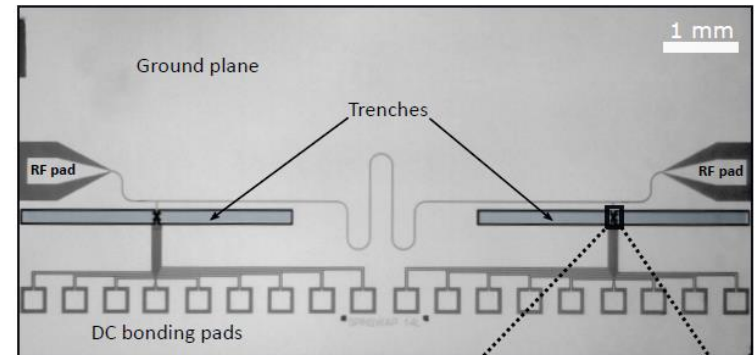
*Charge noise gives an important contribution to decoherence!*

# Towards the strong spin/photon coupling with carbon nanotubes

*T. Cubaynes et al., NPJQI 2019*

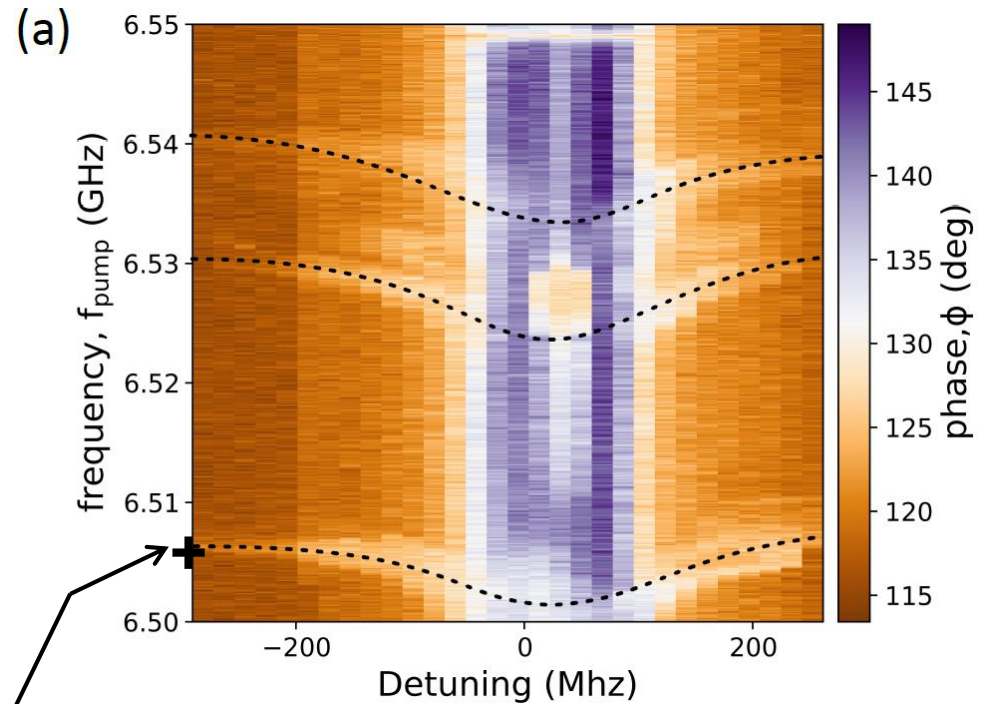
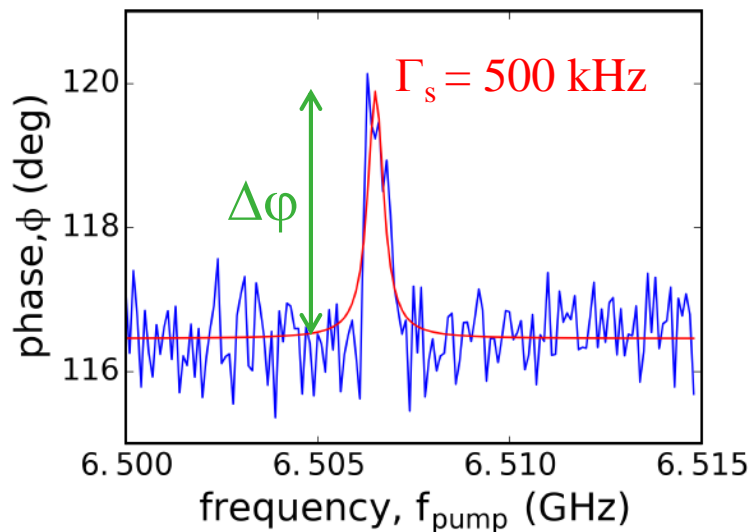
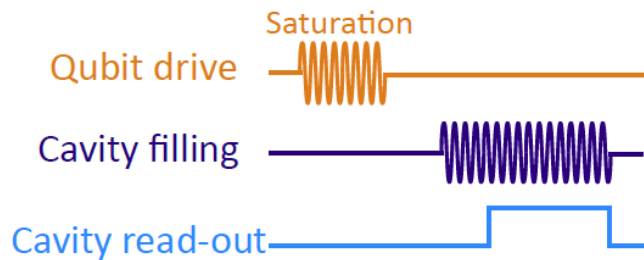


*Stapling technique, see also Pei et al. Nature Nano'12, Weissman et al. Nature Nano'13, Ranjan et al. Nature Comm. '15*



# Reaching the strong spin/photon coupling with carbon nanotubes

T. Cubaynes et al.  
NPJQI 2019



$$\Delta\phi \sim \frac{g_{\perp}^2}{\Lambda_0(\omega_s - \omega_0)}$$

$$\Rightarrow g_{\perp} / (\Lambda_0 + \Gamma_s) = 2$$

# Strong spin/photon coupling

- **Carbon nanotube:**

*Viennot, et al. Science (2015) (Paris)*

$$\Gamma_s = 2.5 \text{ MHz} \quad g_{\perp} / (\Lambda_0 + \Gamma_s) = 0.46$$

- **Suspended carbon nanotube:**

*Cubaynes, et al. NPJQI (2019) (Paris)*

$$\Gamma_s = 0.25 \text{ MHz} \quad g_{\perp} / (\Lambda_0 + \Gamma_s) = 2$$

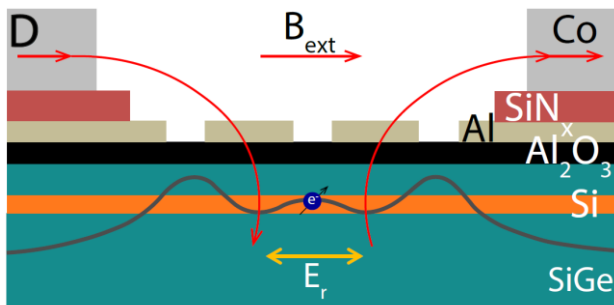
- **2DEG Si/SiGe: anticrossing observed**

*Samkharadze et al., Science (2018) (Delft)*

$$\Gamma_s = 1.8 \text{ MHz} \quad g_{\perp} / (\Lambda_0 + \Gamma_s) = 2.6$$

*Mi et al. Nature (2018) (Princeton)*

$$\Gamma_s = 2.4 \text{ MHz} \quad g_{\perp} / (\Lambda_0 + \Gamma_s) = 1.7$$



*See also: Landig et al.,  
Nature (2018) (Zurich)  
[different principle]*

# Strong spin/photon coupling

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$$\Gamma_s = 2.4 \text{ MHz} \quad g_{\perp} / (\Lambda_0 + \Gamma_s) = 1.7$$

## Possible improvements:

⇒ A better control of spin/charge hybridization

⇒ Use nanotubes with pure C12


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# Keldysh theory for quantum dot circuits

Bruhat, Viennot, Dartailh, Desjardins, Kontos and Cottet, *Phys. Rev. X*, **6**, 021014 (2016)

$$\frac{b_t}{b_{in}} = \frac{t_0}{\omega_{RF} - \omega_0 - i\Lambda_0 - \chi_2(\omega_{RF})}$$

semiclassical cavity transmission  
linear response treatment

electron/photon  
coupling matrix

Nanocircuit charge susceptibility (non interacting multi-dot case) :

$$\chi_2 = -\frac{i}{2} \int_{\omega} \text{Tr} \left[ \tilde{G}_K(\omega) \tilde{g} \left( \tilde{G}_a(\omega - \omega_0) + \tilde{G}_r(\omega + \omega_0) \right) \tilde{g} \right]$$

Keldysh Green's functions:

$$G_r^{j,j'}(t, t') = -i\theta(t) \left\langle \{ \hat{c}_j(t), \hat{c}_{j'}^\dagger(t') \} \right\rangle$$

$$G_a^{j,j'}(t, t') = i\theta(-t) \left\langle \{ \hat{c}_j(t), \hat{c}_{j'}^\dagger(t') \} \right\rangle$$

$$G_K^{j,j'}(t, t') = -i \left\langle [ \hat{c}_j(t), \hat{c}_{j'}^\dagger(t') ] \right\rangle$$

j,j': dot, orbital, and spin indices

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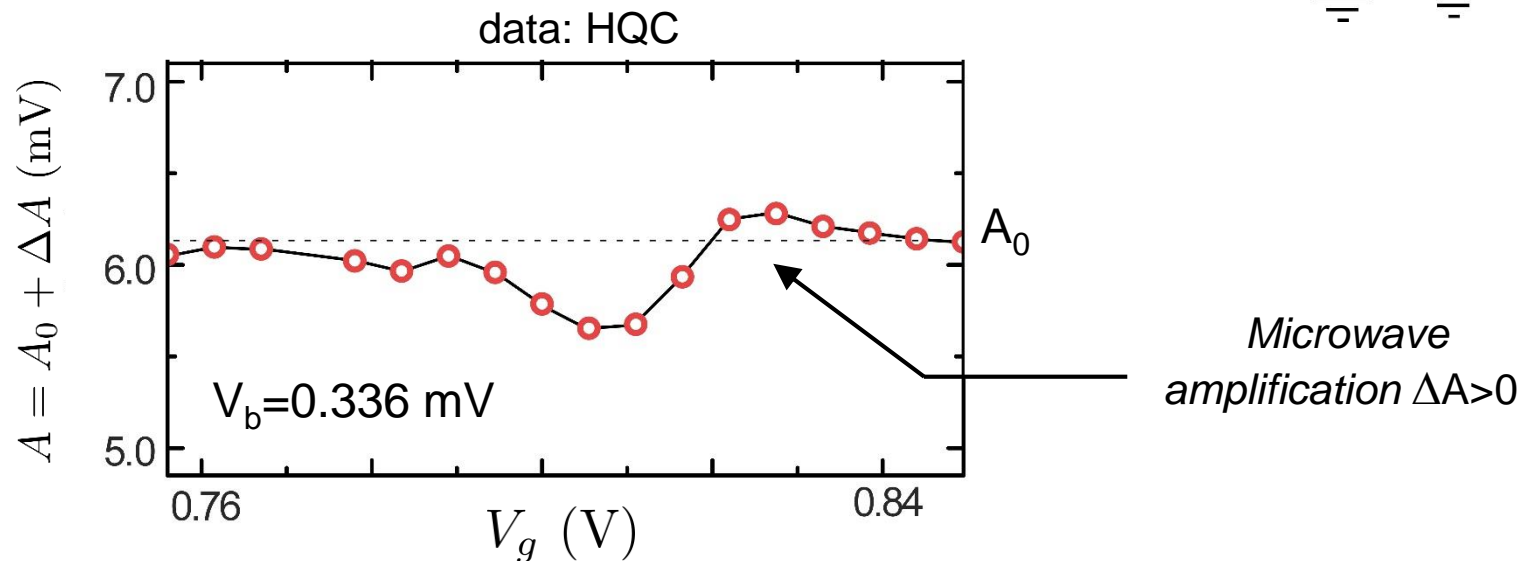
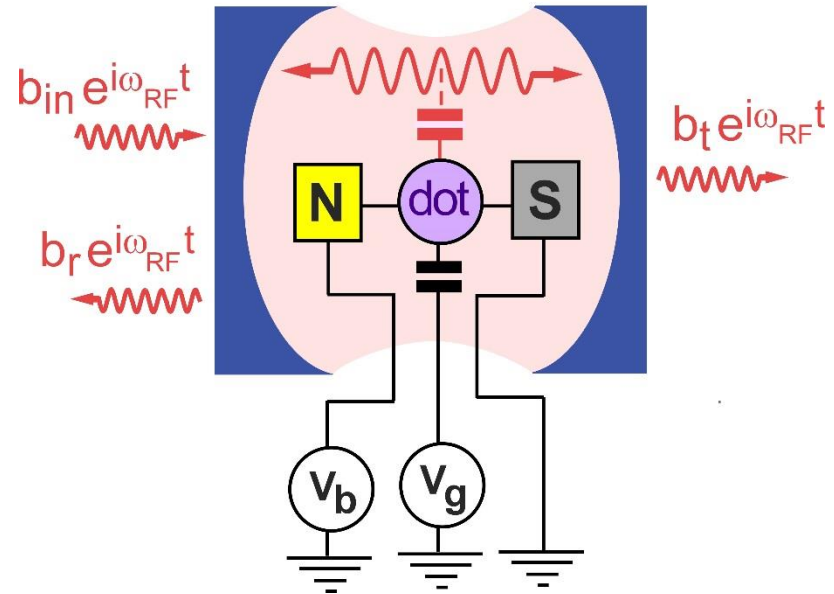
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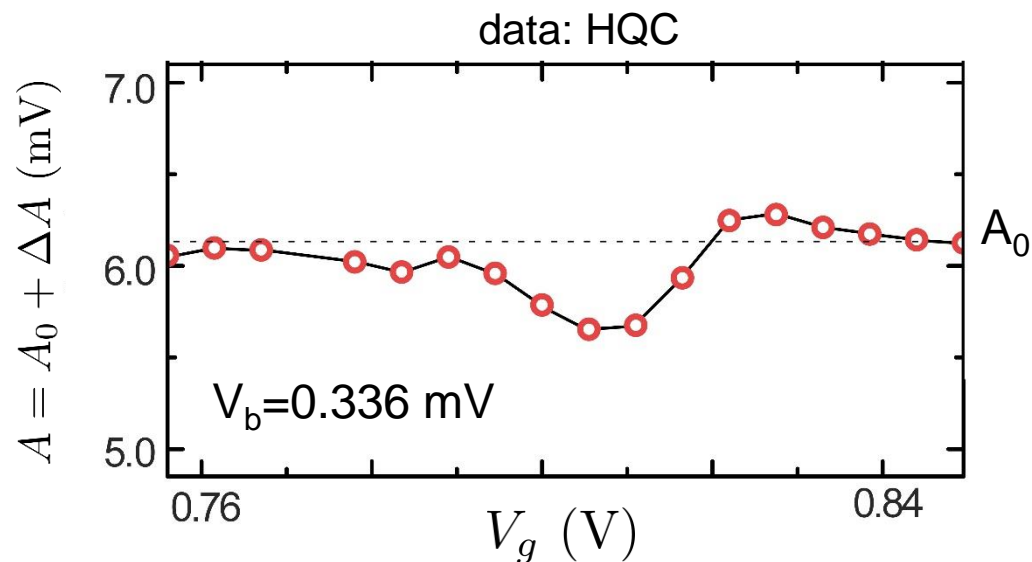
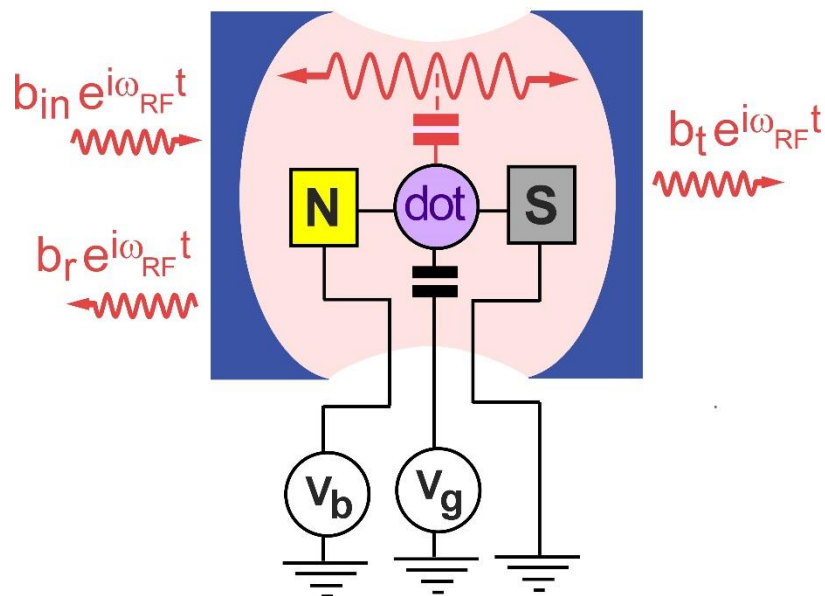
# S/dot/N bijunction at finite $V_b$

*Bruhat, Viennot, Dartiailh, Desjardins, Kontos & Cottet, Phys. Rev. X 6, 021014 (2016)*



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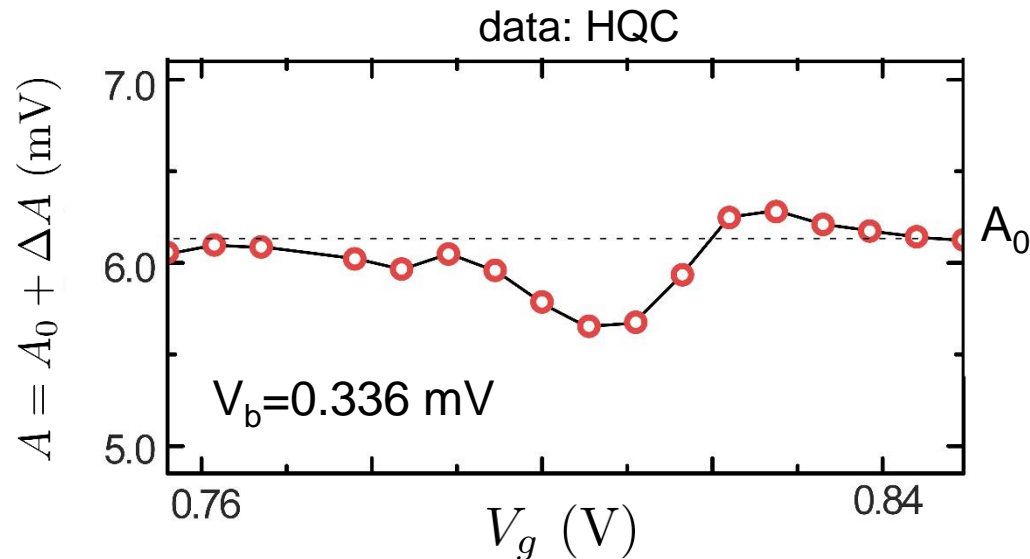
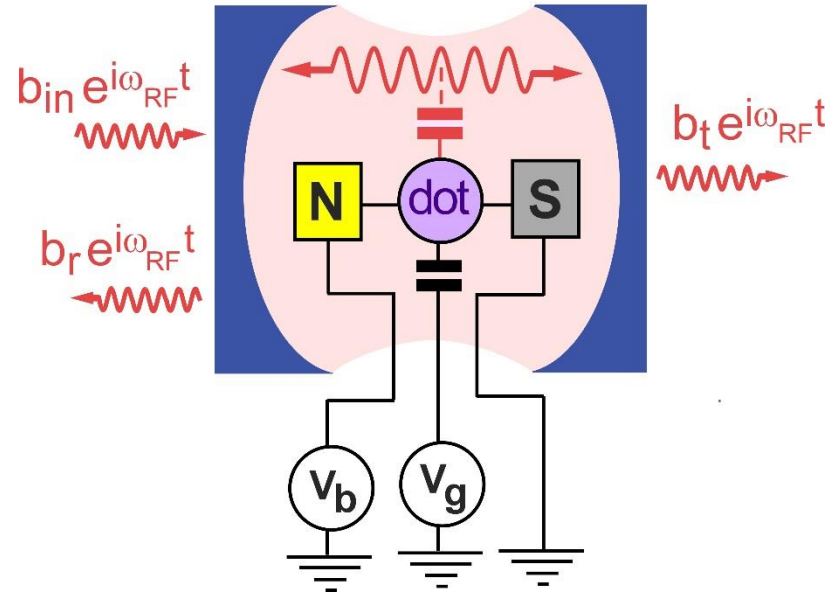
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*No relevant dot internal degree of freedom to explain photon-emission*

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*No relevant dot internal degree of freedom to explain photon-emission*



Look at theory?

# Keldysh theory for quantum dot circuits

*Bruhat, Viennot, Dartailh, Desjardins, Kontos and Cottet, Phys. Rev. X, 6, 021014 (2016)*

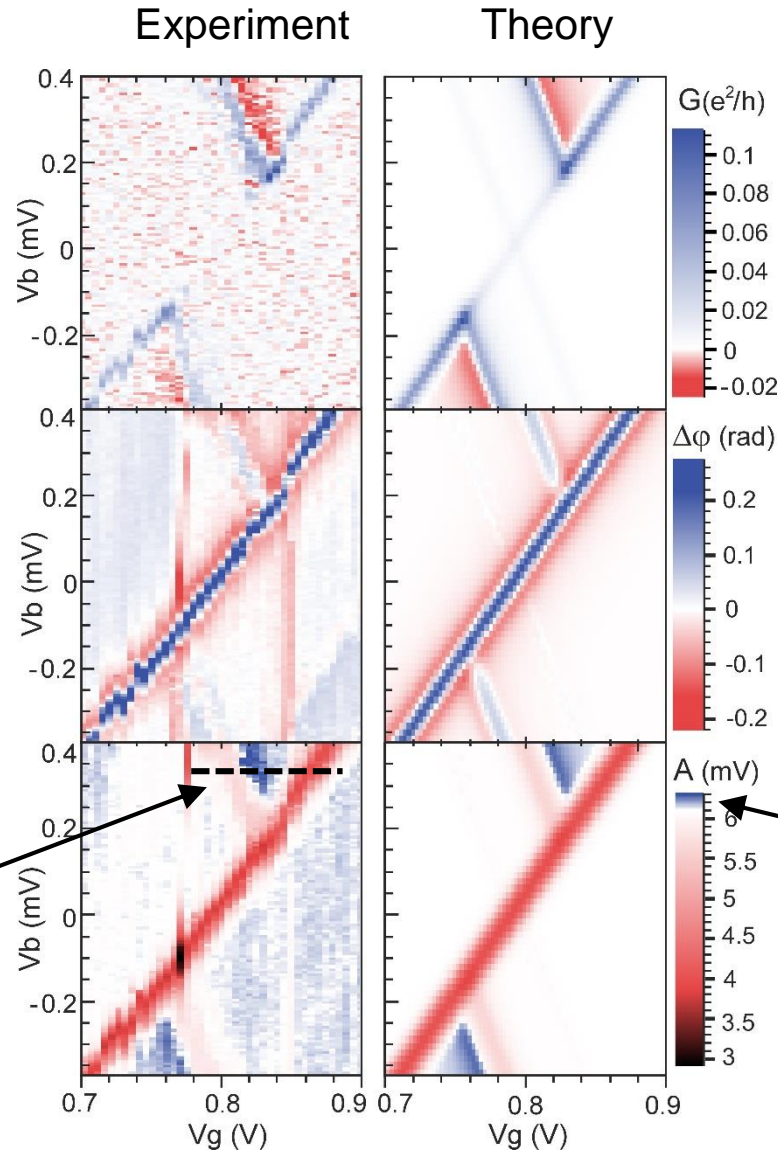
$$\frac{b_t}{b_{in}} = \frac{t_0}{-i\Lambda_0 - g^2 \chi_2(\omega_0)}$$

$$\chi_2 = -\frac{i}{2} \int_{\omega} \text{Tr}_d \left[ \tilde{G}_K(\omega) \tilde{g} \left( \tilde{G}_a(\omega - \omega_0) + \tilde{G}_r(\omega + \omega_0) \right) \tilde{g} \right]$$

Fitting parameters:

- gap of S:  $\Delta$
- temperature:  $T$
- dot/cavity coupling:  $g$
- dot/N tunnel rate:  $\Gamma_N$
- dot/S tunnel rate:  $\Gamma_S$
- BCS peaks broadening:  $\Gamma_b$

# Test of theory at finite bias voltages



Quantitative agreement with the data

CUT

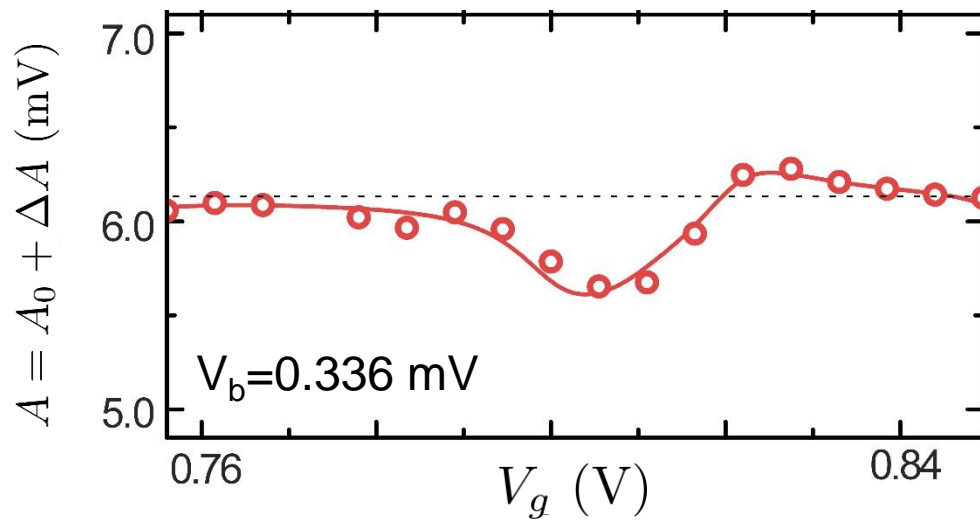
Microwave amplification  $\Delta A > 0$

# Inelastic tunneling between dot and S

*Bruhat et al., Phys. Rev. X (2016)*

$\Gamma_b/2\pi$ :

— 8 GHz



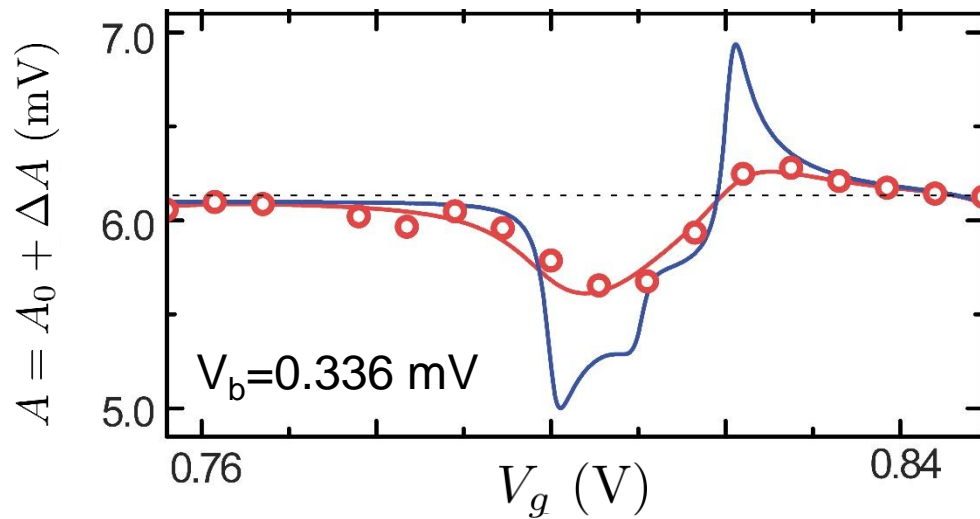


# Inelastic tunneling between dot and S

*Bruhat et al., Phys. Rev. X (2016)*

$\Gamma_b/2\pi$ :

— 1 GHz  
— 8 GHz

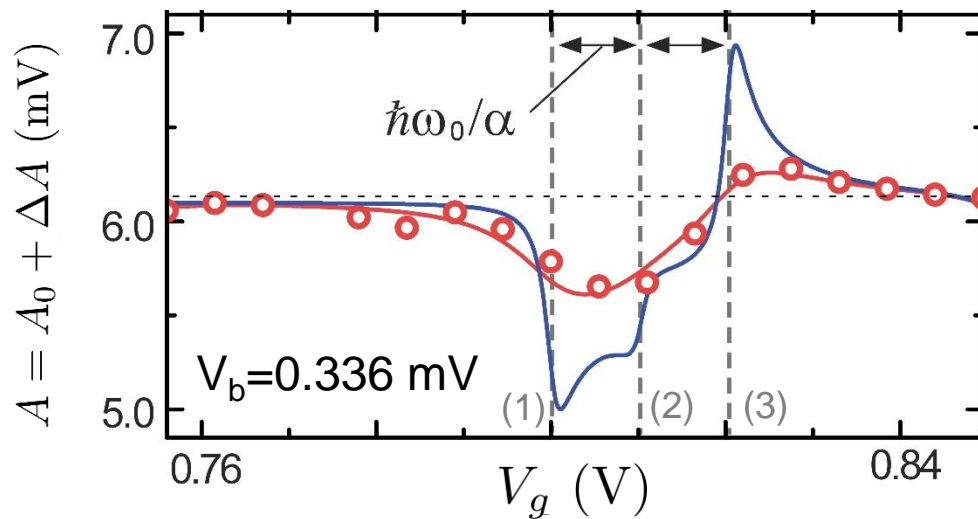


# Inelastic tunneling between dot and S

$\omega_0$  : cavity frequency  
 $\alpha$  : DC-gate lever arm

$\Gamma_b/2\pi$ :

— 1 GHz  
— 8 GHz

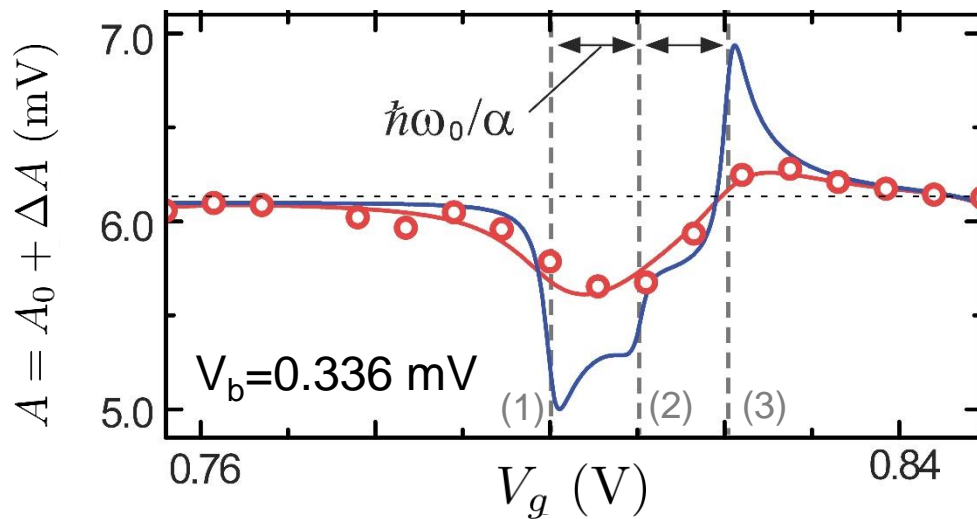


# Inelastic tunneling between dot and S

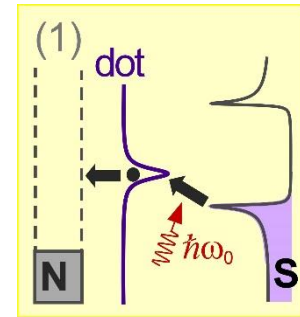
$\omega_0$  : cavity frequency  
 $\alpha$  : DC-gate lever arm

$\Gamma_b/2\pi$ :

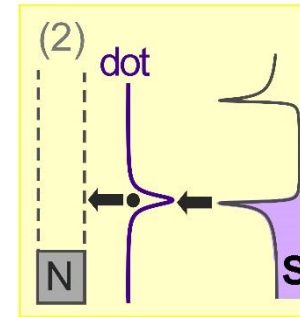
— 1 GHz  
 — 8 GHz



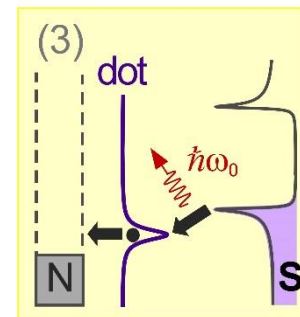
Photon-emission  
 visible only through cavity!



*Photon-absorption*



*Elastic transport*



*Photon-emission*

rate  $\sim 2$  MHz  
 $\sim 0.3$  pA

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- Strong charge/photon coupling
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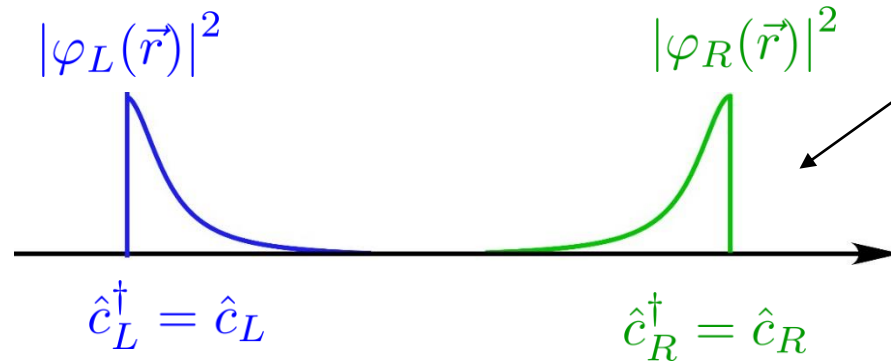
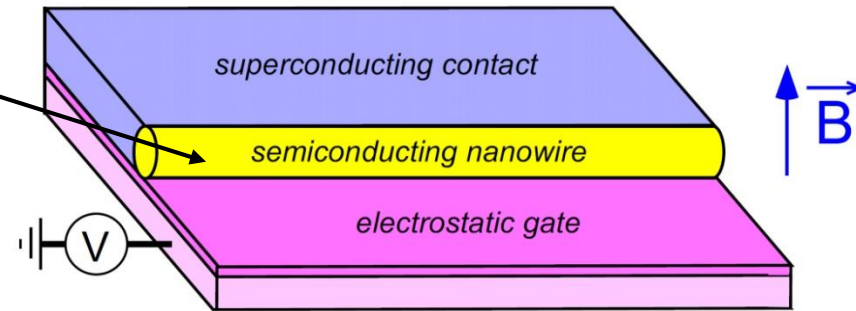
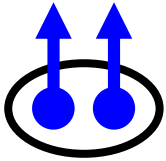
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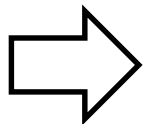
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Synthetic topological superconductivity



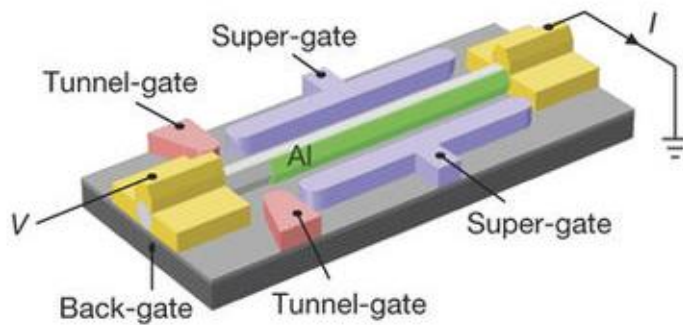
low-energy self-adjoint quasiparticle state



Possible application = **topologically protected quantum computation**

# Observation of zero bias conductance peaks in semiconducting nanowires

Zhang et al., Nature (2018)



See also:

Mourik, et al., Science 336, 1003 (2012)

Williams, et al., PRL (2012)

Das, et al., Nature Phys. (2012).

Deng, et al., Nano Lett. (2012),

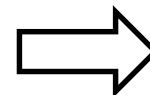
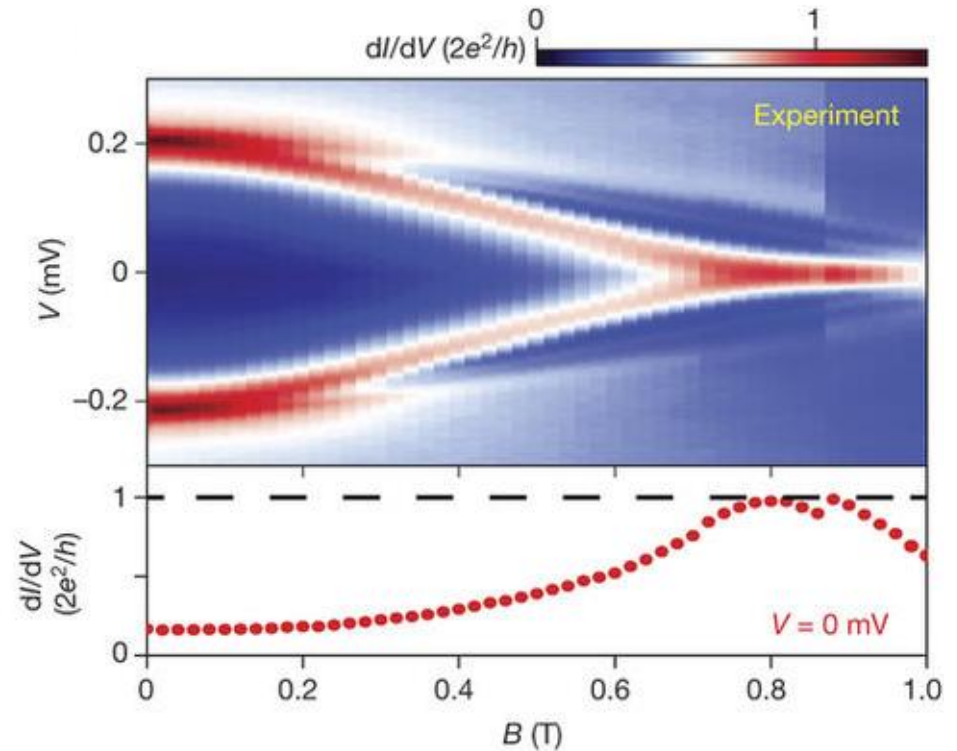
Rokhinson et al., Nature Phys. (2012)

Albrecht et al., Nature (2016)

Deng et al, Science (2016)

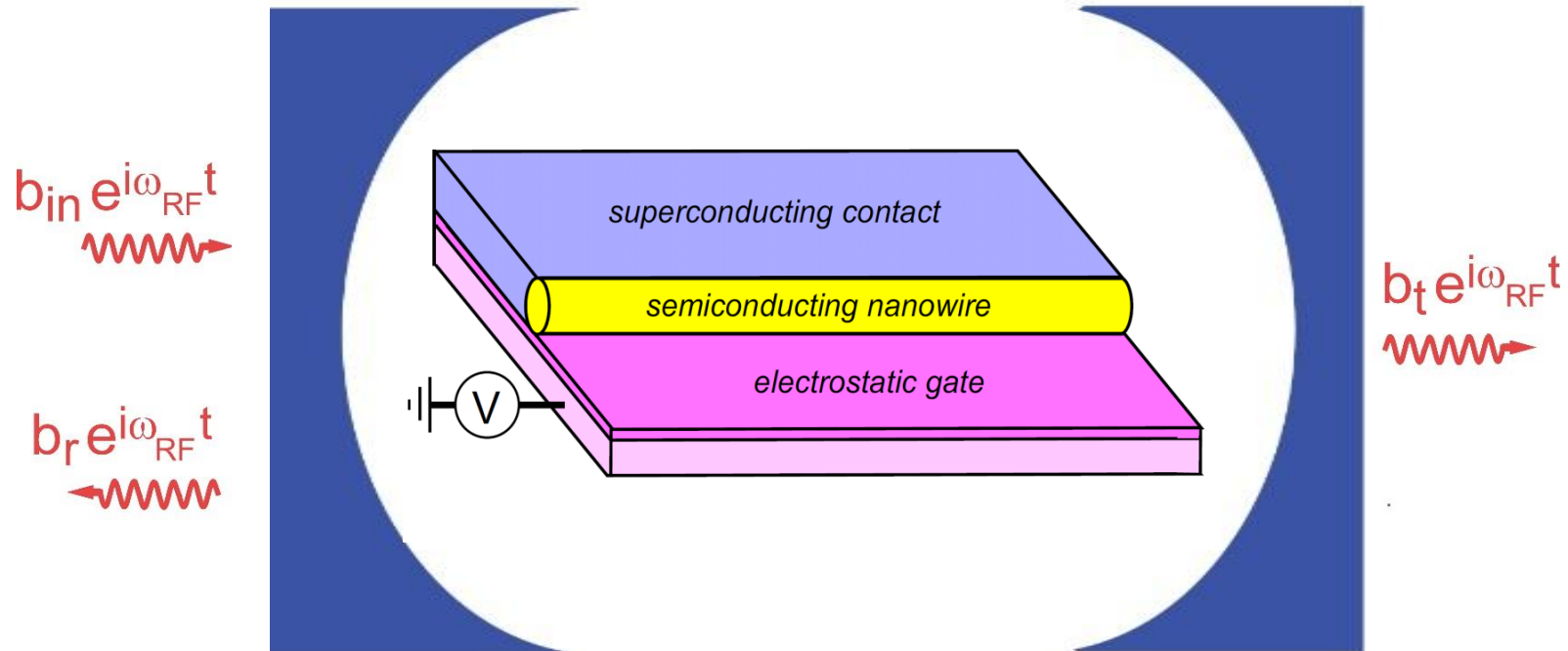
etc...

Density of states measurement:



Majorana bound state

$$\hat{c}^\dagger = \hat{c} ?$$

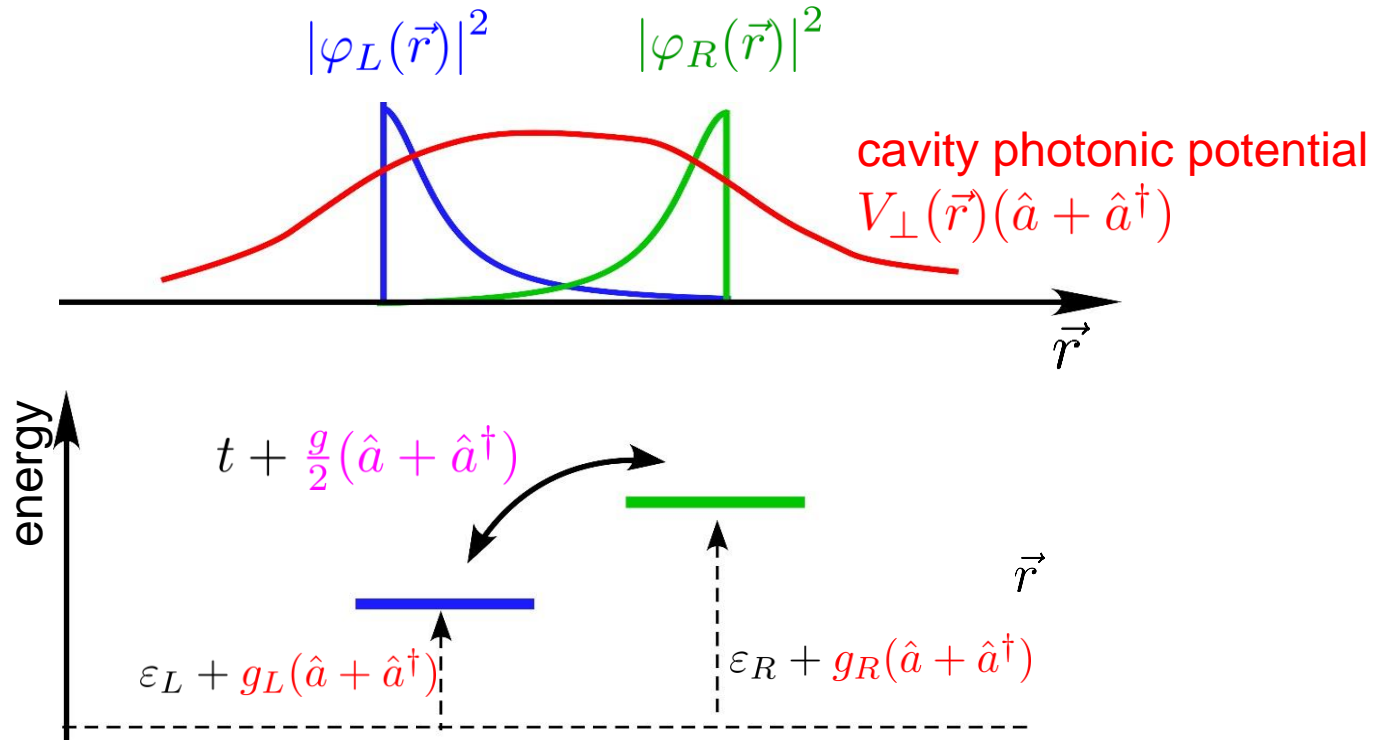


**What can we learn from a Majorana nanocircuit with a cavity?**

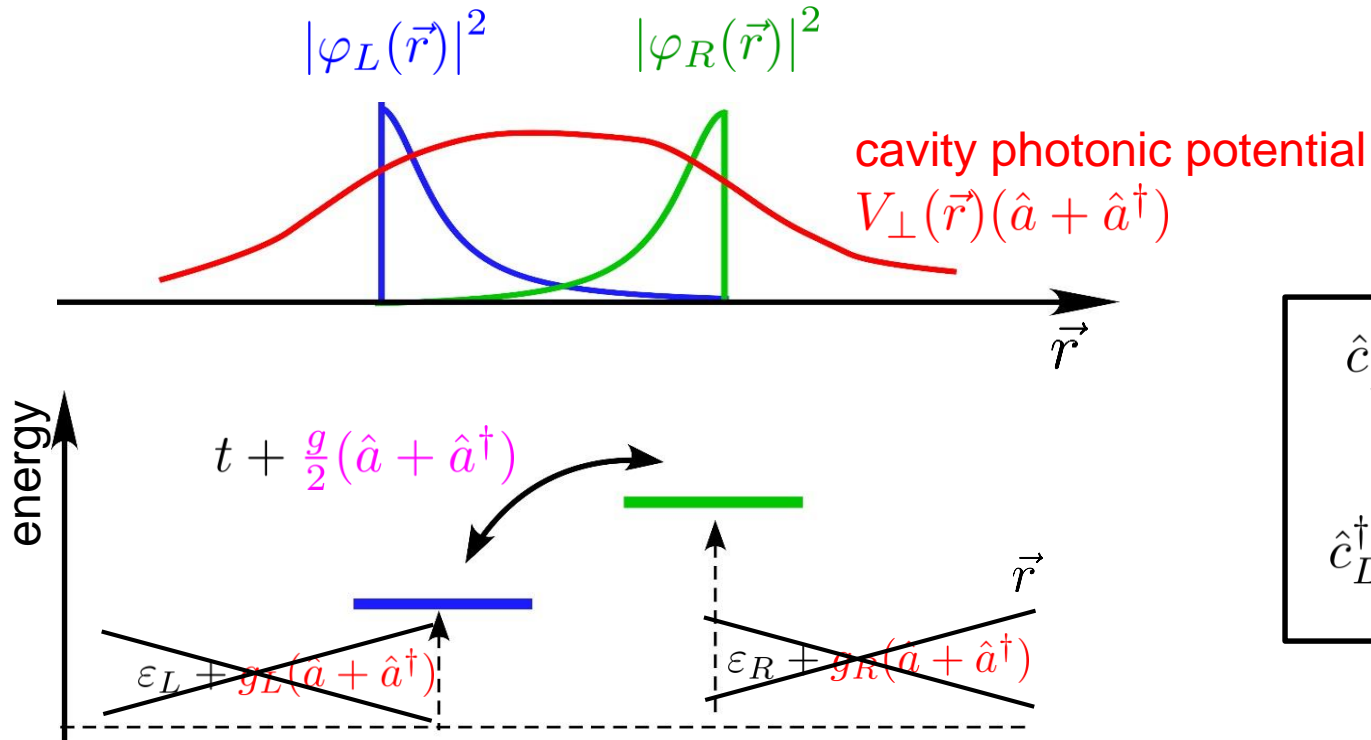
Majorana + cavities, see also:

Hassler et al. *New J. Phys.* (2011), Hyart et al., *PRB* (2013), Müller et al. *PRB* (2013), Ginossar and Grosfeld, *arXiv* 1307.1159, Trif and Tserkovnyak, *PRL* (2012), Schmidt et al. *PRL* (2013). Dmytruk et al. *PRB* (2015), Trif et al. *PRL* (2019)

# Majorana pair in a cavity





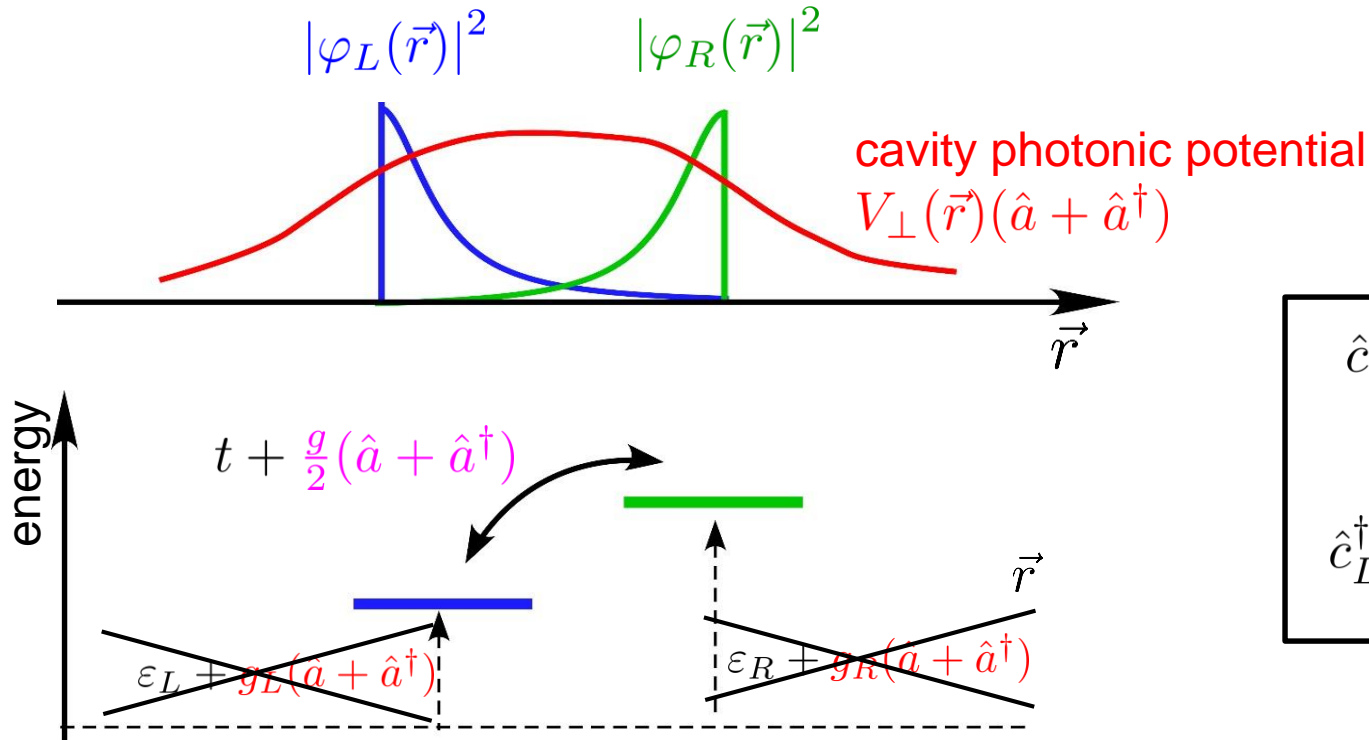


$$\hat{c}_{L(R)}^{\dagger} = \hat{c}_{L(R)}$$

➔

$$\hat{c}_L^{\dagger} \hat{c}_L = \hat{c}_R^{\dagger} \hat{c}_R = 1$$

# Majorana pair in a cavity

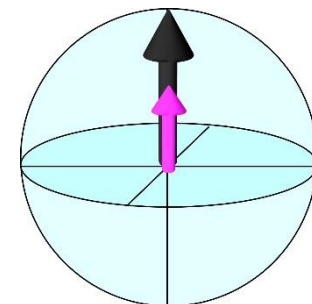


$$\hat{c}_{L(R)}^{\dagger} = \hat{c}_{L(R)}$$

➔

$$\hat{c}_L^{\dagger} \hat{c}_L = \hat{c}_R^{\dagger} \hat{c}_R = 1$$

$$\hat{H} = 2it\hat{c}_L\hat{c}_R + ig(\hat{a} + \hat{a}^{\dagger})\hat{c}_L\hat{c}_R + \hbar\omega_0\hat{a}^{\dagger}\hat{a}$$



no cavity signal?

*Dartiailh, Kontos, Douçot & Cottet, PRL 118, 126803 (2017)*

Theory

$$\frac{b_t}{b_{in}} = \frac{t_0}{\omega_{RF} - \omega_0 - i\Lambda_0 - \chi(\omega_{RF})}$$

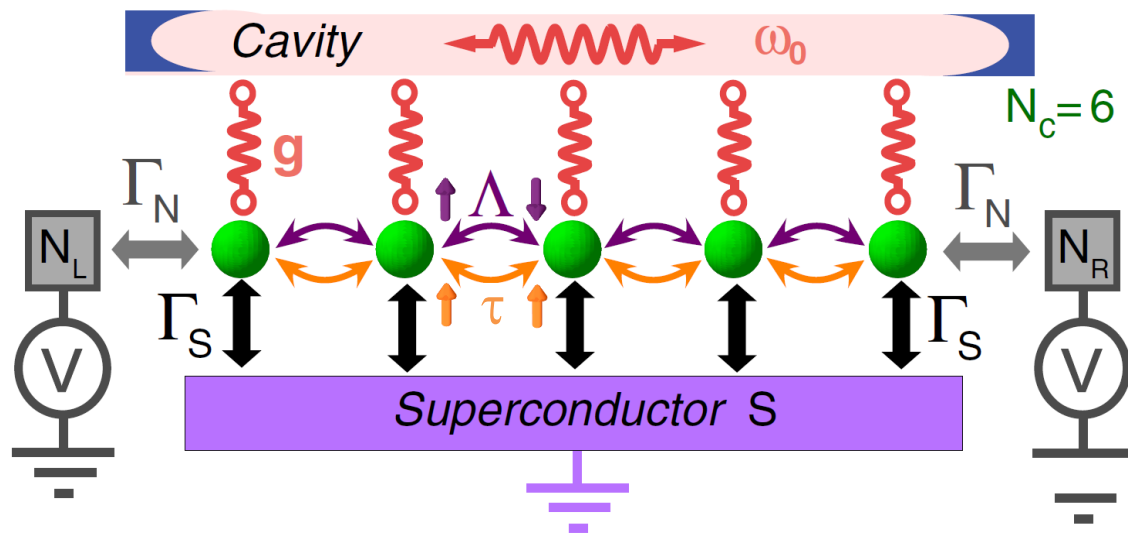
semiclassical cavity transmission  
linear response treatment



Spin-dependent Kitaev chain (tight binding model)

Nanocircuit charge  
susceptibility calculated with  
Keldysh Green's functions

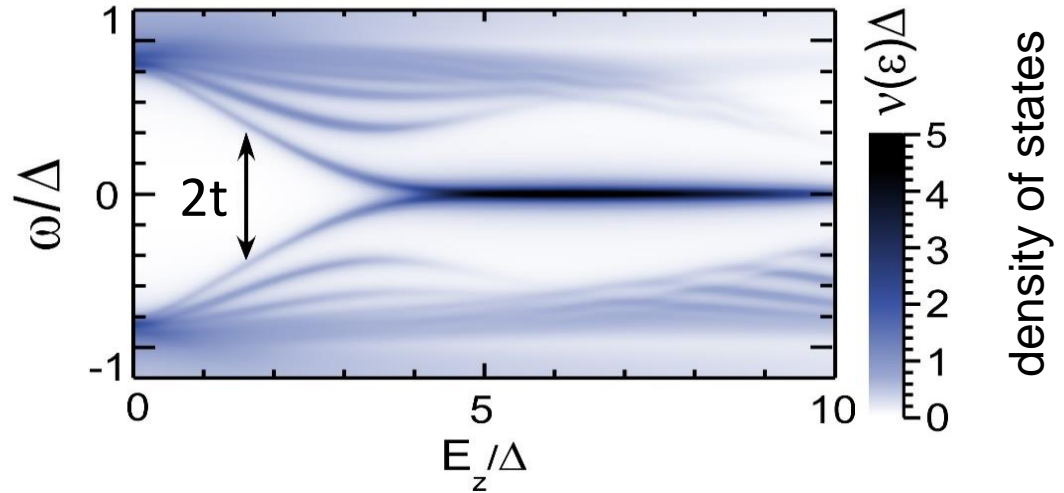
*Bruhat et al.,  
PRX 6, 021014 (2016)*



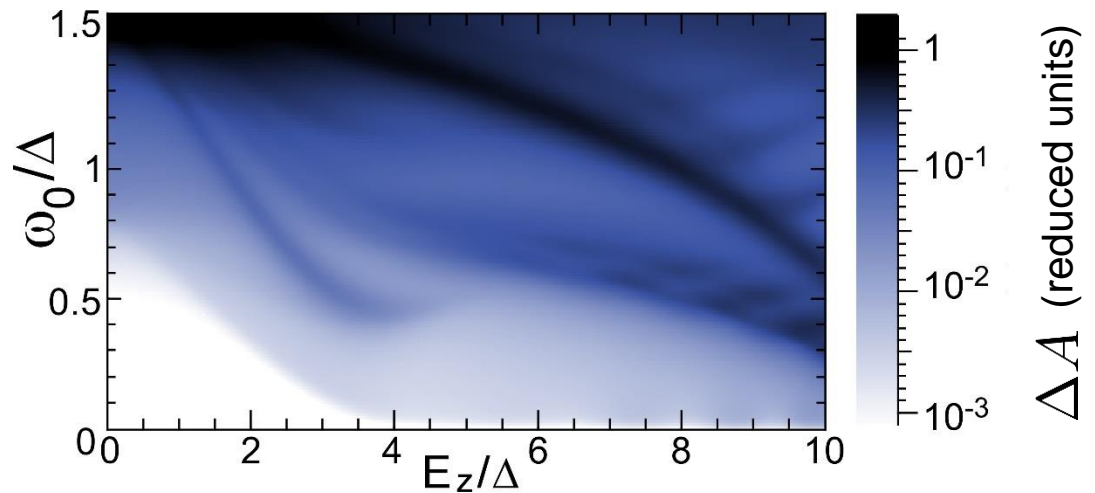
# Dissipative response of the cavity (theory)

*Dartiailh, Kontos, Douçot & Cottet, PRL 118, 126803 (2017)*

*DOS measurable  
With DC current*

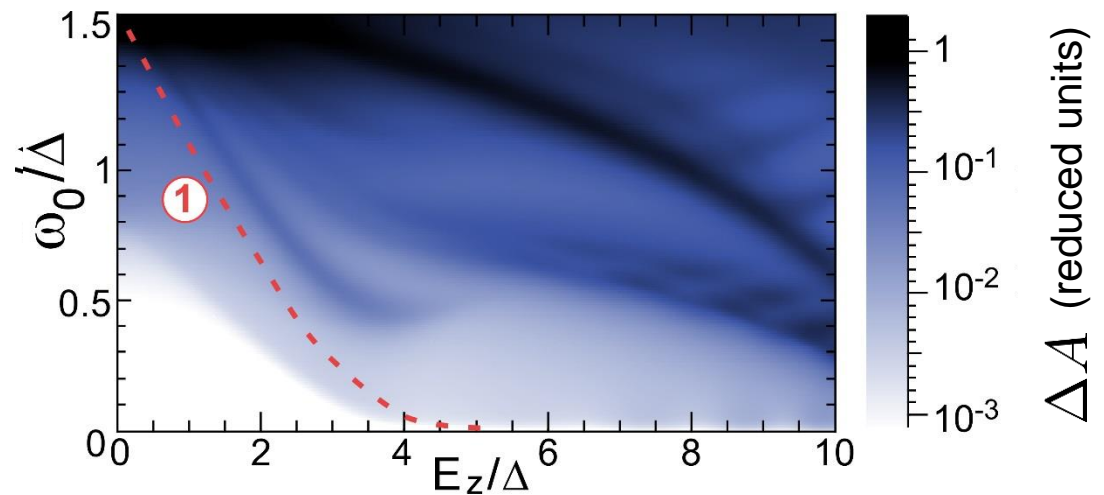
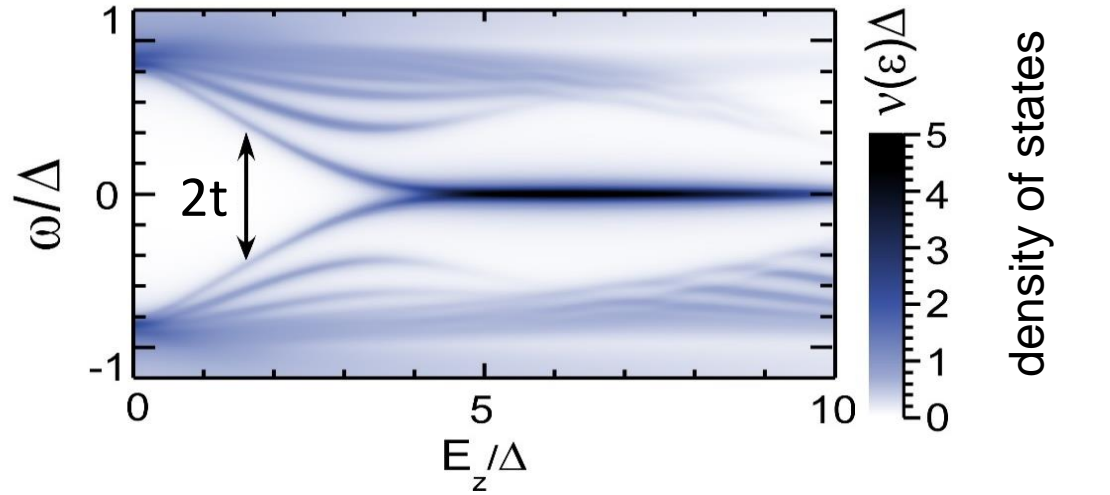
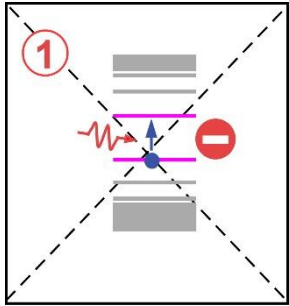


*Cavity signal*



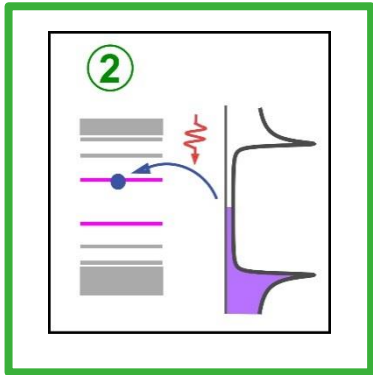
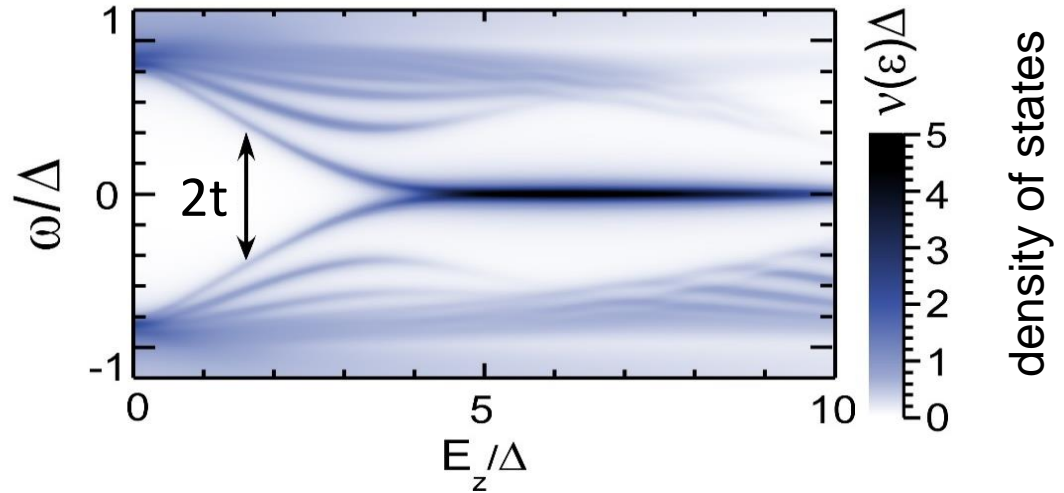
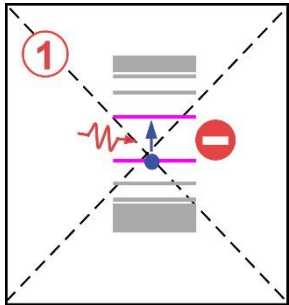
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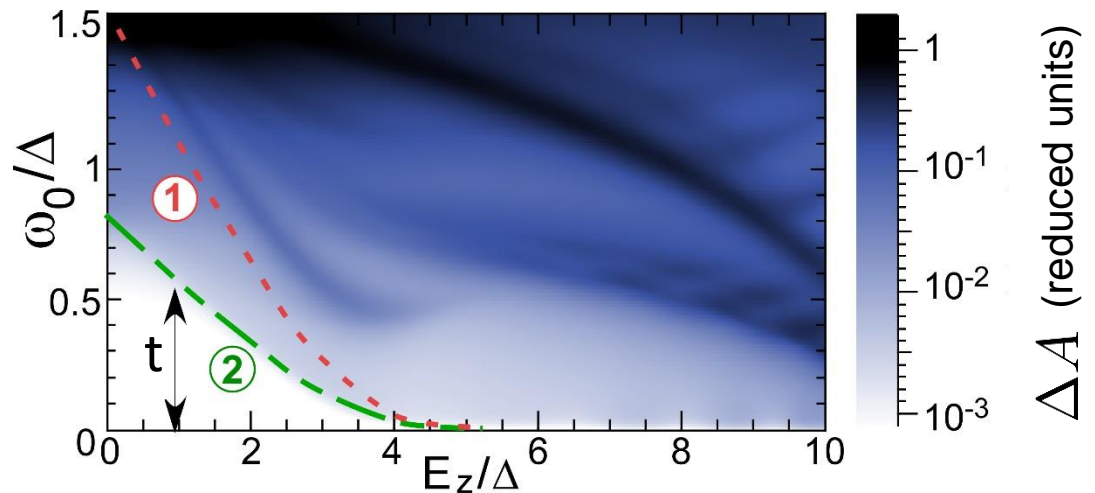


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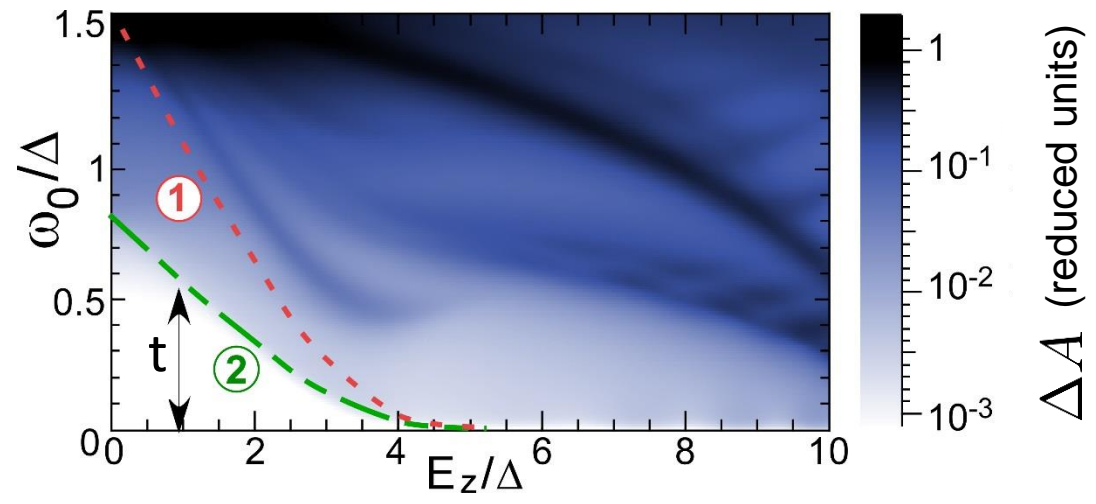
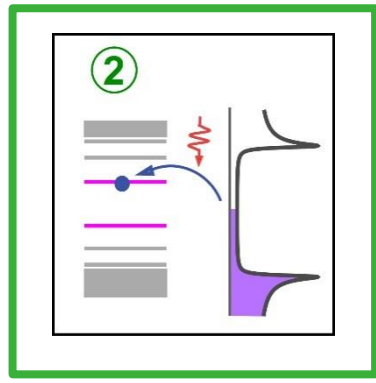
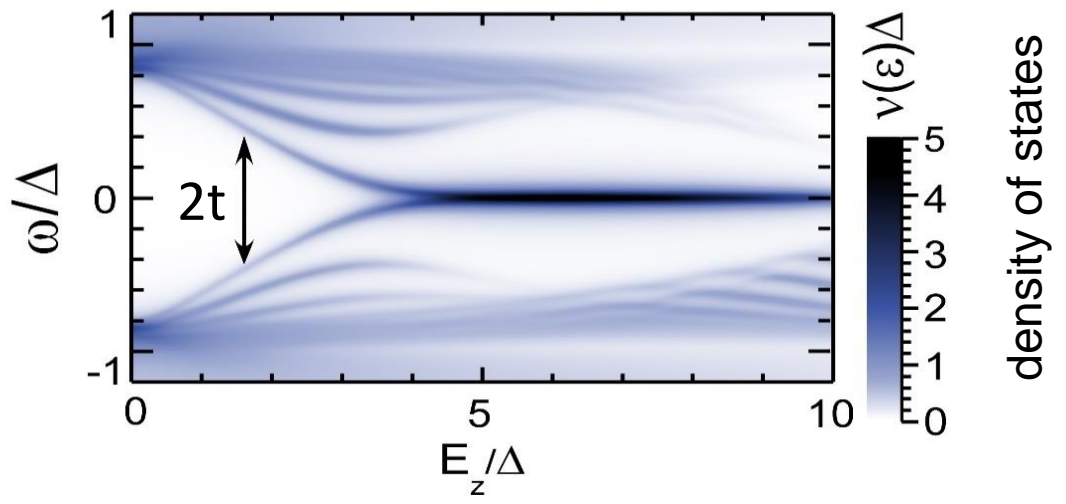
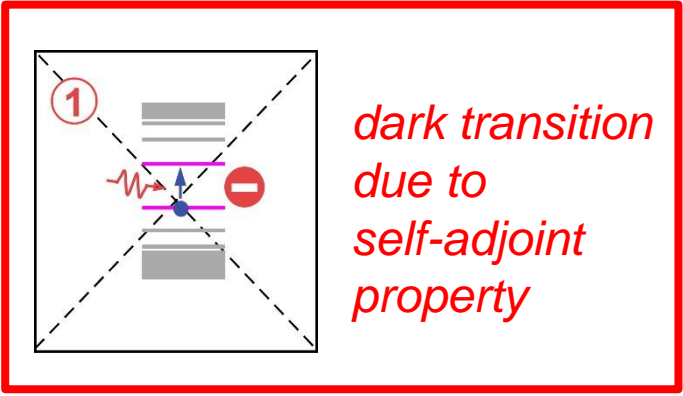


$g \neq 0$



$$\hat{H} = 2it\hat{c}_L\hat{c}_R + ig(\hat{a} + \hat{a}^\dagger)\hat{c}_L\hat{c}_R + \hbar\omega_0\hat{a}^\dagger\hat{a}$$

*Dartiailh, Kontos, Douçot & Cottet, PRL 118, 126803 (2017)*



$$\hat{H} = 2it\hat{c}_L\hat{c}_R + ig(\hat{a} + \hat{a}^\dagger)\hat{c}_L\hat{c}_R + \hbar\omega_0\hat{a}^\dagger\hat{a}$$

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# Quantum non-linear description of Mesoscopic QED

*See preprint, to appear on ArXiv in September 2019...*

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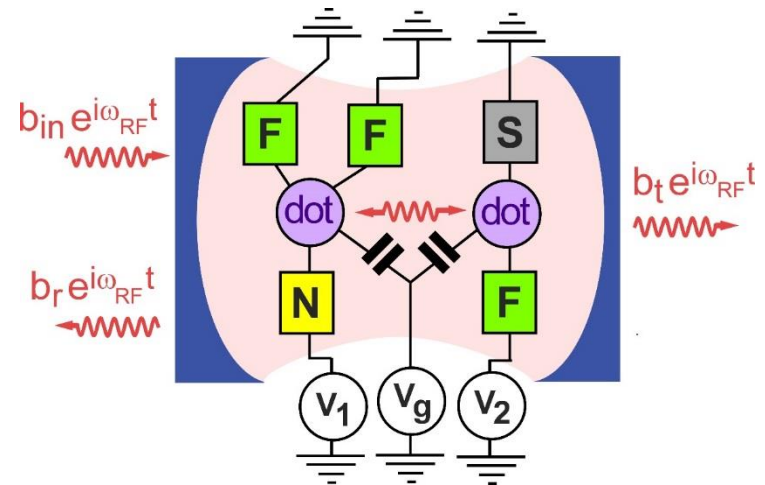


# Preparation of a photonic Schrödinger cat using dissipative tunneling

*See preprint, to appear on ArXiv in September 2019...*

# Summary of results presented

- Microwave cavities are a powerful tool to study/control mesoscopic circuits
- Mesoscopic circuits can be used to prepare non-classical cavity states



## Other cases studied so far in Paris:

- Prediction: Cooper pair splitting in a cavity: *A. Cottet et al. PRB (2012) & (2014)*
- Experiment+Theory: Kondo effect in a cavity: *M. Desjardins et al., Nature 545, 71 (2017)*

**REVIEWS: *A. Cottet, et al. J. Phys.: Condens. Matter 29 433002 (2017)***

The end